

Quantum Logic Gates and Circuits

Physical and Logical Reversibility

Classical and Quantum Logic

- Classical Logic
 - Typically is Irreversible Logic
 - Reversible Logic is a special case
 - Fanout is a powerful use-case
- Quantum Logic
 - Comparison with Classical
 - No Cloning theorem
 - Universal Quantum Logic Gates

Classical Irreversible Logic

- Theory of Nature of Computing (Church, Turing 1936)
- Universality of Primitive Operations

x	y	x AND y
0	0	0
0	1	0
1	0	0
1	1	1

x	y	x OR y
0	0	0
0	1	1
1	0	1
1	1	1

y	NOT y
0	1
1	0

Types of Reversibility

- **Logical Reversibility**

- Ability to reconstruct input from output. Circuit function is a *Bijection*.

- Bijection implies two properties: (1) one-to-one, (2) onto

- **Physical Reversibility**

- Thermodynamic entropy based arguments that relate the loss of information to an increase in dissipated heat.

- Heat dissipation during a computation is generally a sign of physical irreversibility.

Thermodynamics Concepts (Oversimplified)

- Thermodynamics
 - branch of physics that studies the effects of changes in temperature, pressure, and volume in physical systems
- Physical System
 - A region of spacetime and all entities (particles and fields) contained within it. (eg. universe, transistors, circuits, computers - defn from M. Frank)
- Entropy
 - measure of the amount of energy in a physical system that cannot be used to do work - entropy S is multiplied by a temperature to yield an amount of energy. It is a measure of the disorder and randomness present in a system. A quantitative measure of the amount of thermal energy NOT AVAILABLE to do work.

Physical Irreversibility

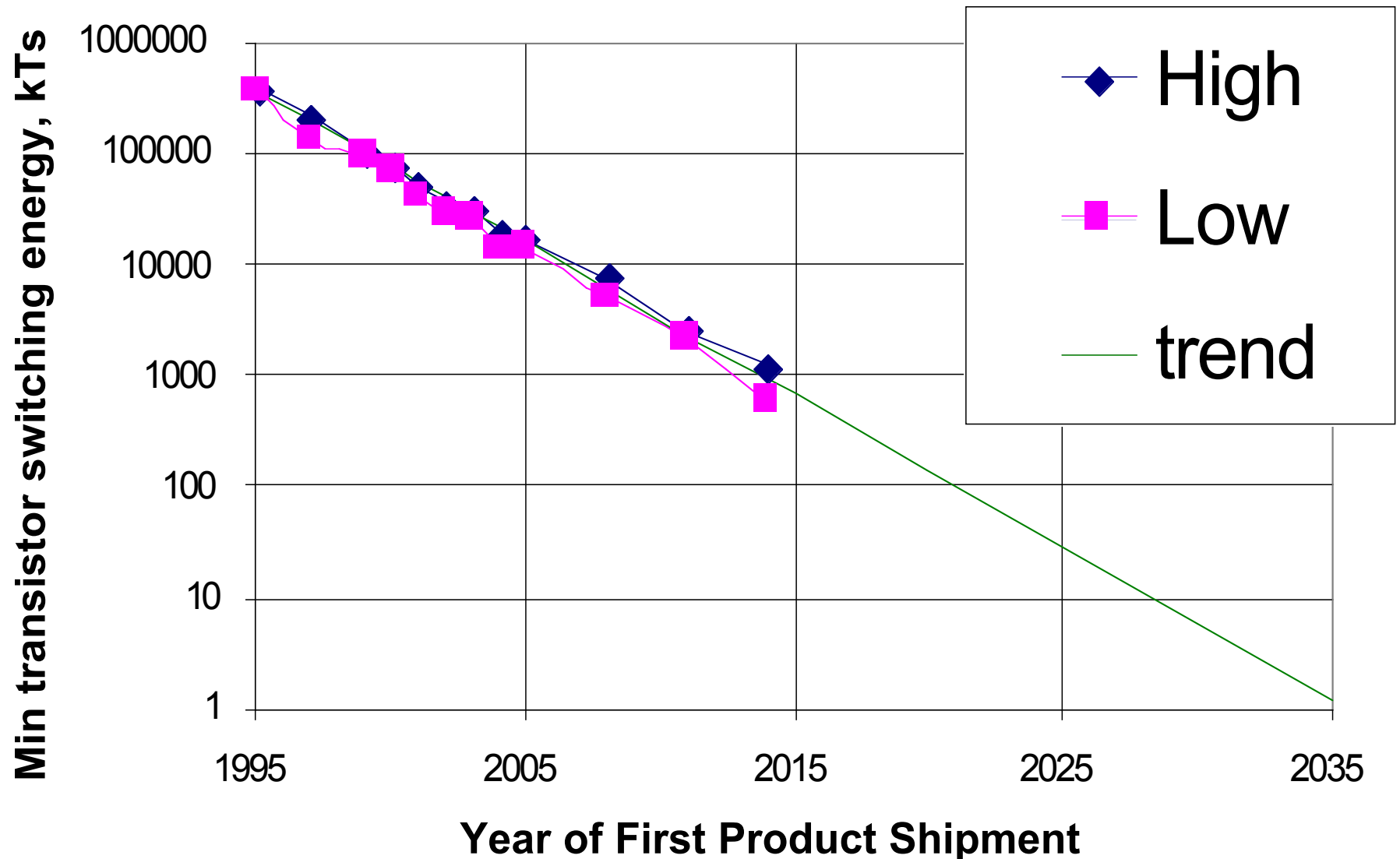
- 2-input NAND Gate
 - one output, two inputs
 - in computing an output, one input is “erased”
 - information irretrievably lost
 - change in entropy of the system is one bit of information - quantitatively this is $\ln 2$
 - conversion to energy increase of $kT \ln 2$ where k is Boltzmann's constant and T is temperature
 - corresponds to energy “lost” to heat dissipation and a sign of physical irreversibility

Developments in Reversibility

- Can a computation be accomplished in a logically reversible fashion? (unlike using a NAND gate - 1970' s)
- Must heat be dissipated during a computation?
 - Feynmann points out (1986) transistor dissipates $10^{10}kT$ joules of heat, DNA copying in a human cell dissipates $100 kT$ joules all far from $0.693 kT$ joule lower bound from erasing a single bit

Min. Transistor Switching Energy Trend*

Trend of minimum transistor switching energy



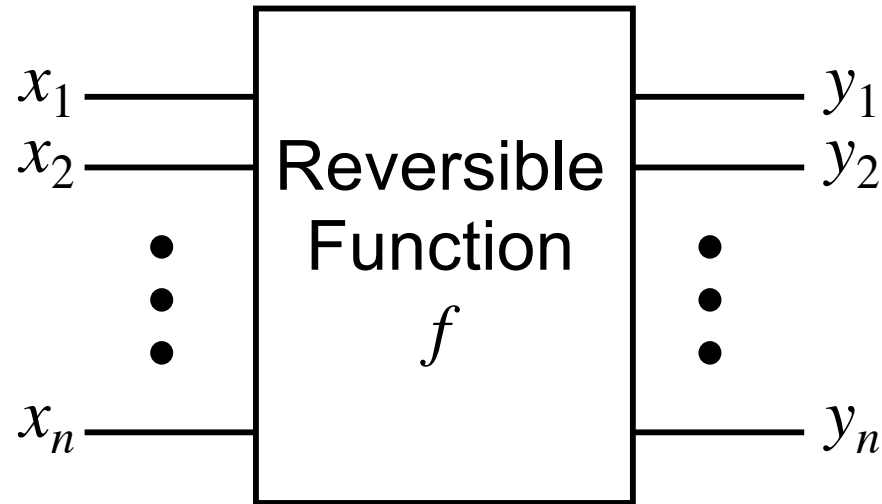
$\frac{1}{2}CV^2$ based on ITRS '99 figures for V_{dd} and minimum transistor gate capacitance. $T=300$ K

*based on chart prepared by M. Frank at Univ. of Fla.

Developments in Reversibility

- 1973 - Bennett proved that classical computation can be accomplished with no energy dissipated per computational step and with reversibility (reversible Turing machine model)
- This triggered a search for physical models for reversible classical computation
- Common Model is a discrete one-to-one binary-valued Boolean function with an equal number of inputs and outputs

Reversible Logic Circuit



- f is a bijective function
- contains symmetry that allows for other forms of representation (transformation matrix)

Quantum Logic Gates and Circuits

Single Qubit “Gates” or Transformation Operators

Quantum Logic Gates and Circuits

- Quantum Gates: Building Blocks of Quantum Computers
- Quantum Gate transforms a Quantum State to a New State
- State Transformations Performed by Gates Described by Hermitian Operators
- Matrix Describing State Transformation is Transfer Matrix

Matrices in Quantum Mechanics

- Normal Matrices, \mathbf{N} :

$$\mathbf{N}\mathbf{N}^\dagger = \mathbf{N}^\dagger\mathbf{N}$$

- Hermitian Matrices, \mathbf{H} :

$$\mathbf{H}\mathbf{H}^\dagger = \mathbf{H}^\dagger\mathbf{H}$$

$$\mathbf{H} = \mathbf{H}^\dagger$$

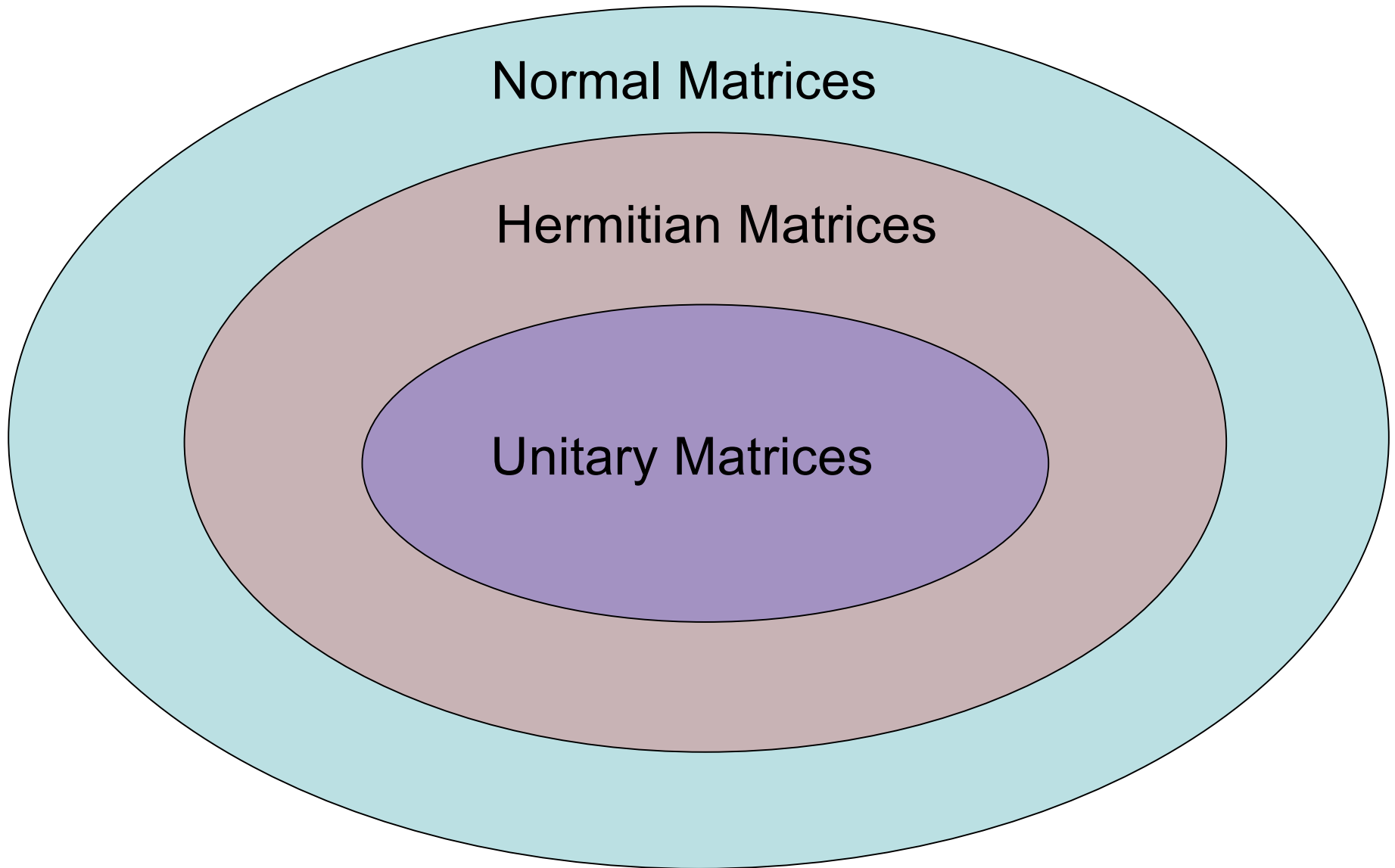
- Unitary Matrices, \mathbf{U} :

$$\mathbf{U}\mathbf{U}^\dagger = \mathbf{U}^\dagger\mathbf{U}$$

$$\mathbf{U} = \mathbf{U}^\dagger$$

$$\mathbf{U}^{-1} = \mathbf{U}^\dagger$$

Matrices in Quantum Mechanics



Practice Problems

- Given that if a square matrix \mathbf{A} satisfies $\mathbf{A}^{-1}=\mathbf{A}^\dagger$, then it must also be Hermitian (*i.e.*, that $\mathbf{A}=\mathbf{A}^\dagger$).
- Prove that the eigenvalues of a Hermitian matrix are always real-valued.
- Prove that the eigenvalues of a Unitary matrix always have unity magnitude (*i.e.*, their norm is always one).
- Prove that the non-trivial eigenvectors of a Hermitian matrix are always orthogonal with one another.

Quantum Logic Gate Matrices

- Matrices are Unitary

$$UU^\dagger = U^\dagger U = \mathbf{I}$$

$$U^{-1} = U^\dagger$$

- Transformations are Reversible

- Unitary Transformations

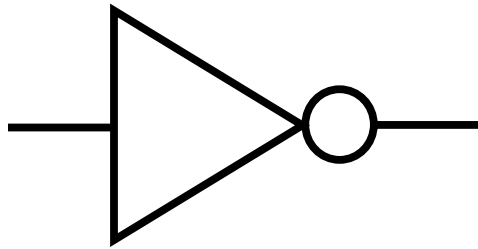
Corresponds to:

- Length Preservation
- Information Preserving Rotation in Vector Space

Quantum Logic Gate Matrices

- Gates have equal number of Inputs and Outputs
- Input/Output States of Quantum Gate Described by Vectors in Hilbert Space
- 1-qubit: $\mathbb{H}^2 \quad \{|0\rangle, |1\rangle\}$
- 2-qubit: $\mathbb{H}^4 = \mathbb{H}^2 \otimes \mathbb{H}^2 \quad \{|00\rangle, |01\rangle, |10\rangle, |11\rangle\}$
- 3-qubit: $\mathbb{H}^8 = \mathbb{H}^2 \otimes \mathbb{H}^2 \otimes \mathbb{H}^2$
 $\{|000\rangle, |001\rangle, |010\rangle, |011\rangle, |100\rangle, |101\rangle, |110\rangle, |111\rangle\}$

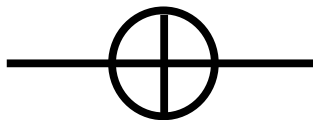
Classical Reversible Gates/Operators



Classic Symbol
Physically Irreversible
Logically Reversible

In	Out
0	1
1	0

Classic Truth Table
Notion of Inputs/Outputs

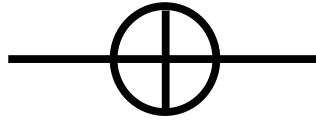


NOT
Symbol for Reversible
NOT Gate
(aka Pauli-X)

$$\begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$$

Matrix Representation
of Pauli-X Gate
Functionality

Reversible NOT Gate

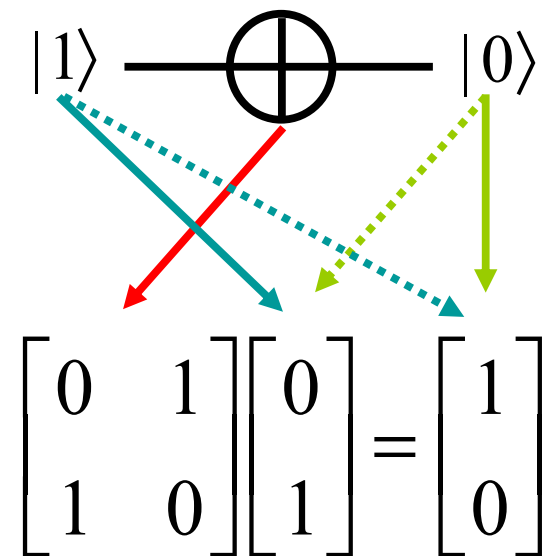
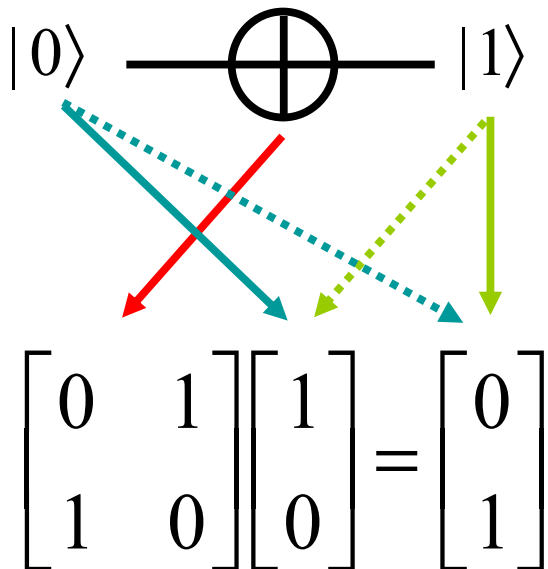


NOT

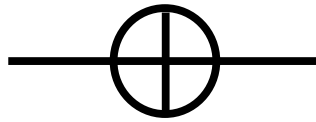
Symbol for Reversible
NOT Gate

$$\begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$$

Matrix Representation
of NOT Gate
Functionality



Reversible NOT Gate



NOT

Symbol for Reversible
NOT Gate

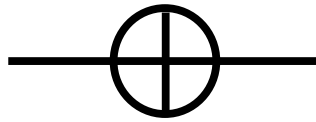
$$\mathbf{A} = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$$

Matrix Representation
of NOT Gate
Functionality

$$\mathbf{A}^\dagger = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$$

$$\mathbf{A}^\dagger \mathbf{A} = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} = \mathbf{I}$$

Reversible NOT Gate



NOT

Symbol for Reversible
NOT Gate

$$\begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} = |0\rangle\langle 1| + |1\rangle\langle 0|$$

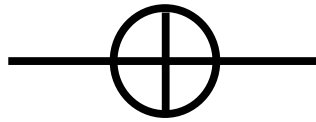
Matrix Representation
of NOT Gate
Functionality

$$\begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \begin{bmatrix} 1 \\ 0 \end{bmatrix} = \begin{bmatrix} 0 \\ 1 \end{bmatrix}$$

Dirac Notation Example

$$\begin{aligned} & (|0\rangle\langle 1| + |1\rangle\langle 0|)|0\rangle \\ &= |0\rangle\langle 1|0\rangle + |1\rangle\langle 0|0\rangle \\ &= |0\rangle(0) + |1\rangle(1) = |1\rangle \end{aligned}$$

Reversible NOT Gate



NOT

Symbol for Reversible
NOT Gate

$$\begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} = |0\rangle\langle 1| + |1\rangle\langle 0|$$

Matrix Representation
of NOT Gate
Functionality

“INPUTS”

“OUTPUTS”

$$\begin{array}{c} |0\rangle \\ |1\rangle \end{array} \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$$

The diagram shows the NOT gate's functionality. On the left, the output states $|0\rangle$ and $|1\rangle$ are listed vertically. A blue arrow labeled “OUTPUTS” points to these states. On the right, the input states $|0\rangle$ and $|1\rangle$ are listed vertically. A red arrow labeled “INPUTS” points to these states. A permutation matrix is shown between the input and output states, mapping $|0\rangle$ to $|1\rangle$ and $|1\rangle$ to $|0\rangle$.

Permutation Matrix
Transformations of
Qubits

Derivation of \mathbf{I}_2 or σ_0

- This Operator Performs an Identity Transformation of the Basis Vectors:

$$|0\rangle \mapsto |0\rangle \qquad |1\rangle \mapsto |1\rangle$$

- Computed as:

$$\sigma_0 = \mathbf{I} = |0\rangle\langle 0| + |1\rangle\langle 1|$$

$$\sigma_0 = \mathbf{I} = \begin{bmatrix} 1 \\ 0 \end{bmatrix} \otimes \begin{bmatrix} 1 & 0 \end{bmatrix} + \begin{bmatrix} 0 \\ 1 \end{bmatrix} \otimes \begin{bmatrix} 0 & 1 \end{bmatrix}$$

$$\sigma_0 = \mathbf{I} = \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix} + \begin{bmatrix} 0 & 0 \\ 0 & 1 \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$$

Derivation of X or σ_X

- This Operator “Flips” or “Negates” a Qubit:

$$|0\rangle \mapsto |1\rangle \quad |1\rangle \mapsto |0\rangle$$

- Computed as:

$$\sigma_1 = \mathbf{X} = |0\rangle\langle 1| + |1\rangle\langle 0|$$

$$\sigma_1 = \mathbf{X} = \begin{bmatrix} 1 \\ 0 \end{bmatrix} \otimes \begin{bmatrix} 0 & 1 \end{bmatrix} + \begin{bmatrix} 0 \\ 1 \end{bmatrix} \otimes \begin{bmatrix} 1 & 0 \end{bmatrix}$$

$$\sigma_1 = \mathbf{X} = \begin{bmatrix} 0 & 1 \\ 0 & 0 \end{bmatrix} + \begin{bmatrix} 0 & 0 \\ 1 & 0 \end{bmatrix} = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$$

Derivation of Y or σ_Y

- This Operator Multiplies a Qubit by i (90-degree phase shift) then “Flips” or “Negates” it:

$$|0\rangle \mapsto i|1\rangle \quad |1\rangle \mapsto -i|0\rangle$$

- Computed as:

$$\sigma_2 = Y = -i|0\rangle\langle 1| + i|1\rangle\langle 0|$$

$$\sigma_2 = Y = -i \begin{bmatrix} 1 \\ 0 \end{bmatrix} \otimes \begin{bmatrix} 0 & 1 \end{bmatrix} + i \begin{bmatrix} 0 \\ 1 \end{bmatrix} \otimes \begin{bmatrix} 1 & 0 \end{bmatrix}$$

$$\sigma_2 = Y = -i \begin{bmatrix} 0 & 1 \\ 0 & 0 \end{bmatrix} + i \begin{bmatrix} 0 & 0 \\ 1 & 0 \end{bmatrix} = \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix}$$

Derivation of \mathbf{Z} or σ_z

- This Operator is an Identity with a 180 Degree Phase Shift Operation:

$$|0\rangle \mapsto |0\rangle \quad |1\rangle \mapsto -|1\rangle$$

- Computed as:

$$\sigma_3 = \mathbf{Z} = |0\rangle\langle 0| - |1\rangle\langle 1|$$

$$\sigma_3 = \mathbf{Z} = \begin{bmatrix} 1 \\ 0 \end{bmatrix} \otimes \begin{bmatrix} 1 & 0 \end{bmatrix} - \begin{bmatrix} 0 \\ 1 \end{bmatrix} \otimes \begin{bmatrix} 0 & 1 \end{bmatrix}$$

$$\sigma_3 = \mathbf{Z} = \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix} - \begin{bmatrix} 0 & 0 \\ 0 & 1 \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$$

Pauli Operator Examples

- Assume the Following:

$$|\varphi\rangle = \sigma_i |\psi\rangle = \sigma_i [\alpha_0 |0\rangle + \alpha_1 |1\rangle]$$

$$|\varphi\rangle = \sigma_0 |\psi\rangle = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} \alpha_0 \\ \alpha_1 \end{bmatrix} = \begin{bmatrix} \alpha_0 \\ \alpha_1 \end{bmatrix} = \alpha_0 |0\rangle + \alpha_1 |1\rangle$$

$$|\varphi\rangle = \sigma_1 |\psi\rangle = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \begin{bmatrix} \alpha_0 \\ \alpha_1 \end{bmatrix} = \begin{bmatrix} \alpha_1 \\ \alpha_0 \end{bmatrix} = \alpha_1 |0\rangle + \alpha_0 |1\rangle$$

$$|\varphi\rangle = \sigma_2 |\psi\rangle = \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix} \begin{bmatrix} \alpha_0 \\ \alpha_1 \end{bmatrix} = i \begin{bmatrix} -\alpha_1 \\ \alpha_0 \end{bmatrix} = -i\alpha_1 |0\rangle + i\alpha_0 |1\rangle$$

$$|\varphi\rangle = \sigma_3 |\psi\rangle = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \begin{bmatrix} \alpha_0 \\ \alpha_1 \end{bmatrix} = \begin{bmatrix} \alpha_0 \\ -\alpha_1 \end{bmatrix} = \alpha_0 |0\rangle - \alpha_1 |1\rangle$$

Quantum Logic Gates and Circuits

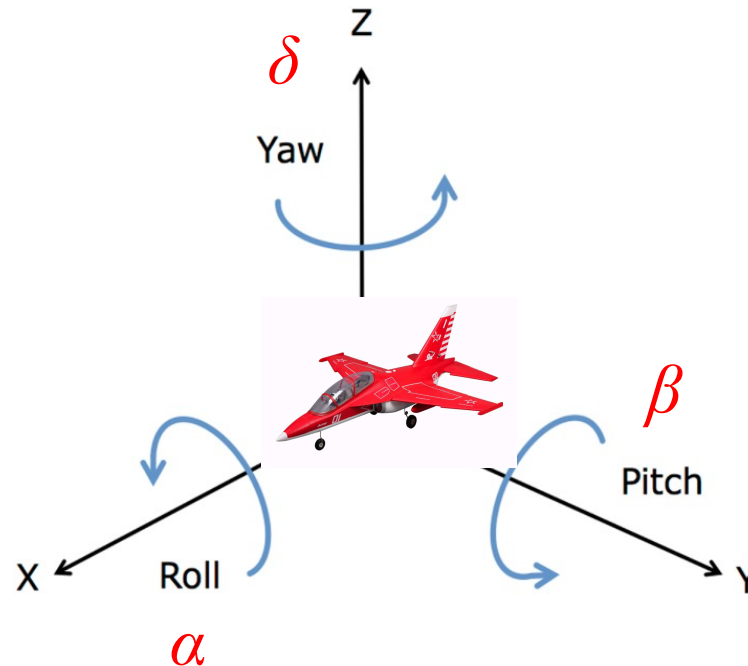
Single Qubit “Gates” or Transformation Operators

GEOMETRIC INTERPRETATIONS

Orientation and Definitions from Aerospace

- 3-D Rotation Matrices are Elements of the non-Abelian $SO(3)$ Group with Direct Matrix Multiplication as the Group Product Operation

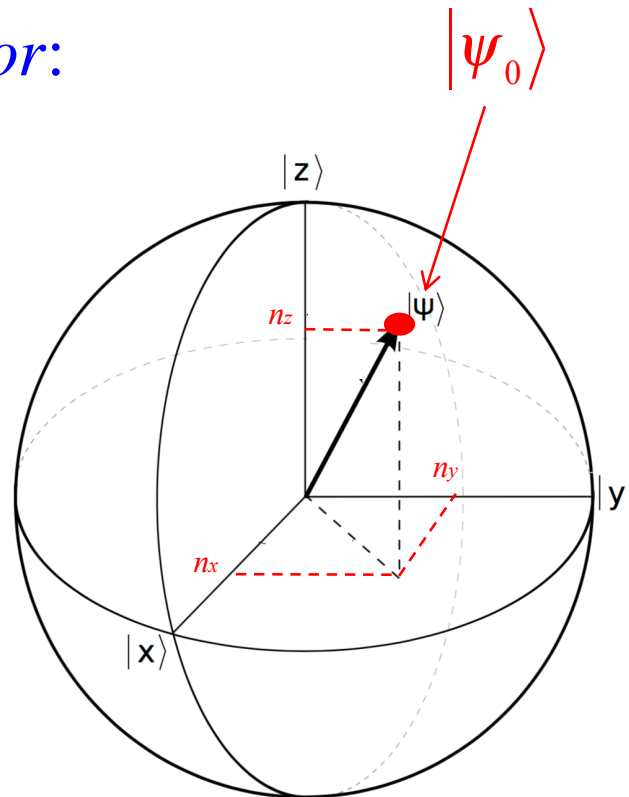
$$\mathbf{R}(\omega) = \mathbf{R}_X(\alpha)\mathbf{R}_Y(\beta)\mathbf{R}_Z(\delta)$$



Qubit on the Bloch Sphere Surface (Cartesian or rectangular)

Cartesian Bloch Vector:

$$\hat{\mathbf{n}}_{|\psi_0\rangle} = (n_x, n_y, n_z)^T$$



- Location on Surface of Sphere in Rectangular Coordinates:

$$|\psi_0\rangle = n_x |x\rangle + n_y |y\rangle + n_z |z\rangle$$

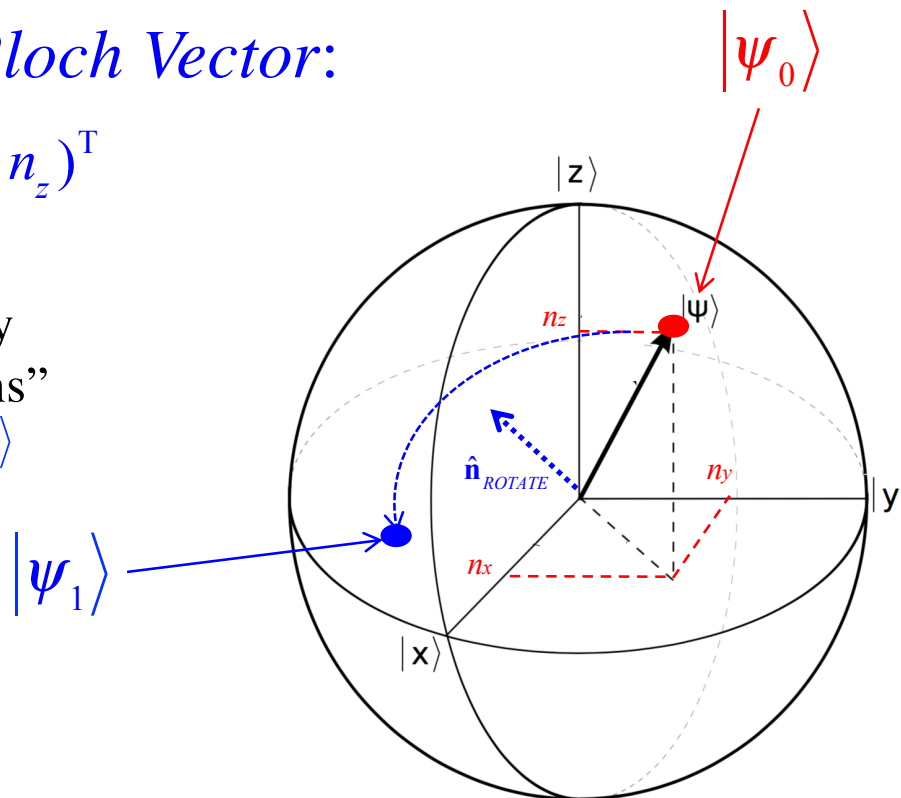
- Non-traditional Basis, $(|x\rangle, |y\rangle, |z\rangle)$

Qubit Rotation on the Bloch Sphere Surface (Cartesian or rectangular)

Cartesian Bloch Vector:

$$\hat{\mathbf{n}}_{|\psi_0\rangle} = (n_x, n_y, n_z)^T$$

There are many
Different “Paths”
from $|\psi_0\rangle$ to $|\psi_1\rangle$



$$|\psi_0\rangle \rightarrow |\psi_1\rangle$$

$$|\psi_1\rangle = \mathbf{R} |\psi_0\rangle$$

- Geodesic or “Great Circle” Path is Shown with Blue Dashed Line
- Shortest Path Rotation
- Great Circle Rotation is Rotation of Angle ω about an Axis of Rotation: $\hat{\mathbf{n}}_{ROTATE}$

- \mathbf{R} is Generally Defined about a General “Axis of Rotation”
 - $\hat{\mathbf{n}}_{ROTATE} = \hat{\mathbf{n}}_{|\psi_0\rangle} \times \hat{\mathbf{n}}_{|\psi_1\rangle}$ (3D Vector Cross Product, Right-hand Rule)
- \mathbf{R} is a Single 3×3 Rotation Matrix – Unconventional in QIS
 - $\mathbf{R}(\hat{\mathbf{n}}_{ROTATE}, \omega)$ (Explicit Formula in Backup Slides, Simpler using Quaternions)

Rotation Operator, \mathbf{R} , Product of Elemental Rotations

- \mathbf{R} can be Defined as a Product of Rotations about Three Axes, (x,y,z) :

$$|\psi_0\rangle \rightarrow |\psi_1\rangle$$

$$|\psi_1\rangle = \mathbf{R}|\psi_0\rangle$$

- Two forms of Elemental Rotation Decompositions
 - Product of 3 Rotation Matrices about Each Axis: Tait-Bryan Rotations
 - Product of 3 Rotation Matrices about Two of the Axes: Euler Rotations

- Tait-Bryan Forms:

$$\mathbf{R} = \mathbf{R}_x(\alpha)\mathbf{R}_y(\beta)\mathbf{R}_z(\delta), \mathbf{R} = \mathbf{R}_y(\beta)\mathbf{R}_z(\delta)\mathbf{R}_x(\alpha), \mathbf{R} = \mathbf{R}_z(\delta)\mathbf{R}_x(\alpha)\mathbf{R}_y(\beta),$$

$$\mathbf{R} = \mathbf{R}_x(\alpha)\mathbf{R}_z(\delta)\mathbf{R}_y(\beta), \mathbf{R} = \mathbf{R}_z(\delta)\mathbf{R}_y(\beta)\mathbf{R}_x(\alpha), \mathbf{R} = \mathbf{R}_y(\beta)\mathbf{R}_x(\alpha)\mathbf{R}_z(\delta)$$

- Euler Forms:

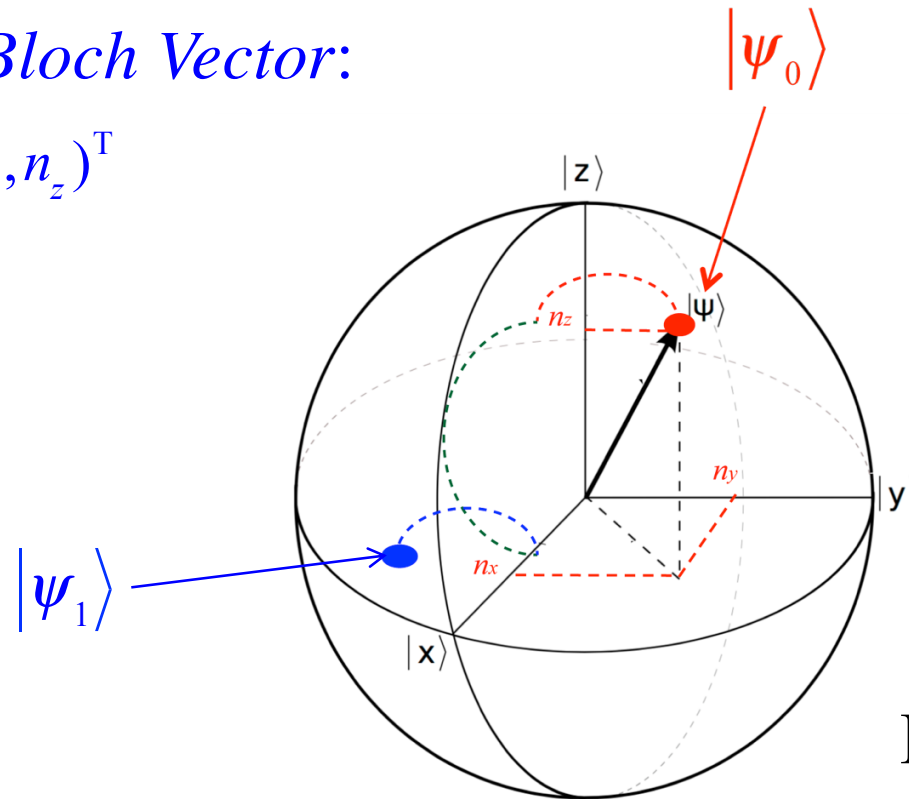
$$\mathbf{R} = \mathbf{R}_z(\delta_1)\mathbf{R}_x(\alpha)\mathbf{R}_z(\delta_2), \mathbf{R} = \mathbf{R}_x(\alpha_1)\mathbf{R}_y(\beta)\mathbf{R}_x(\alpha_2), \mathbf{R} = \mathbf{R}_y(\beta_1)\mathbf{R}_z(\delta)\mathbf{R}_y(\beta_2),$$

$$\mathbf{R} = \mathbf{R}_z(\delta_1)\mathbf{R}_y(\beta)\mathbf{R}_z(\delta_2), \mathbf{R} = \mathbf{R}_x(\alpha_1)\mathbf{R}_z(\delta)\mathbf{R}_x(\alpha_2), \mathbf{R} = \mathbf{R}_y(\beta_1)\mathbf{R}_x(\alpha)\mathbf{R}_y(\beta_2)$$

Qubit Rotation on the Bloch Sphere Surface (Elemental Rotation)

Cartesian Bloch Vector:

$$\hat{\mathbf{n}}_{|\psi_0\rangle} = (n_x, n_y, n_z)^T$$



$$|\psi_0\rangle \rightarrow |\psi_1\rangle$$

$$|\psi_1\rangle = \mathbf{R} |\psi_0\rangle$$

eg., Tait-Bryan form: (z-y-x)

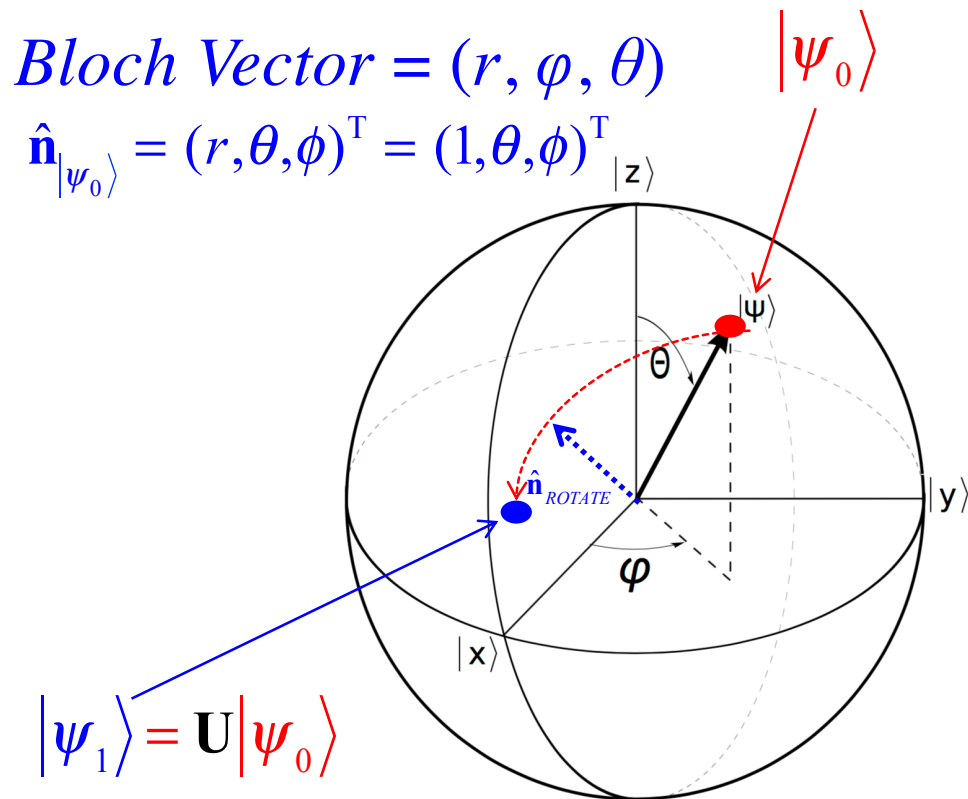
$$\mathbf{R} = \mathbf{R}_X(\alpha) \mathbf{R}_Y(\beta) \mathbf{R}_Z(\delta)$$

$$\mathbf{R}_X(\alpha) = \begin{bmatrix} 1 & 0 & 0 \\ 0 & \cos \alpha & -\sin \alpha \\ 0 & \sin \alpha & \cos \alpha \end{bmatrix}$$

$$\mathbf{R}_Y(\beta) = \begin{bmatrix} \cos \beta & 0 & \sin \beta \\ 0 & 1 & 0 \\ -\sin \beta & 0 & \cos \beta \end{bmatrix}$$

$$\mathbf{R}_Z(\delta) = \begin{bmatrix} \cos \delta & -\sin \delta & 0 \\ \sin \delta & \cos \delta & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

Moving on the Bloch Sphere Surface (Spherical)



$$r = \|\psi_0\| = \langle \psi_0 | \psi_0 \rangle = \sqrt{n_x^2 + n_y^2 + n_z^2} = 1$$

$$\varphi = \tan^{-1}\left(\frac{n_y}{n_x}\right) \quad \theta = \cos^{-1}\left(\frac{n_z}{\sqrt{n_x^2 + n_y^2 + n_z^2}}\right)$$

$$n_x = \langle \psi_0 | \psi_0 \rangle \sin \theta \cos \varphi = \sin \theta \cos \varphi$$

$$n_y = \langle \psi_0 | \psi_0 \rangle \sin \theta \sin \varphi = \sin \theta \sin \varphi$$

$$n_z = \langle \psi_0 | \psi_0 \rangle \cos \theta = \cos \theta$$

- Location on Surface of Sphere in Rectangular Coordinates:

$$|\psi_0\rangle = n_x|x\rangle + n_y|y\rangle + n_z|z\rangle = (\sin \theta \cos \varphi)|x\rangle + (\sin \theta \sin \varphi)|y\rangle + (\cos \theta)|z\rangle$$

$$= e^{i\gamma} \left[\cos\left(\frac{\theta}{2}\right)|0\rangle + e^{i\phi} \sin\left(\frac{\theta}{2}\right)|1\rangle \right]$$

*Global Phase γ is not Important
in Terms of Information Content*

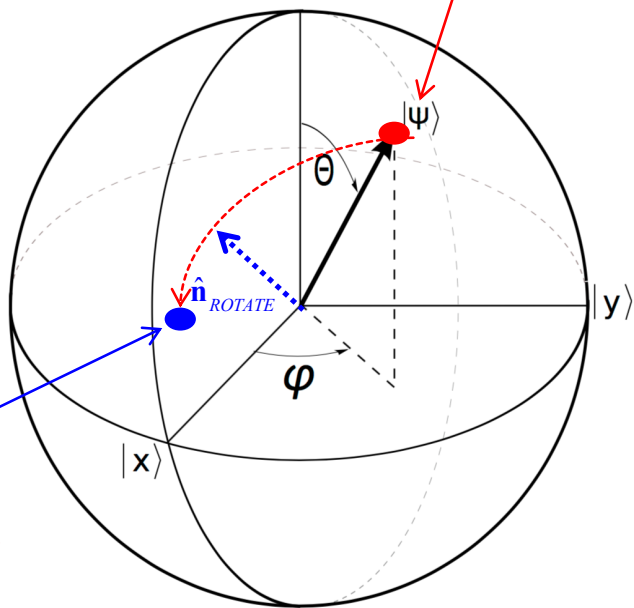
Moving on the Bloch Sphere Surface (Spherical)

Bloch Vector = (r, φ, θ)

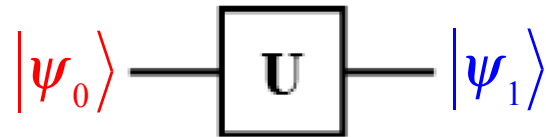
$$\hat{\mathbf{n}}_{|\psi_0\rangle} = (r, \theta, \phi)^T = (1, \theta, \phi)^T_{|z\rangle}$$

$|\psi_0\rangle$

$|\psi_0\rangle$ Evolves in Time to $|\psi_1\rangle$



$$|\psi_1\rangle = \mathbf{U}|\psi_0\rangle$$



TIME

$$\mathbf{U}_{3D} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & u_{00} & u_{01} \\ 0 & u_{10} & u_{11} \end{bmatrix} = \mathbf{R}_X(\alpha)\mathbf{R}_Y(\beta)\mathbf{R}_Z(\delta)$$

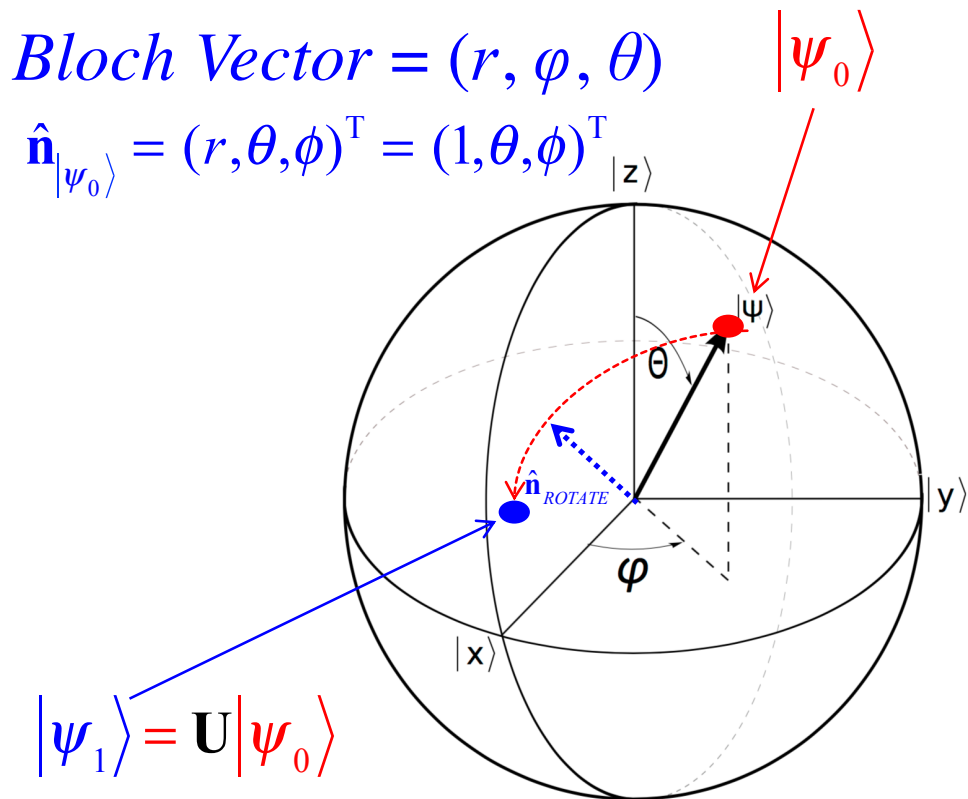
Euler or Tait-Bryan Decompositions

$$\mathbf{R}_X(\alpha) = \begin{bmatrix} 1 & 0 & 0 \\ 0 & \cos\frac{\alpha}{2} & -i\sin\frac{\alpha}{2} \\ 0 & -i\sin\frac{\alpha}{2} & \cos\frac{\alpha}{2} \end{bmatrix}$$

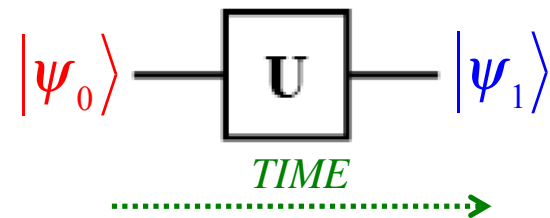
$$\mathbf{R}_Y(\beta) = \begin{bmatrix} 1 & 0 & 0 \\ 0 & \cos\frac{\beta}{2} & -\sin\frac{\beta}{2} \\ 0 & \sin\frac{\beta}{2} & \cos\frac{\beta}{2} \end{bmatrix}$$

$$\mathbf{R}_Z(\delta) = \begin{bmatrix} 1 & 0 & 0 \\ 0 & e^{-i\frac{\delta}{2}} & 0 \\ 0 & 0 & e^{i\frac{\delta}{2}} \end{bmatrix}$$

Moving on the Bloch Sphere Surface (Spherical)



$|\psi_0\rangle$ Evolves in Time to $|\psi_1\rangle$



$$\mathbf{U} = \begin{bmatrix} u_{00} & u_{01} \\ u_{10} & u_{11} \end{bmatrix} = \mathbf{R}_X(\alpha)\mathbf{R}_Y(\beta)\mathbf{R}_Z(\delta)$$

Euler or Tait-Bryan
Decompositions

$$\mathbf{R}_X(\alpha) = \begin{bmatrix} \cos\frac{\alpha}{2} & -i\sin\frac{\alpha}{2} \\ -i\sin\frac{\alpha}{2} & \cos\frac{\alpha}{2} \end{bmatrix}$$

$$\mathbf{R}_Y(\beta) = \begin{bmatrix} \cos\frac{\beta}{2} & -\sin\frac{\beta}{2} \\ \sin\frac{\beta}{2} & \cos\frac{\beta}{2} \end{bmatrix}$$

$$\mathbf{R}_Z(\delta) = \begin{bmatrix} e^{-i\frac{\delta}{2}} & 0 \\ 0 & e^{i\frac{\delta}{2}} \end{bmatrix}$$

Rotation Identities (Euler Decompositions)

Rotate γ Degrees
about **X**-axis

$$\mathbf{R}_X(\gamma) = \mathbf{R}_Z\left(-\frac{\pi}{2}\right)\mathbf{R}_Y(\gamma)\mathbf{R}_Z\left(\frac{\pi}{2}\right)$$

$$\mathbf{R}_X(\gamma) = \mathbf{R}_Y\left(\frac{\pi}{2}\right)\mathbf{R}_Z(\gamma)\mathbf{R}_Y\left(-\frac{\pi}{2}\right)$$

Rotate γ Degrees
about **Y**-axis

$$\mathbf{R}_Y(\gamma) = \mathbf{R}_X\left(-\frac{\pi}{2}\right)\mathbf{R}_Z(\gamma)\mathbf{R}_X\left(\frac{\pi}{2}\right)$$

$$\mathbf{R}_Y(\gamma) = \mathbf{R}_Z\left(\frac{\pi}{2}\right)\mathbf{R}_X(\gamma)\mathbf{R}_Z\left(-\frac{\pi}{2}\right)$$

Rotate γ Degrees
about **Z**-axis

$$\mathbf{R}_Z(\gamma) = \mathbf{R}_X\left(\frac{\pi}{2}\right)\mathbf{R}_Y(\gamma)\mathbf{R}_X\left(-\frac{\pi}{2}\right)$$

$$\mathbf{R}_Z(\gamma) = \mathbf{R}_Y\left(-\frac{\pi}{2}\right)\mathbf{R}_X(\gamma)\mathbf{R}_Y\left(\frac{\pi}{2}\right)$$

- **Notes:**

- These Examples Restrict Decomposed Euler Angles to $(\pm(\pi/2), \gamma, \pm(\pi/2))$
- Angles are Restricted to the Interval $[-\pi, +\pi]$
- Rotations of $\pm 2m\pi$ for $m=0,1,2,\dots$ are Omitted,
 $\mathbf{R}_k(\pm 2m\pi) = \mathbf{R}_k(0) = \mathbf{I}$

Rotation Identities (Tait-Bryan Decompositions)

Rotate γ Degrees
about **X**-axis

Rotate γ Degrees
about **Y**-axis

Rotate γ Degrees
about **Z**-axis

$\mathbf{R}_X(\gamma) = \mathbf{R}_X(\alpha)\mathbf{R}_Y(\beta)\mathbf{R}_Z(\delta)$	$\mathbf{R}_Y(\gamma) = \mathbf{R}_X(\alpha)\mathbf{R}_Y(\beta)\mathbf{R}_Z(\delta)$	$\mathbf{R}_Z(\gamma) = \mathbf{R}_X(\alpha)\mathbf{R}_Y(\beta)\mathbf{R}_Z(\delta)$
$\mathbf{R}_X(\gamma) = \mathbf{R}_Y(\beta)\mathbf{R}_Z(\delta)\mathbf{R}_X(\alpha)$	$\mathbf{R}_Y(\gamma) = \mathbf{R}_Y(\beta)\mathbf{R}_Z(\delta)\mathbf{R}_X(\alpha)$	$\mathbf{R}_Z(\gamma) = \mathbf{R}_Y(\beta)\mathbf{R}_Z(\delta)\mathbf{R}_X(\alpha)$
$\mathbf{R}_X(\gamma) = \mathbf{R}_Z(\delta)\mathbf{R}_X(\alpha)\mathbf{R}_Y(\beta)$	$\mathbf{R}_Y(\gamma) = \mathbf{R}_Z(\delta)\mathbf{R}_X(\alpha)\mathbf{R}_Y(\beta)$	$\mathbf{R}_Z(\gamma) = \mathbf{R}_Z(\delta)\mathbf{R}_X(\alpha)\mathbf{R}_Y(\beta)$
$\mathbf{R}_X(\gamma) = \mathbf{R}_X(\alpha)\mathbf{R}_Z(\delta)\mathbf{R}_Y(\beta)$	$\mathbf{R}_Y(\gamma) = \mathbf{R}_X(\alpha)\mathbf{R}_Z(\delta)\mathbf{R}_Y(\beta)$	$\mathbf{R}_Z(\gamma) = \mathbf{R}_X(\alpha)\mathbf{R}_Z(\delta)\mathbf{R}_Y(\beta)$
$\mathbf{R}_X(\gamma) = \mathbf{R}_Z(\delta)\mathbf{R}_Y(\beta)\mathbf{R}_X(\alpha)$	$\mathbf{R}_Y(\gamma) = \mathbf{R}_Z(\delta)\mathbf{R}_Y(\beta)\mathbf{R}_X(\alpha)$	$\mathbf{R}_Z(\gamma) = \mathbf{R}_Z(\delta)\mathbf{R}_Y(\beta)\mathbf{R}_X(\alpha)$
$\mathbf{R}_X(\gamma) = \mathbf{R}_Y(\beta)\mathbf{R}_X(\alpha)\mathbf{R}_Z(\delta)$	$\mathbf{R}_Y(\gamma) = \mathbf{R}_Y(\beta)\mathbf{R}_X(\alpha)\mathbf{R}_Z(\delta)$	$\mathbf{R}_Z(\gamma) = \mathbf{R}_Y(\beta)\mathbf{R}_X(\alpha)\mathbf{R}_Z(\delta)$

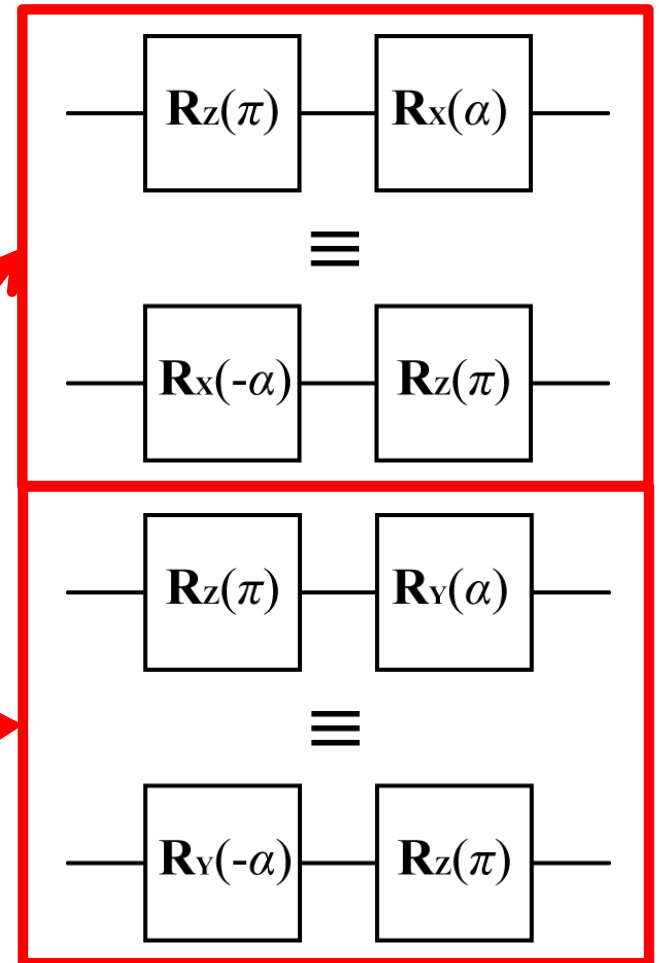
- Multiple Satisfying Values of (α, β, δ) for a Given γ Angle
- Tait-Bryan Increase the Rotation Operations that can be Reduced using Other Identities and Decompositions

Some Useful Identities for $SO(3)$ Elements

- Useful for Combining and Eliminating Elemental Rotations
- Choose Appropriate 3D Angles and Decomposition Types to Allow Advantageous use of Identities
- Use Other Identities (such as):

- $\mathbf{R}_k(\alpha)\mathbf{R}_k(-\alpha) = \mathbf{I}$
- $\mathbf{R}_x(\alpha)\mathbf{R}_z(\pi) = \mathbf{R}_z(\pi)\mathbf{R}_x(-\alpha)$
- $\mathbf{R}_y(\alpha)\mathbf{R}_z(\pi) = \mathbf{R}_z(\pi)\mathbf{R}_y(-\alpha)$
- $\mathbf{R}_k(\alpha)\mathbf{R}_k(\beta) = \mathbf{R}_k(\alpha + \beta)$
- $\mathbf{R}_k(\pm 2m\pi) = \mathbf{R}_k(0) = \mathbf{I}$

$$k \in \{\mathbf{X}, \mathbf{Y}, \mathbf{Z}\}$$

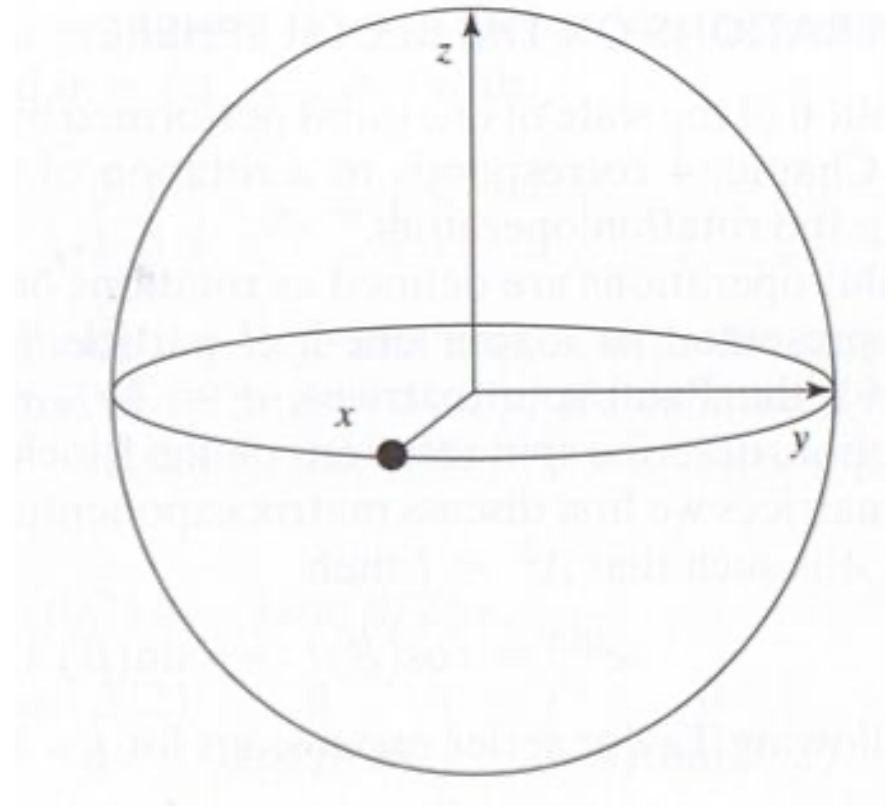
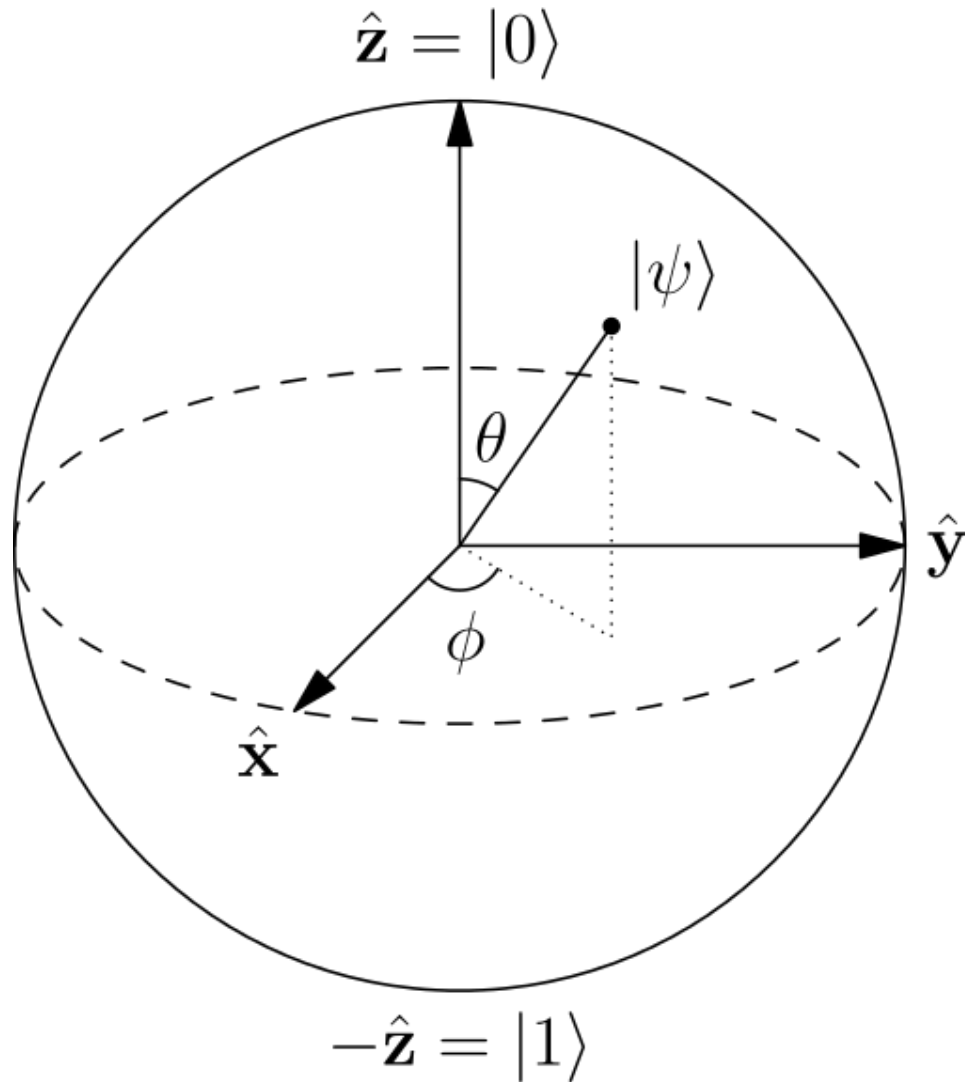


Quantum Logic Gates and Circuits

Bloch Sphere Geometry

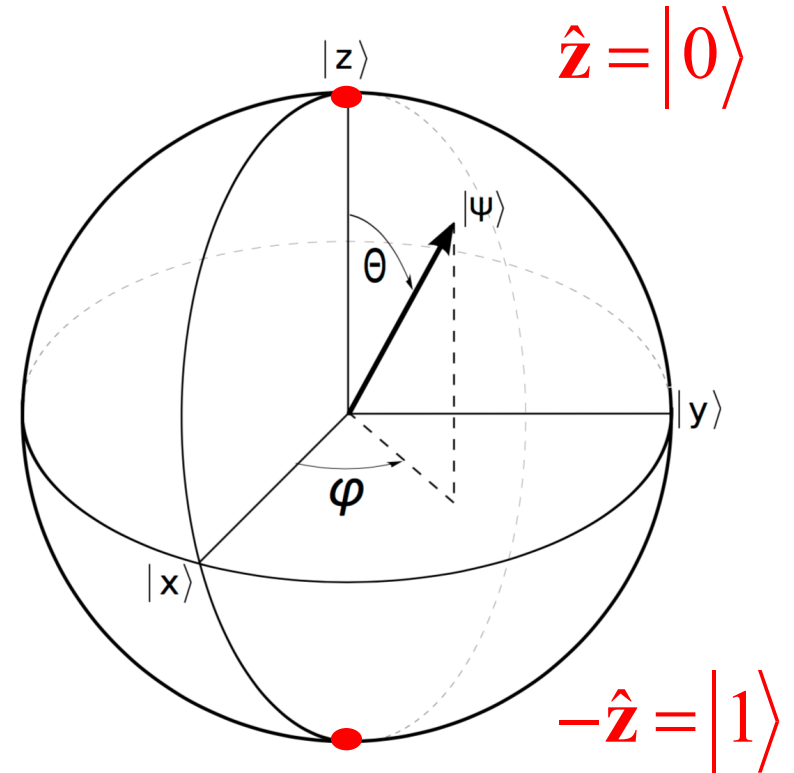
Bloch Sphere

- Geometric Representation of SINGLE qubit quantum state



Bloch Sphere

- Unit Radius Sphere
- Point on Surface Represents the State of a Qubit
- Qubit: Quantum State of Single “Information Carrier”
- The “Wave Function”
- Solution to Schrodinger Wave Equation Under Certain Assumptions
- Photonic Information Carriers include Location (spatial mode) and Polarization
- So-called “Computational Basis” Indicated in Red Font

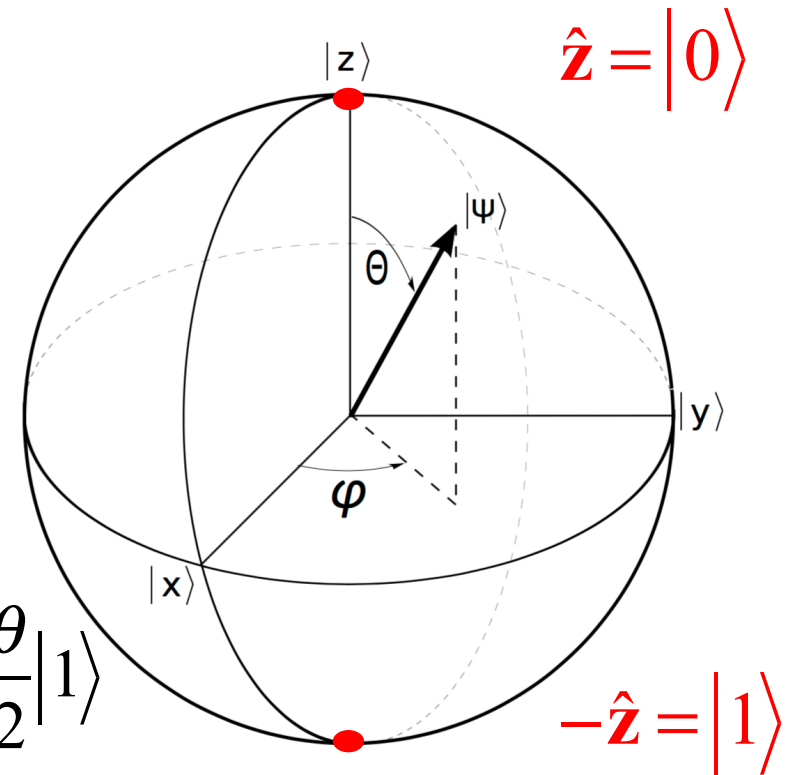


Bloch Sphere

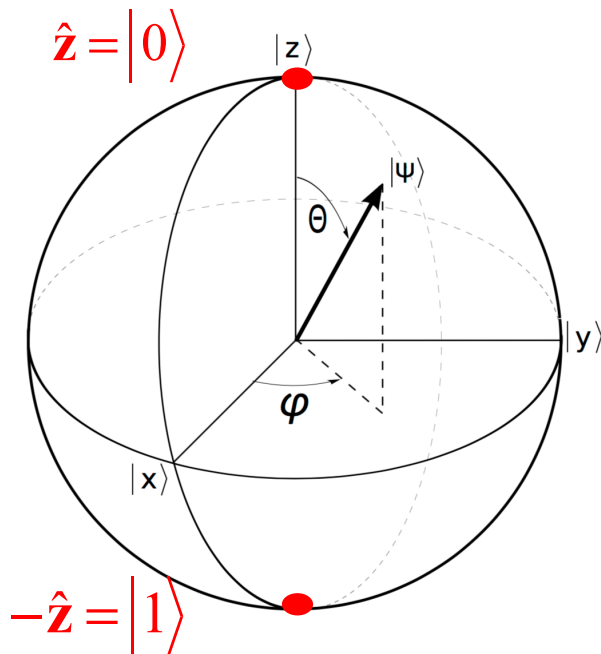
- Unit Radius Sphere
- Point on Surface Represents the State of a Qubit
- Qubit: Quantum State of Single “Information Carrier”
- The qubit:
- In General,

$$|\psi\rangle = \cos\frac{\theta}{2}|0\rangle + e^{i\phi}\sin\frac{\theta}{2}|1\rangle$$

$$|\psi\rangle = \alpha|0\rangle + \beta|1\rangle \quad |\alpha|^2 + |\beta|^2 = 1$$



Bloch Sphere as Qubit State Space



- Unit Radius Sphere
- ϕ : x - y Plane Angle Represents Phase
- θ : x - z Plane Angle Represents Superposition of Computation Basis (**Z**-Basis)

- Point on Surface of Sphere Rectangular Coordinates:

$$\hat{\mathbf{n}} = n_x \hat{\mathbf{x}} + n_y \hat{\mathbf{y}} + n_z \hat{\mathbf{z}} \quad \hat{\mathbf{n}} \cdot \hat{\mathbf{x}} = \cos \phi \quad \hat{\mathbf{n}} \cdot \hat{\mathbf{y}} = \sin \theta \quad \hat{\mathbf{n}} \cdot \hat{\mathbf{z}} = \cos \theta$$

$$\hat{\mathbf{n}} = (\sin \theta \cos \phi, \sin \theta \sin \phi, \cos \theta) \quad \text{Bloch Vector}$$

$$= \hat{\mathbf{x}} \sin \theta \cos \phi + \hat{\mathbf{y}} \sin \theta \sin \phi + \hat{\mathbf{z}} \cos \theta$$

$$= n_x \hat{\mathbf{x}} + n_y \hat{\mathbf{y}} + n_z \hat{\mathbf{z}}$$

- Point on Surface in Spherical Coordinates: $\hat{\mathbf{n}} = (1, \theta, \phi)$
- Phase ϕ “irrelevant” in Terms of Information Content

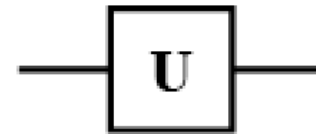
Quantum Logic Gates and Circuits

Physical Meaning of Single Qubit Gate

Moving on the Sphere Surface

Generic Single Qubit Gate

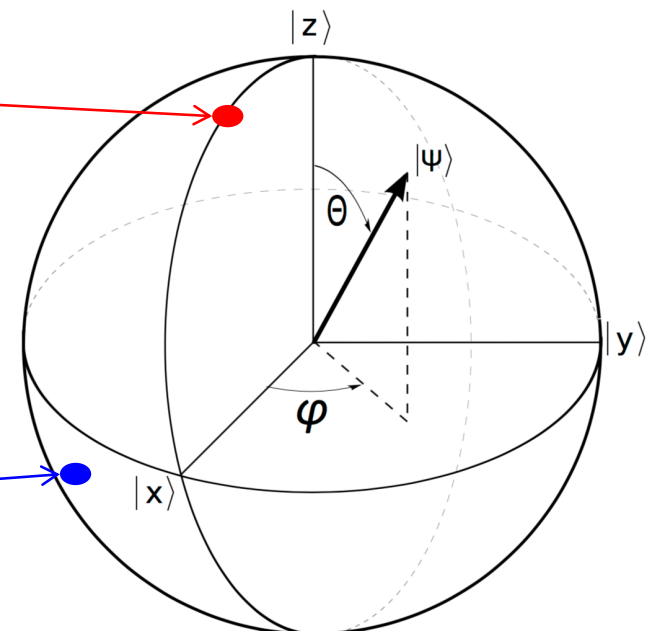
$$\mathbf{U} = \begin{bmatrix} u_{00} & u_{01} \\ u_{10} & u_{11} \end{bmatrix}$$



- Denoted as \mathbf{U} since Unitary Matrix
 - Hermitian: $\mathbf{U}^\dagger = \mathbf{U}$ (\mathbf{U}^\dagger denotes the self-adjoint of \mathbf{U})
 - Adjoint means transpose and conjugate
 - Self-Inverse: $\mathbf{U}^2 = \mathbf{I}$, $\mathbf{U}^{-1} = \mathbf{U}$

“Ket Gamma-nought Evolves to Ket Gamma-sub-one”

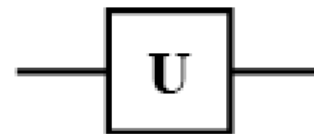
$$|\gamma_1\rangle = \mathbf{U}|\gamma_0\rangle$$



Schrodinger's Wave Equation

Generic Single Qubit Gate

$$\mathbf{U} = \begin{bmatrix} u_{00} & u_{01} \\ u_{10} & u_{11} \end{bmatrix}$$



$$\mathcal{H}|\psi(\vec{r}, t)\rangle = i\hbar \frac{\partial}{\partial t} |\psi(\vec{r}, t)\rangle$$

$$\mathcal{H}|\psi(t)\rangle = i\hbar \frac{\partial}{\partial t} |\psi(t)\rangle$$

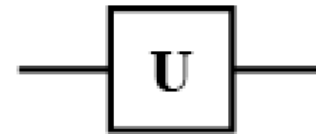
$$|\psi(t)\rangle = e^{-i\mathcal{H}t/\hbar} |\psi(0)\rangle$$

$$\mathbf{U} = e^{-i\mathcal{H}t/\hbar}$$

Relation of \mathbf{U} to Qubit Wave Function

Generic Single Qubit Gate

$$\mathbf{U} = \begin{bmatrix} u_{00} & u_{01} \\ u_{10} & u_{11} \end{bmatrix}$$



- Schrodinger's Time-varying Wave Equation,

$$\mathcal{H}|\psi(\vec{r}, t)\rangle = i\hbar \frac{\partial}{\partial t} |\psi(\vec{r}, t)\rangle \quad |\psi(t)\rangle = e^{-i\mathcal{H}t/\hbar} |\psi(0)\rangle$$
$$\mathbf{U} = e^{-i\mathcal{H}t/\hbar}$$

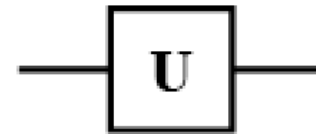
- Hamiltonian \mathcal{H} , is Self-adjoint
- \mathbf{U} is Unitary since $\mathbf{U}^\dagger \mathbf{U} = \mathbf{I}$
- Hence, the Solution or Wave Function Evolves over Finite Time Interval (Gate Delay) as: $|\psi(t)\rangle = \mathbf{U}(t)|\psi(0)\rangle$
 - *i.e.*, a product of unitary operators is finite
- For Case of Time Independent Hamiltonian, \mathcal{H} : $\mathbf{U} = e^{-i\mathcal{H}t/\hbar}$

$$|\psi(t)\rangle = \mathbf{U}|\psi(0)\rangle \quad |\psi_1\rangle = \mathbf{U}|\psi_0\rangle$$

Relation of \mathbf{U} to Qubit Wave Function (cont.)

Generic Single Qubit Gate

$$\mathbf{U} = \begin{bmatrix} u_{00} & u_{01} \\ u_{10} & u_{11} \end{bmatrix}$$



- Quantum Informatics:
 - Wave Function or Quantum State Represents Information
 - Quantum State of Single Carrier is a “Qubit”
 - Quantum State Evolution over Finite Time is a “Gate”
 - Gate is Represented by Unitary Operator, \mathbf{U}
- Theoretical Model of Quantum Gate: $|\psi_1\rangle = \mathbf{U}|\psi_0\rangle$
- Example of General Single-gate Evolution:

$$|\psi_1\rangle = \mathbf{U}|\psi_0\rangle = \begin{bmatrix} u_{00} & u_{01} \\ u_{10} & u_{11} \end{bmatrix} \begin{bmatrix} \alpha \\ \beta \end{bmatrix} = \begin{bmatrix} \alpha u_{00} + u_{01} \beta \\ \alpha u_{10} + u_{11} \beta \end{bmatrix} = (\alpha u_{00} + u_{01} \beta) |0\rangle + (\alpha u_{10} + u_{11} \beta) |1\rangle$$

Quantum Informatics: Evolutions are Computations

- Series of Operators (Gates) over Time Evolve Qubits
- For Single Qubit, Conceptually is a Path on the Surface of Bloch Sphere
- Discretized as a Series of Locations at Discrete Times
 - Initial Location is Quantum State at Time Zero
 - Each Operation (Gate) “Moves” or “Rotates” Quantum State Vector to New Position on Bloch Sphere
 - Final Position is Result of Computation
- Series of Computations are Conceptually:
 - A Program on a Quantum Computer
 - A Cascade of Quantum Logic Gates
- Graphical Representation is a Quantum Logic Circuit or a Quantum Program

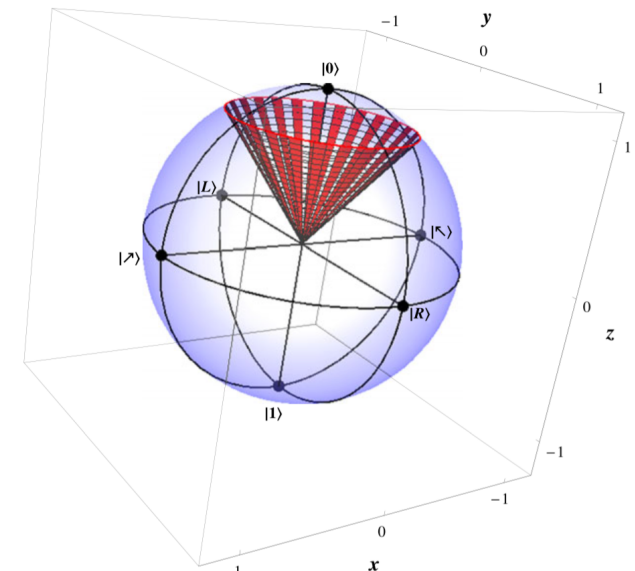
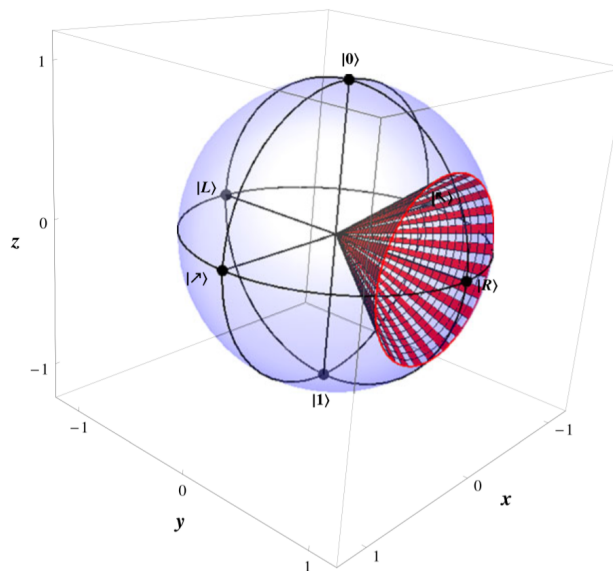
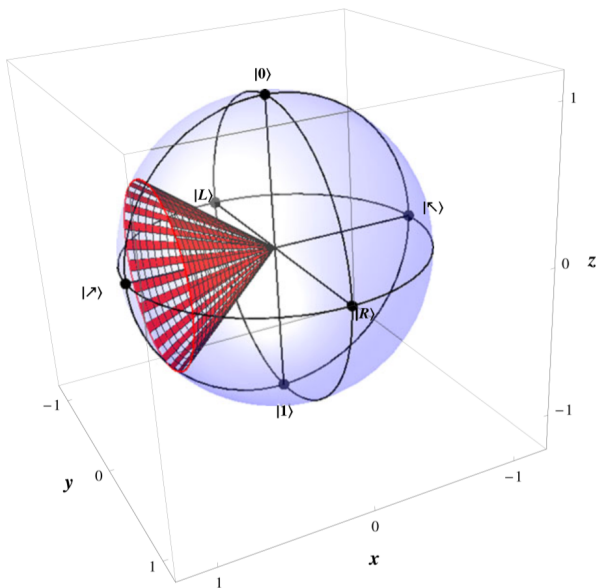
Moving on the Sphere Surface

- Moving by Rotating About an Axis
- Defines a “Cone” on Surface About Each Axis:

$$\mathbf{R}_Y(\alpha) = \begin{bmatrix} \cos(\alpha/2) & -\sin(\alpha/2) \\ \sin(\alpha/2) & \cos(\alpha/2) \end{bmatrix}$$

$$\mathbf{R}_X(\alpha) = \begin{bmatrix} \cos(\alpha/2) & -i\sin(\alpha/2) \\ -i\sin(\alpha/2) & \cos(\alpha/2) \end{bmatrix}$$

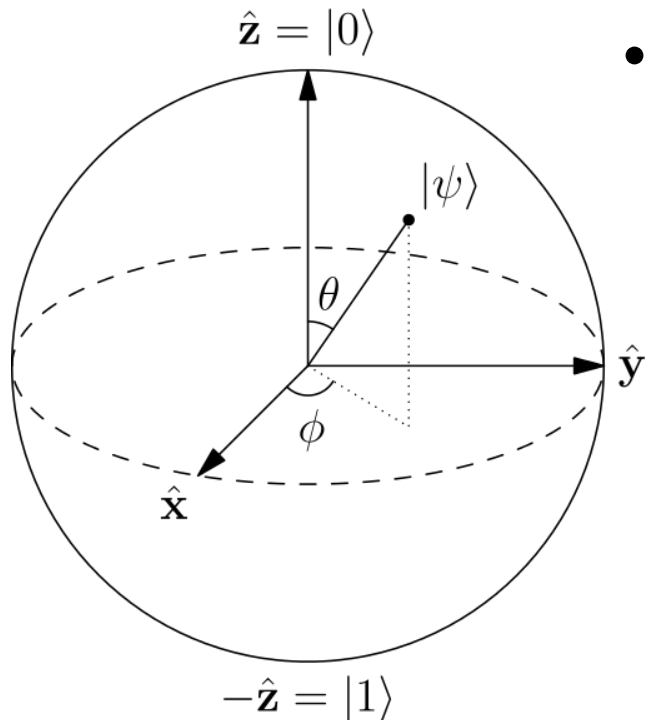
$$\mathbf{R}_Z(\alpha) = \begin{bmatrix} e^{-i(\alpha/2)} & 0 \\ 0 & e^{i(\alpha/2)} \end{bmatrix}$$



Parameterized Bloch Sphere

- Representing Qubits in terms of Bloch Sphere Angles

$$\begin{aligned}
 |\psi\rangle &= \cos\left(\frac{\theta}{2}\right)|0\rangle + e^{i\phi}\sin\left(\frac{\theta}{2}\right)|1\rangle & 0 \leq \theta \leq \pi \\
 &= \cos\left(\frac{\theta}{2}\right)|0\rangle + (\cos\phi + i\sin\phi)\sin\left(\frac{\theta}{2}\right)|1\rangle & 0 \leq \phi \leq 2\pi
 \end{aligned}$$



- Qubits can have an Arbitrary Phase Shift, γ , that is Irrelevant and is NOT Represented on Bloch Sphere

$$\begin{aligned}
 |\psi\rangle &= e^{i\gamma} \left(\cos\left(\frac{\theta}{2}\right)|0\rangle + e^{i\phi}\sin\left(\frac{\theta}{2}\right)|1\rangle \right) \\
 &\Rightarrow \cos\left(\frac{\theta}{2}\right)|0\rangle + e^{i\phi}\sin\left(\frac{\theta}{2}\right)|1\rangle
 \end{aligned}$$

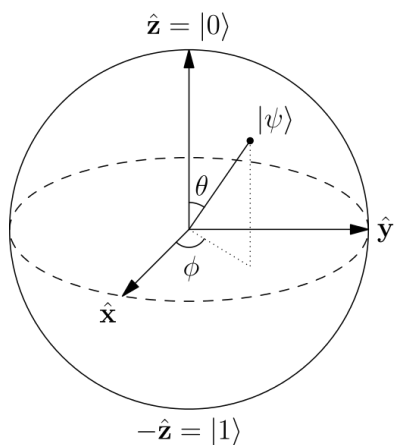
Parameterized Bloch Sphere

- Unitary Transfer Matrix that Causes Qubit to Rotate to New Position on Bloch Sphere through a Time-based Evolution can be Parameterized by Bloch Sphere Angles

$$\mathbf{R}(\theta, \phi) = |\psi\rangle\langle\psi| = \begin{bmatrix} \cos\frac{\theta}{2} & \\ & e^{i\phi} \sin\frac{\theta}{2} \end{bmatrix} \otimes \begin{bmatrix} \cos\frac{\theta}{2} & e^{-i\phi} \sin\frac{\theta}{2} \\ & \end{bmatrix}$$

$$0 \leq \theta \leq \pi$$

$$0 \leq \phi \leq 2\pi$$



$$= \begin{bmatrix} \cos^2\frac{\theta}{2} & e^{-i\phi} \cos\frac{\theta}{2} \sin\frac{\theta}{2} \\ e^{i\phi} \cos\frac{\theta}{2} \sin\frac{\theta}{2} & \sin^2\frac{\theta}{2} \end{bmatrix}$$

$$= \frac{1}{2} \begin{bmatrix} 1 + \cos\theta & \cos\phi \sin\theta - i \sin\phi \sin\theta \\ \cos\phi \sin\theta + i \sin\phi \cos\theta & 1 - \cos\theta \end{bmatrix}$$

Quantum Logic Gates and Circuits

The Pauli Gates

Parameterized Bloch Sphere

- General Rotation Matrix $\mathbf{R}(\theta, \phi)$ can be Expressed in Terms of the Pauli Matrices as:

$$\mathbf{R}(\theta, \phi) = |\psi\rangle\langle\psi| = \frac{1}{2}(\mathbf{I} + \mathbf{X}\cos\phi\sin\theta + \mathbf{Y}\sin\phi\sin\theta + \mathbf{Z}\cos\theta) = \frac{1}{2}(\mathbf{I} + \vec{\mathbf{n}} \cdot \vec{\boldsymbol{\sigma}})$$

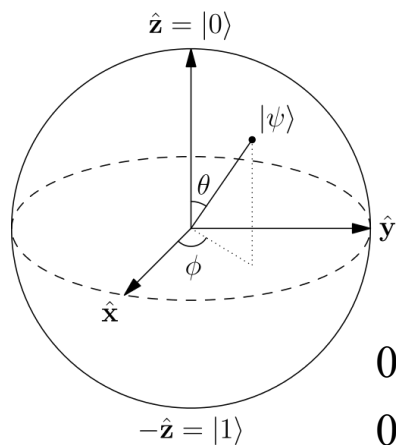
- Position of Qubit is (unit) Bloch Vector

$$\vec{\mathbf{n}} = \hat{\mathbf{n}} = (n_x, n_y, n_z) = (\cos\phi\sin\theta, \sin\phi\sin\theta, \cos\theta) = n_x\hat{\mathbf{x}} + n_y\hat{\mathbf{y}} + n_z\hat{\mathbf{z}}$$

- The “Pauli Vector” is a Vector of Matrices:

$$\vec{\boldsymbol{\sigma}} = \hat{\boldsymbol{\sigma}} = (\mathbf{X}, \mathbf{Y}, \mathbf{Z})$$

- The Pauli Matrices:



$$0 \leq \theta \leq \pi$$

$$0 \leq \phi \leq 2\pi$$

$$\mathbf{X} = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \quad \mathbf{Y} = \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix} \quad \mathbf{Z} = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$$

Parameterized Bloch Sphere

- The Bloch Vector is Rotated about the Axes in the Bloch Sphere via the Following Rotation Matrices:

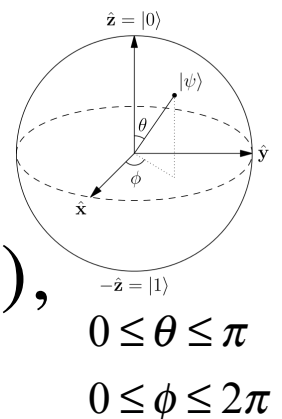
$$\mathbf{R}_x(\theta) = e^{-i\frac{\theta}{2}\mathbf{X}} \quad \mathbf{R}_y(\theta) = e^{-i\frac{\theta}{2}\mathbf{Y}} \quad \mathbf{R}_z(\theta) = e^{-i\frac{\theta}{2}\mathbf{Z}}$$

- Since the Pauli Matrices are Unitary (eg. $\mathbf{A}^2 = \mathbf{I}$), can use Euler's Equation in Matrix Form:

$$e^{\pm i\theta\mathbf{A}} = \cos(\theta)\mathbf{I} \pm i\sin(\theta)\mathbf{A}$$

- Since the Pauli Matrices Satisfy $\mathbf{X}^2 = \mathbf{Y}^2 = \mathbf{Z}^2 = \mathbf{I}$, Rotation Operators Become:

$$\mathbf{R}_x(\theta) = e^{-i\frac{\theta}{2}\mathbf{X}} = \cos\frac{\theta}{2}\mathbf{I} - i\sin\frac{\theta}{2}\mathbf{X} = \begin{bmatrix} \cos\frac{\theta}{2} & -i\sin\frac{\theta}{2} \\ -i\sin\frac{\theta}{2} & \cos\frac{\theta}{2} \end{bmatrix}$$



Parameterized Bloch Sphere

- Since the Pauli Matrices Satisfy $\mathbf{X}^2=\mathbf{Y}^2=\mathbf{Z}^2=\mathbf{I}$,
Rotation Operators Become:

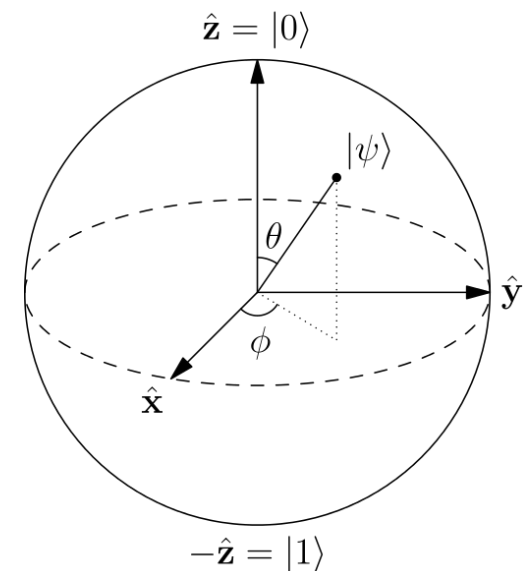
$$\mathbf{R}_x(\theta) = e^{-i\frac{\theta}{2}\mathbf{X}} = \cos\frac{\theta}{2}\mathbf{I} - i\sin\frac{\theta}{2}\mathbf{X} = \begin{bmatrix} \cos\frac{\theta}{2} & -i\sin\frac{\theta}{2} \\ -i\sin\frac{\theta}{2} & \cos\frac{\theta}{2} \end{bmatrix}$$

$$\mathbf{R}_y(\theta) = e^{-i\frac{\theta}{2}\mathbf{Y}} = \cos\frac{\theta}{2}\mathbf{I} - i\sin\frac{\theta}{2}\mathbf{Y} = \begin{bmatrix} \cos\frac{\theta}{2} & \sin\frac{\theta}{2} \\ \sin\frac{\theta}{2} & \cos\frac{\theta}{2} \end{bmatrix}$$

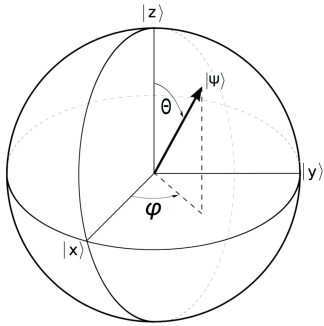
$$\mathbf{R}_z(\theta) = e^{-i\frac{\theta}{2}\mathbf{Z}} = \cos\frac{\theta}{2}\mathbf{I} - i\sin\frac{\theta}{2}\mathbf{Z} = \begin{bmatrix} e^{-i\frac{\theta}{2}} & 0 \\ 0 & e^{i\frac{\theta}{2}} \end{bmatrix}$$

$$0 \leq \theta \leq \pi$$

$$0 \leq \phi \leq 2\pi$$



The Pauli Gates: Rotation Angle is π



- The Trivial Case:

$$\sigma_0 = \mathbf{I} = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \quad \begin{array}{l} |0\rangle \rightarrow |0\rangle \\ |1\rangle \rightarrow |1\rangle \end{array}$$

————

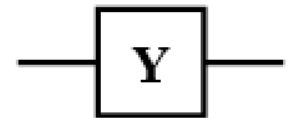
- Pauli-**X** Rotates About **X** by π Radians:

$$\sigma_1 = \sigma_X = \mathbf{X} = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \quad \begin{array}{l} |0\rangle \rightarrow |1\rangle \\ |1\rangle \rightarrow |0\rangle \end{array} \quad \begin{array}{l} \text{"NOT"} \\ \text{"bit flip"} \end{array}$$



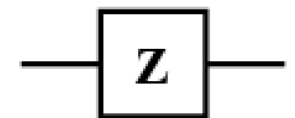
- Pauli-**Y** Rotates About **Y** by π Radians:

$$\sigma_2 = \sigma_Y = \mathbf{Y} = \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix} \quad \begin{array}{l} |0\rangle \rightarrow i|1\rangle \\ |1\rangle \rightarrow -i|0\rangle \end{array}$$



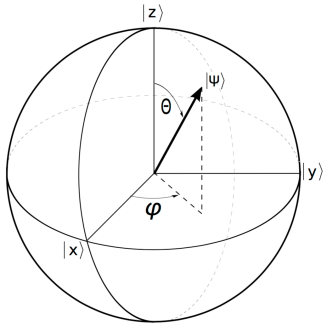
- Pauli-**Z** Rotates About **Z** by π Radians:

$$\sigma_3 = \sigma_Z = \mathbf{Z} = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \quad \begin{array}{l} |0\rangle \rightarrow |0\rangle \\ |1\rangle \rightarrow -|1\rangle \end{array} \quad \text{"phase flip"}$$



$$\mathbf{X}^2 = \mathbf{Y}^2 = \mathbf{Z}^2 = -i\mathbf{XYZ} = \mathbf{I}$$

Parameterized Rotation



- Arbitrary Rotation About Axis j :

$$\mathbf{J}_\alpha = \cos \alpha \mathbf{I} + i \sin \alpha \mathbf{J}$$

- j Axis is x , y , or z Axis
- $\mathbf{J} = \mathbf{X}$, \mathbf{Y} , or \mathbf{Z}
- Theoretical “Gate” or Operator
- Bloch Sphere is Unit Radius
- Point on Surface of Sphere Rectangular Coordinates:

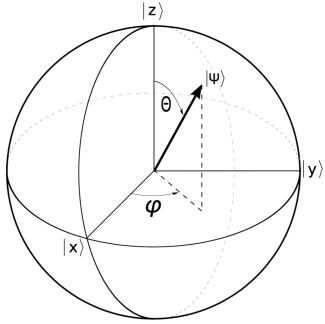
$$\hat{\mathbf{n}} = n_x \hat{\mathbf{x}} + n_y \hat{\mathbf{y}} + n_z \hat{\mathbf{z}}$$

$$\hat{\mathbf{n}} \cdot \hat{\mathbf{x}} = \cos \varphi$$

$$\hat{\mathbf{n}} \cdot \hat{\mathbf{y}} = \sin \theta$$

$$\hat{\mathbf{n}} \cdot \hat{\mathbf{z}} = \cos \theta$$

Moving on the Sphere Surface



Using Pauli Gates for General Rotations

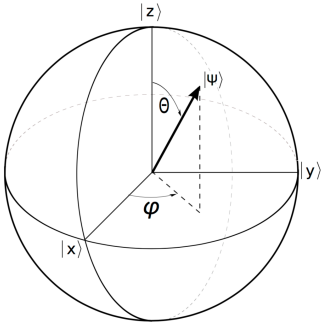
- We can Move Anywhere by using Pauli Rotations
 - And one more operator

$$\hat{\mathbf{n}} = n_x \hat{x} + n_y \hat{y} + n_z \hat{z} \quad \text{Bloch Vector}$$

$$\vec{\sigma} = (\sigma_x, \sigma_y, \sigma_z) = \mathbf{X}\hat{x} + \mathbf{Y}\hat{y} + \mathbf{Z}\hat{z} \quad \text{Pauli Vector}$$

$$S(\hat{\mathbf{n}}) = \hat{\mathbf{n}} \cdot \vec{\sigma} = n_x \sigma_x + n_y \sigma_y + n_z \sigma_z$$

The Pauli Gate Algebra



$$\mathbf{H} = \frac{\mathbf{X} + \mathbf{Z}}{\sqrt{2}}$$

$$\mathbf{XY} = -\mathbf{YZ} = i\mathbf{Z}$$

$$\mathbf{YZ} = -\mathbf{ZY} = i\mathbf{X}$$

$$\mathbf{ZX} = -\mathbf{XZ} = i\mathbf{Y}$$

$$\mathbf{HXH} = \mathbf{Z}$$

$$\mathbf{HYH} = -\mathbf{Y}$$

$$\mathbf{HZH} = \mathbf{X}$$

$$\mathbf{X}^2 = \mathbf{Y}^2 = \mathbf{Z}^2 = \mathbf{H}^2 = \mathbf{I}$$

$$\mathbf{XXX} = \mathbf{X}$$

$$\mathbf{XYX} = -\mathbf{Y}$$

$$\mathbf{XZX} = -\mathbf{Z}$$

$$\mathbf{X}^2 = \mathbf{Y}^2 = \mathbf{Z}^2 = -i\mathbf{XYZ} = \mathbf{I}$$

$$\mathbf{I} = -i\mathbf{XYZ}$$

$$(\mathbf{I})\mathbf{Z} = (-i\mathbf{XYZ})\mathbf{Z}$$

$$\mathbf{Z} = -i\mathbf{XY}$$

$$(\mathbf{Z})\mathbf{Y} = (-i\mathbf{XY})\mathbf{Y}$$

$$\mathbf{ZY} = -i\mathbf{X}$$

$$(\mathbf{ZY})\mathbf{X} = (-i\mathbf{X})\mathbf{X}$$

$$\mathbf{Z Y X} = -i\mathbf{I}$$

$$\mathbf{Z}(\mathbf{Z Y X}) = \mathbf{Z}(-i\mathbf{I})$$

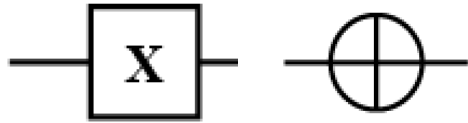
$$\mathbf{Y X} = -i\mathbf{Z}$$

Pauli Operators (Theory)

- Pauli Group \mathbb{P} is Set of Pauli Operators with Coefficients $\{\pm 1, \pm i\}$
- Single Qubit Pauli Group \mathbb{P}_1 is:
$$\mathbb{P}_1 = \{\mathbf{I}, \pm\mathbf{X}, \pm\mathbf{Y}, \pm\mathbf{Z}, \pm i\mathbf{I}, \pm i\mathbf{X}, \pm i\mathbf{Y}, \pm i\mathbf{Z}\} = \{\mathbf{P} \in \mathbb{P}_1\}$$
- Multi-qubit (n) Pauli Group Consists of Elements that are Products of n Pauli Operators
- Clifford Group, $\mathbf{C} = \{\mathbf{C}_i \in \mathbf{C}\}$, is Group of Transformations that Leave the Pauli Group Invariant $\mathbf{C}\mathbf{P}\mathbf{C}^\dagger = \mathbf{P}'$ where $\mathbf{P}' \in \mathbb{P}_1$ with
- Prominent Members of Clifford Group are: Hadamard (\mathbf{H}), Phase (\mathbf{S}), and Controlled-NOT (\mathbf{CNOT} , Feynman, Controlled- \mathbf{X})
- Gates in \mathbf{P} and \mathbf{C} are not Universal Unless \mathbf{T} Included

Euler Elemental Rotation Decomposition Examples for Pauli Operators

- Global Phase Shifts of δ are Represented as $(e^{i\delta})\mathbf{I}$



$$\mathbf{X} = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$$

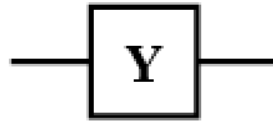
$$\mathbf{X} = \mathbf{R}_X(\pi) \left(e^{i\frac{\pi}{2}} \right) \mathbf{I}$$

$$\mathbf{X} = \mathbf{R}_Z(\pi) \mathbf{R}_Y(\pi) \left(e^{-i\frac{\pi}{2}} \right) \mathbf{I}$$

$$\mathbf{X} = \mathbf{R}_Y(\pi) \mathbf{R}_Z(\pi) \left(e^{i\frac{\pi}{2}} \right) \mathbf{I}$$

$$\mathbf{X} = \mathbf{R}_Z\left(-\frac{\pi}{2}\right) \mathbf{R}_Y(\pi) \mathbf{R}_Z\left(\frac{\pi}{2}\right) \left(e^{i\frac{\pi}{2}} \right) \mathbf{I}$$

$$\mathbf{X} = \mathbf{R}_Y\left(\frac{\pi}{2}\right) \mathbf{R}_Z(\pi) \mathbf{R}_Y\left(-\frac{\pi}{2}\right) \left(e^{i\frac{\pi}{2}} \right) \mathbf{I}$$



$$\mathbf{Y} = \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix}$$

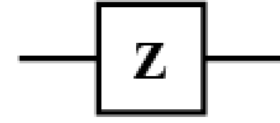
$$\mathbf{Y} = \mathbf{R}_Y(\pi) \left(e^{i\frac{\pi}{2}} \right) \mathbf{I}$$

$$\mathbf{Y} = \mathbf{R}_Z(\pi) \mathbf{R}_X(\pi) \left(e^{-i\frac{\pi}{2}} \right) \mathbf{I}$$

$$\mathbf{Y} = \mathbf{R}_X(\pi) \mathbf{R}_Z(\pi) \left(e^{i\frac{\pi}{2}} \right) \mathbf{I}$$

$$\mathbf{Y} = \mathbf{R}_Z\left(\frac{\pi}{2}\right) \mathbf{R}_X(\pi) \mathbf{R}_Z\left(-\frac{\pi}{2}\right) \left(e^{i\frac{\pi}{2}} \right) \mathbf{I}$$

$$\mathbf{Y} = \mathbf{R}_X\left(-\frac{\pi}{2}\right) \mathbf{R}_Z(\pi) \mathbf{R}_X\left(\frac{\pi}{2}\right) \left(e^{i\frac{\pi}{2}} \right) \mathbf{I}$$



$$\mathbf{Z} = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$$

$$\mathbf{Z} = \mathbf{R}_Z(\pi) \left(e^{i\frac{\pi}{2}} \right) \mathbf{I}$$

$$\mathbf{Z} = \mathbf{R}_X(\pi) \mathbf{R}_Y(\pi) \left(e^{i\frac{\pi}{2}} \right) \mathbf{I}$$

$$\mathbf{Z} = \mathbf{R}_Y(\pi) \mathbf{R}_X(\pi) \left(e^{-i\frac{\pi}{2}} \right) \mathbf{I}$$

$$\mathbf{Z} = \mathbf{R}_X\left(\frac{\pi}{2}\right) \mathbf{R}_Y(\pi) \mathbf{R}_X\left(-\frac{\pi}{2}\right) \left(e^{i\frac{\pi}{2}} \right) \mathbf{I}$$

$$\mathbf{Z} = \mathbf{R}_Y\left(-\frac{\pi}{2}\right) \mathbf{R}_X(\pi) \mathbf{R}_Y\left(\frac{\pi}{2}\right) \left(e^{i\frac{\pi}{2}} \right) \mathbf{I}$$

Quantum Logic Gates and Circuits

Single Qubit Gates that Modify the Probability Amplitude or Produce Pure (Relative) Phase Shifts

Hadamard Operator

- This Operator is Commonly used to Maximize Superposition of a Qubit in a Basis State

$$\mathbf{H} = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix}$$

- Example: $|\psi\rangle = \alpha_0 |0\rangle + \alpha_1 |1\rangle$

$$\mathbf{H} |\psi\rangle = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} \begin{bmatrix} \alpha_0 \\ \alpha_1 \end{bmatrix} = \frac{\alpha_0}{\sqrt{2}} (|0\rangle + |1\rangle) + \frac{\alpha_1}{\sqrt{2}} (|0\rangle - |1\rangle)$$

Hadamard Operator

- This Operator is Commonly used to Maximize Superposition of a Qubit in a Basis State

$$\mathbf{H} = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix}$$

- Example: $|\psi\rangle = 0|0\rangle + 1|1\rangle = |1\rangle$

$$\mathbf{H}|\psi\rangle = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} \begin{bmatrix} 0 \\ 1 \end{bmatrix} = \begin{bmatrix} 1/\sqrt{2} \\ -1/\sqrt{2} \end{bmatrix} = (1/\sqrt{2})|0\rangle - (1/\sqrt{2})|1\rangle$$

$$\text{Prob}[|0\rangle \text{ measured}] = (1/\sqrt{2})^2 = 50\%$$

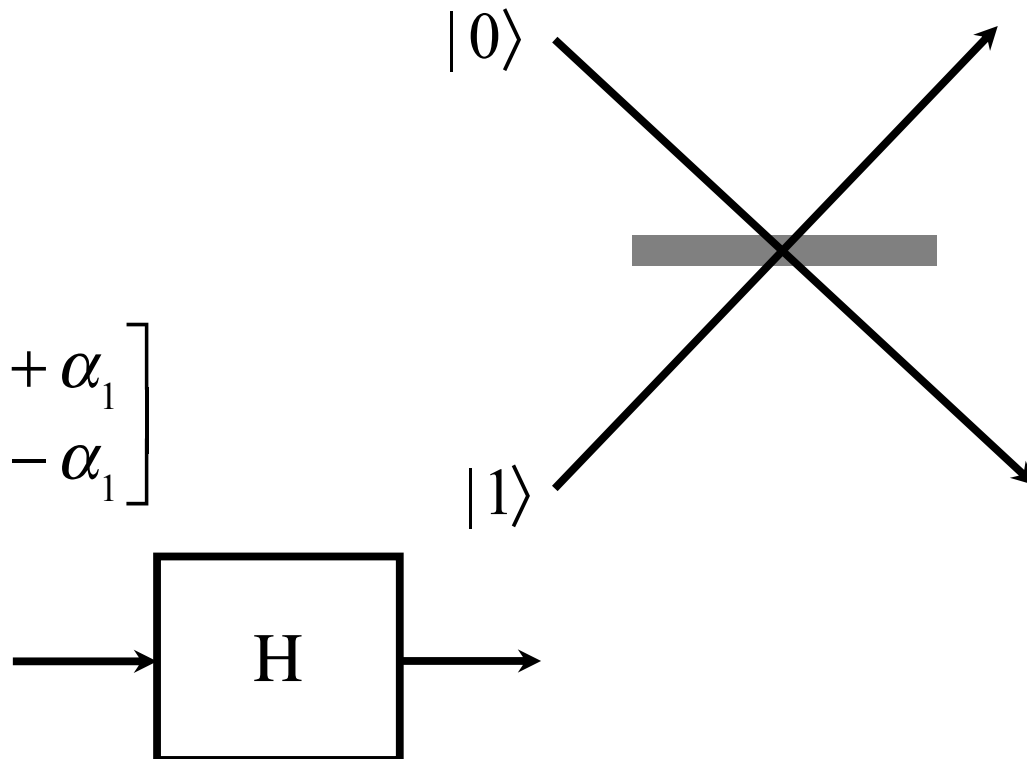
$$\text{Prob}[|1\rangle \text{ measured}] = (1/\sqrt{2})^2 = 50\%$$

Beam Splitter Example

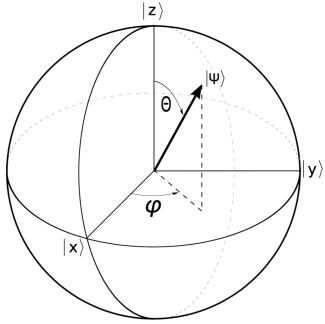
- 50-50 Beam Splitter Performs a Hadamard Transform on Particles (location/spatially encoded information)
- Beam Splitters have been Constructed for Quantum Particles other than Photons

$$|\psi\rangle = \alpha_0 |0\rangle + \alpha_1 |1\rangle$$

$$|\varphi\rangle = \mathbf{H} |\psi\rangle = \frac{1}{\sqrt{2}} \begin{bmatrix} \alpha_0 + \alpha_1 \\ \alpha_0 - \alpha_1 \end{bmatrix}$$



Other Single-Qubit Gates



- Hadamard (90-degree rotation about axis parallel to the xy -plane)

$$\mathbf{H} = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} |0\rangle \rightarrow \frac{|0\rangle + |1\rangle}{\sqrt{2}} \quad |1\rangle \rightarrow \frac{|0\rangle - |1\rangle}{\sqrt{2}} \quad \text{---} \boxed{\mathbf{H}} \text{---}$$

$$\mathbf{H} = \mathbf{R}_X(\pi) \mathbf{R}_Y\left(\frac{\pi}{2}\right) \mathbf{Ph}\left(\frac{\pi}{2}\right) \quad \mathbf{H} = \mathbf{R}_Y\left(\frac{\pi}{2}\right) \mathbf{R}_Z(\pi) \mathbf{Ph}\left(\frac{\pi}{2}\right)$$

- Square Root of NOT (Square root of \mathbf{X}):

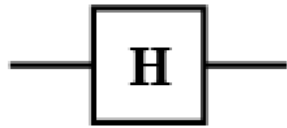
$$\mathbf{V} = \frac{1}{2} \begin{bmatrix} 1+i & 1-i \\ 1-i & 1+i \end{bmatrix} |0\rangle \rightarrow \frac{(1+i)|0\rangle + (1-i)|1\rangle}{2} \quad |1\rangle \rightarrow \frac{(1-i)|0\rangle + (1+i)|1\rangle}{2} \quad \text{---} \boxed{\mathbf{V}} \text{---}$$

$$\mathbf{V} = \mathbf{R}_X\left(\frac{\pi}{2}\right) \mathbf{Ph}\left(\frac{\pi}{4}\right) \quad \mathbf{V} = \mathbf{R}_Z\left(-\frac{\pi}{2}\right) \mathbf{R}_Y\left(\frac{\pi}{2}\right) \mathbf{R}_Z\left(-\frac{\pi}{2}\right) \mathbf{Ph}\left(\frac{\pi}{4}\right)$$

Euler Elemental Rotation Decomposition

Examples for Hadamard Operator

- Global Phase Shifts of δ are Represented as $(e^{i\delta}\mathbf{I})$



$$\mathbf{H} = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix}$$

$(x-z-x)$

$$\mathbf{H} = \mathbf{R}_x\left(\frac{\pi}{2}\right)\mathbf{R}_z\left(\frac{\pi}{2}\right)\mathbf{R}_x\left(\frac{\pi}{2}\right)\left(e^{i\frac{\pi}{2}}\mathbf{I}\right)$$

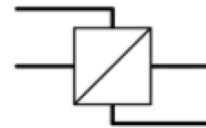
$$\mathbf{H} = \mathbf{R}_x\left(-\frac{\pi}{2}\right)\mathbf{R}_z\left(-\frac{\pi}{2}\right)\mathbf{R}_x\left(-\frac{\pi}{2}\right)\left(e^{-i\frac{\pi}{2}}\mathbf{I}\right)$$

$(y-z-y)$

$$\mathbf{H} = \mathbf{R}_y(\pi)\mathbf{R}_z(\pi)\mathbf{R}_y\left(\frac{\pi}{2}\right)\left(e^{i\frac{\pi}{2}}\mathbf{I}\right)$$

$$\mathbf{H} = \mathbf{R}_z(\pi)\mathbf{R}_y\left(-\frac{\pi}{2}\right)\left(e^{i\frac{\pi}{2}}\mathbf{I}\right)$$

$$\mathbf{H} = \mathbf{R}_y\left(\frac{\pi}{2}\right)\mathbf{R}_z(\pi)\left(e^{i\frac{\pi}{2}}\mathbf{I}\right)$$



$(y-x-y)$

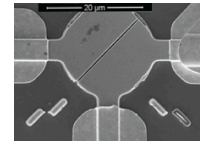
$$\mathbf{H} = \mathbf{R}_x(\pi)\mathbf{R}_y\left(\frac{\pi}{2}\right)\left(e^{i\frac{\pi}{2}}\mathbf{I}\right)$$

$$\mathbf{H} = \mathbf{R}_y(\pi)\mathbf{R}_x(\pi)\mathbf{R}_y\left(-\frac{\pi}{2}\right)\left(e^{-i\frac{\pi}{2}}\mathbf{I}\right)$$

$(z-x-z)$

$$\mathbf{H} = \mathbf{R}_z\left(\frac{\pi}{2}\right)\mathbf{R}_x\left(\frac{\pi}{2}\right)\mathbf{R}_z\left(\frac{\pi}{2}\right)\left(e^{i\frac{\pi}{2}}\mathbf{I}\right)$$

$$\mathbf{H} = \mathbf{R}_z\left(-\frac{\pi}{2}\right)\mathbf{R}_x\left(-\frac{\pi}{2}\right)\mathbf{R}_z\left(-\frac{\pi}{2}\right)\left(e^{-i\frac{\pi}{2}}\mathbf{I}\right)$$



SMU Photonic Quantum Circuit

Single Qubit Operations (Square Root of X)

Transformation Matrix:

$$\frac{1}{\sqrt{2}} \begin{bmatrix} e^{i\pi/4} & e^{-i\pi/4} \\ e^{-i\pi/4} & e^{i\pi/4} \end{bmatrix}$$

From Euler's Identity:

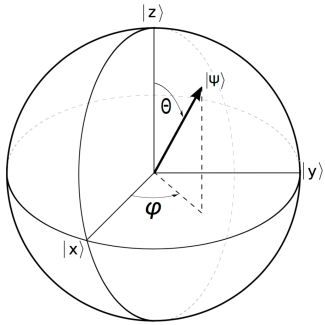
$$e^{i\pi/4} = \frac{1}{\sqrt{2}}(1+i) \qquad e^{-i\pi/4} = \frac{1}{\sqrt{2}}(1-i)$$

$$\frac{1}{\sqrt{2}} \begin{bmatrix} e^{i\pi/4} & e^{-i\pi/4} \\ e^{-i\pi/4} & e^{i\pi/4} \end{bmatrix} = \frac{1}{2} \begin{bmatrix} 1+i & 1-i \\ 1-i & 1+i \end{bmatrix}$$

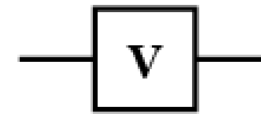
Two Gates in Series (Square of Matrix):

$$\left(\frac{1}{2}\right)^2 \begin{bmatrix} 1+i & 1-i \\ 1-i & 1+i \end{bmatrix}^2 = \frac{1}{4} \begin{bmatrix} 0 & 4 \\ 4 & 0 \end{bmatrix} = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$$


Other Single-Qubit Gates



$$\mathbf{V} = \frac{1}{2} \begin{bmatrix} 1+i & 1-i \\ 1-i & 1+i \end{bmatrix}$$




- **Square Root of NOT (Square root of X):**

$$\mathbf{V} = \frac{1}{2} \begin{bmatrix} 1+i & 1-i \\ 1-i & 1+i \end{bmatrix} \quad |0\rangle \rightarrow \frac{(1+i)|0\rangle + (1-i)|1\rangle}{2} \quad |1\rangle \rightarrow \frac{(1-i)|0\rangle + (1+i)|1\rangle}{2}$$


$$\mathbf{V} = \mathbf{R}_X\left(\frac{\pi}{2}\right) \mathbf{Ph}\left(\frac{\pi}{4}\right) \quad \mathbf{V} = \mathbf{R}_Z\left(-\frac{\pi}{2}\right) \mathbf{R}_Y\left(\frac{\pi}{2}\right) \mathbf{R}_Z\left(-\frac{\pi}{2}\right) \mathbf{Ph}\left(\frac{\pi}{4}\right)$$

- **Another Square Root of NOT Gate:**

$$\mathbf{V}^\dagger = \frac{1}{2} \begin{bmatrix} 1-i & 1+i \\ 1+i & 1-i \end{bmatrix} \quad |0\rangle \rightarrow \frac{(1-i)|0\rangle + (1+i)|1\rangle}{2} \quad |1\rangle \rightarrow \frac{(1+i)|0\rangle + (1-i)|1\rangle}{2}$$


Euler Elemental Rotation Decomposition

Examples for Square-root of NOT Gates

- Global Phase Shifts of δ are Represented as $(e^{i\delta})\mathbf{I}$



$$V = \frac{1}{2} \begin{bmatrix} 1+i & 1-i \\ 1-i & 1+i \end{bmatrix}$$

$$V^\dagger = \frac{1}{2} \begin{bmatrix} 1-i & 1+i \\ 1+i & 1-i \end{bmatrix}$$

$$V = \mathbf{R}_X\left(\frac{\pi}{2}\right) \left(e^{i\frac{\pi}{4}} \mathbf{I} \right)$$

$$V^\dagger = \mathbf{R}_X\left(-\frac{\pi}{2}\right) \left(e^{-i\frac{\pi}{4}} \mathbf{I} \right)$$

$$V = \mathbf{R}_Z\left(-\frac{\pi}{2}\right) \mathbf{R}_Y\left(\frac{\pi}{2}\right) \mathbf{R}_Z\left(\frac{\pi}{2}\right) \left(e^{i\frac{\pi}{4}} \mathbf{I} \right)$$

$$V^\dagger = \mathbf{R}_Z\left(-\frac{\pi}{2}\right) \mathbf{R}_Y\left(-\frac{\pi}{2}\right) \mathbf{R}_Z\left(\frac{\pi}{2}\right) \left(e^{-i\frac{\pi}{4}} \mathbf{I} \right)$$

$$V = \mathbf{R}_Y\left(\frac{\pi}{2}\right) \mathbf{R}_Z\left(\frac{\pi}{2}\right) \mathbf{R}_Y\left(-\frac{\pi}{2}\right) \left(e^{i\frac{\pi}{4}} \mathbf{I} \right)$$

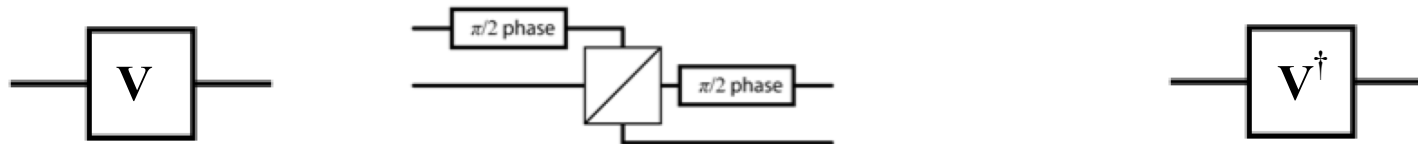
$$V^\dagger = \mathbf{R}_Y\left(\frac{\pi}{2}\right) \mathbf{R}_Z\left(-\frac{\pi}{2}\right) \mathbf{R}_Y\left(-\frac{\pi}{2}\right) \left(e^{-i\frac{\pi}{4}} \mathbf{I} \right)$$

$$V = \mathbf{R}_Z\left(-\frac{\pi}{2}\right) \mathbf{H} \mathbf{R}_Z\left(-\frac{\pi}{2}\right) \left(e^{-i\frac{\pi}{4}} \mathbf{I} \right)$$

(Example of V in Terms
of H Operator)

Decomposition Examples for Square-root of NOT Gates with Hadamard

- Global Phase Shifts of δ are Represented as $(e^{i\delta})\mathbf{I}$



$$\mathbf{V} = \frac{1}{2} \begin{bmatrix} 1+i & 1-i \\ 1-i & 1+i \end{bmatrix}$$

$$\mathbf{V}^\dagger = \frac{1}{2} \begin{bmatrix} 1-i & 1+i \\ 1+i & 1-i \end{bmatrix}$$

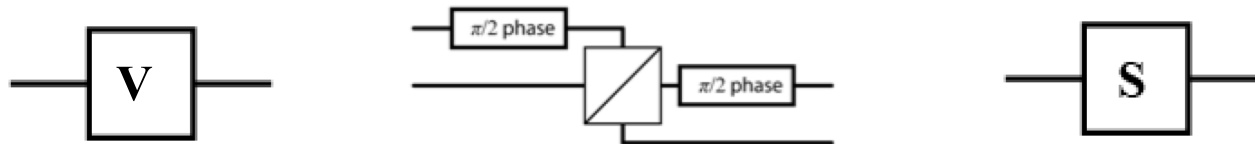
$$\mathbf{V} = \mathbf{R}_z\left(-\frac{\pi}{2}\right) \mathbf{H} \mathbf{R}_z\left(-\frac{\pi}{2}\right) \left(e^{-i\frac{\pi}{4}} \mathbf{I} \right)$$

$$\mathbf{V} = \mathbf{H} \mathbf{S} \mathbf{H} = \frac{1}{2} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & i \end{bmatrix} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} = \frac{1}{2} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} \begin{bmatrix} 1 & 1 \\ i & -i \end{bmatrix} = \frac{1}{2} \begin{bmatrix} 1+i & 1-i \\ 1-i & 1+i \end{bmatrix}$$

$$\mathbf{V}^\dagger = \mathbf{H} \mathbf{S}^\dagger \mathbf{H} = \frac{1}{2} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & -i \end{bmatrix} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} = \frac{1}{2} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} \begin{bmatrix} 1 & 1 \\ -i & i \end{bmatrix} = \frac{1}{2} \begin{bmatrix} 1-i & 1+i \\ 1+i & 1-i \end{bmatrix}$$

Decomposition Examples for Square-root of NOT Gates with S

- Global Phase Shifts of δ are Represented as $(e^{i\delta})\mathbf{I}$



$$\mathbf{V} = \frac{1}{2} \begin{bmatrix} 1+i & 1-i \\ 1-i & 1+i \end{bmatrix}$$

$$\mathbf{H} = \left(e^{-i\frac{\pi}{4}} \right) \mathbf{S} \mathbf{V} \mathbf{S}$$

$$\mathbf{S} = \begin{bmatrix} 1 & 0 \\ 0 & i \end{bmatrix}$$

$$\begin{aligned} \mathbf{S} \mathbf{V} \mathbf{S} &= \begin{bmatrix} 1 & 0 \\ 0 & i \end{bmatrix} \frac{1}{2} \begin{bmatrix} 1+i & 1-i \\ 1-i & 1+i \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & i \end{bmatrix} = \frac{1}{2} \begin{bmatrix} 1 & 0 \\ 0 & i \end{bmatrix} \begin{bmatrix} 1+i & i(1-i) \\ 1-i & i(1+i) \end{bmatrix} \\ &= \frac{1}{2} \begin{bmatrix} 1+i & i(1-i) \\ i(1-i) & -(1+i) \end{bmatrix} = \frac{1}{2} \begin{bmatrix} i(1-i) & i(1-i) \\ i(1-i) & -i(1-i) \end{bmatrix} = \frac{1}{2} \begin{bmatrix} 1+i & 1+i \\ 1+i & -(1+i) \end{bmatrix} \\ &= \left(\frac{1}{2} + i\frac{1}{2} \right) \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} = \left(e^{-i\frac{\pi}{4}} \right) \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} = \left(e^{-i\frac{\pi}{4}} \right) \mathbf{H} \end{aligned}$$

Decomposition Examples for Square-root of NOT Gates with X

- Global Phase Shifts of δ are Represented as $(e^{i\delta})\mathbf{I}$



$$\mathbf{V} = \frac{1}{2} \begin{bmatrix} 1+i & 1-i \\ 1-i & 1+i \end{bmatrix}$$

$$\mathbf{X} = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$$

$$\mathbf{VX} = \mathbf{V}^\dagger$$

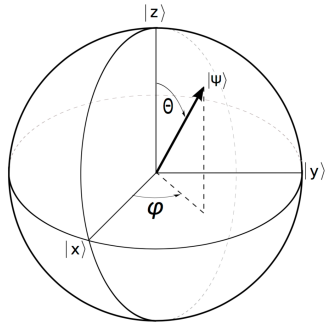
$$\mathbf{V}^\dagger \mathbf{X} = \mathbf{V}$$

$$\mathbf{VX} = \left(\frac{1}{2} \right) \begin{bmatrix} 1+i & 1-i \\ 1-i & 1+i \end{bmatrix} \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} = \left(\frac{1}{2} \right) \begin{bmatrix} 1-i & 1+i \\ 1+i & 1-i \end{bmatrix} = \mathbf{V}^\dagger$$

$$\mathbf{V}^\dagger \mathbf{X} = \left(\frac{1}{2} \right) \begin{bmatrix} 1-i & 1+i \\ 1+i & 1-i \end{bmatrix} \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} = \left(\frac{1}{2} \right) \begin{bmatrix} 1+i & 1-i \\ 1-i & 1+i \end{bmatrix} = \mathbf{V}$$

Single-Qubit Phase Shift Gates

Fixed Rotations About z-axis



- General Case (no effect on state value):

$$\mathbf{P}_\phi = \begin{bmatrix} 1 & 0 \\ 0 & e^{i\phi} \end{bmatrix} \quad |0\rangle \rightarrow |0\rangle \quad |1\rangle \rightarrow e^{i\phi}|1\rangle \quad \text{---} \boxed{\phi} \text{---}$$

- Phase Shift by $\pi/2$ (Phase Gate):

$$\mathbf{S} = \begin{bmatrix} 1 & 0 \\ 0 & i \end{bmatrix} = e^{i(\pi/4)} \begin{bmatrix} e^{-i(\pi/4)} & 0 \\ 0 & e^{i(\pi/4)} \end{bmatrix} \quad \mathbf{S}^2 = \mathbf{Z} \quad \text{---} \boxed{\mathbf{S}} \text{---}$$

- Phase Shift by $\pi/4$ (“ $\pi/8$ Gate” or “ 22.5° gate”):

$$\mathbf{T} = \begin{bmatrix} 1 & 0 \\ 0 & e^{i(\pi/4)} \end{bmatrix} = e^{i(\pi/8)} \begin{bmatrix} e^{-i(\pi/8)} & 0 \\ 0 & e^{i(\pi/8)} \end{bmatrix} \quad \mathbf{T}^2 = \mathbf{S} \quad \text{---} \boxed{\mathbf{T}} \text{---}$$

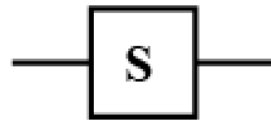
- General Phase Shift by δ :

$$\mathbf{Ph}_\delta = e^{i\delta} \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \quad |0\rangle \rightarrow e^{i\delta}|0\rangle \quad |1\rangle \rightarrow e^{i\delta}|1\rangle \quad \text{---} \boxed{\delta} \text{---}$$

Euler Elemental Rotation Decomposition

Examples for Relative Phase Shift Gates

- Global Phase Shifts of δ are Represented as $(e^{i\delta}\mathbf{I})$

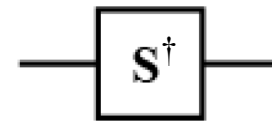


$$\mathbf{S} = \begin{bmatrix} 1 & 0 \\ 0 & i \end{bmatrix}$$

$$\mathbf{S} = \mathbf{R}_Z\left(\frac{\pi}{2}\right)\left(e^{i\frac{\pi}{4}}\mathbf{I}\right)$$

$$\mathbf{S} = \mathbf{R}_X\left(\frac{\pi}{2}\right)\mathbf{R}_Y\left(\frac{\pi}{2}\right)\mathbf{R}_X\left(-\frac{\pi}{2}\right)\left(e^{i\frac{\pi}{4}}\mathbf{I}\right)$$

$$\mathbf{S} = \mathbf{R}_Y\left(-\frac{\pi}{2}\right)\mathbf{R}_X\left(\frac{\pi}{2}\right)\mathbf{R}_Y\left(\frac{\pi}{2}\right)\left(e^{i\frac{\pi}{4}}\mathbf{I}\right)$$



$$\mathbf{S}^\dagger = \begin{bmatrix} 1 & 0 \\ 0 & -i \end{bmatrix}$$

$$\mathbf{S}^\dagger = \mathbf{R}_Z\left(-\frac{\pi}{2}\right)\left(e^{-i\frac{\pi}{4}}\mathbf{I}\right)$$

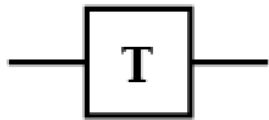
$$\mathbf{S}^\dagger = \mathbf{R}_X\left(\frac{\pi}{2}\right)\mathbf{R}_Y\left(-\frac{\pi}{2}\right)\mathbf{R}_X\left(-\frac{\pi}{2}\right)\left(e^{-i\frac{\pi}{4}}\mathbf{I}\right)$$

$$\mathbf{S}^\dagger = \mathbf{R}_Y\left(-\frac{\pi}{2}\right)\mathbf{R}_X\left(-\frac{\pi}{2}\right)\mathbf{R}_Y\left(\frac{\pi}{2}\right)\left(e^{-i\frac{\pi}{4}}\mathbf{I}\right)$$

Euler Elemental Rotation Decomposition

Examples for $\pi/8$ Gates

- Global Phase Shifts of δ are Represented as $(e^{i\delta}\mathbf{I})$

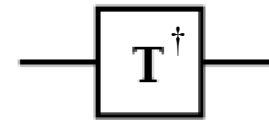


$$\mathbf{T} = \begin{bmatrix} 1 & 0 \\ 0 & e^{i(\pi/4)} \end{bmatrix}$$

$$\mathbf{T} = \mathbf{R}_Z\left(\frac{\pi}{4}\right)\left(e^{i\frac{\pi}{8}}\mathbf{I}\right)$$

$$\mathbf{T} = \mathbf{R}_X\left(\frac{\pi}{2}\right)\mathbf{R}_Y\left(\frac{\pi}{4}\right)\mathbf{R}_X\left(-\frac{\pi}{2}\right)\left(e^{i\frac{\pi}{8}}\mathbf{I}\right)$$

$$\mathbf{T} = \mathbf{R}_Y\left(-\frac{\pi}{2}\right)\mathbf{R}_X\left(\frac{\pi}{4}\right)\mathbf{R}_Y\left(\frac{\pi}{2}\right)\left(e^{i\frac{\pi}{8}}\mathbf{I}\right)$$



$$\mathbf{T}^\dagger = \begin{bmatrix} 1 & 0 \\ 0 & e^{-i(\pi/4)} \end{bmatrix}$$

$$\mathbf{T}^\dagger = \mathbf{R}_Z\left(-\frac{\pi}{4}\right)\left(e^{-i\frac{\pi}{8}}\mathbf{I}\right)$$

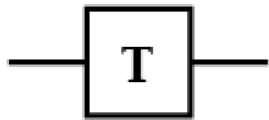
$$\mathbf{T}^\dagger = \mathbf{R}_X\left(\frac{\pi}{2}\right)\mathbf{R}_Y\left(-\frac{\pi}{4}\right)\mathbf{R}_X\left(-\frac{\pi}{2}\right)\left(e^{-i\frac{\pi}{8}}\mathbf{I}\right)$$

$$\mathbf{T}^\dagger = \mathbf{R}_Y\left(-\frac{\pi}{2}\right)\mathbf{R}_X\left(-\frac{\pi}{4}\right)\mathbf{R}_Y\left(\frac{\pi}{2}\right)\left(e^{-i\frac{\pi}{8}}\mathbf{I}\right)$$

Euler Elemental Rotation Decomposition

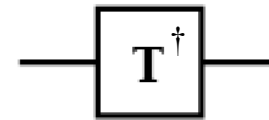
Examples for $\pi/8$ Gates

- Global Phase Shifts of δ are Represented as $(e^{i\delta}\mathbf{I})$



$$\mathbf{T} = \begin{bmatrix} 1 & 0 \\ 0 & e^{i(\pi/4)} \end{bmatrix}$$

$$\mathbf{T}^\dagger = \mathbf{SSST} = \mathbf{SZT}$$



$$\mathbf{T}^\dagger = \begin{bmatrix} 1 & 0 \\ 0 & e^{-i(\pi/4)} \end{bmatrix}$$

$$\mathbf{SS} = \begin{bmatrix} 1 & 0 \\ 0 & i \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & i \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} = \mathbf{Z}$$

$$\mathbf{T}^\dagger = \begin{bmatrix} 1 & 0 \\ 0 & i \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & e^{i\pi/4} \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ 0 & i \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & -e^{i\pi/4} \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ 0 & -ie^{i\pi/4} \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ 0 & e^{-i\pi/2} e^{i\pi/4} \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ 0 & e^{-i\pi/4} \end{bmatrix}$$

$$\mathbf{T} = \mathbf{SSST}^\dagger = \mathbf{SZT}^\dagger$$

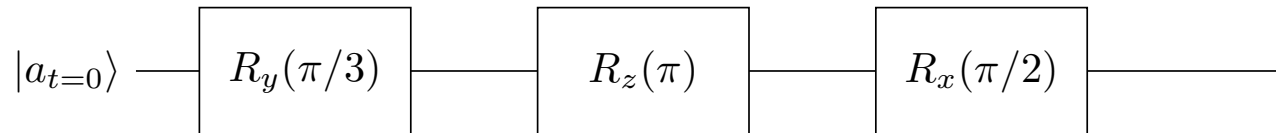
$$\mathbf{SS} = \begin{bmatrix} 1 & 0 \\ 0 & i \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & i \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} = \mathbf{Z}$$

$$\mathbf{T} = \begin{bmatrix} 1 & 0 \\ 0 & i \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & e^{-i\pi/4} \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ 0 & i \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & -e^{-i\pi/4} \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ 0 & -ie^{-i\pi/4} \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ 0 & e^{-i\pi/2} e^{-i\pi/4} \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ 0 & e^{-i3\pi/4} \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ 0 & e^{i\pi/4} \end{bmatrix}$$

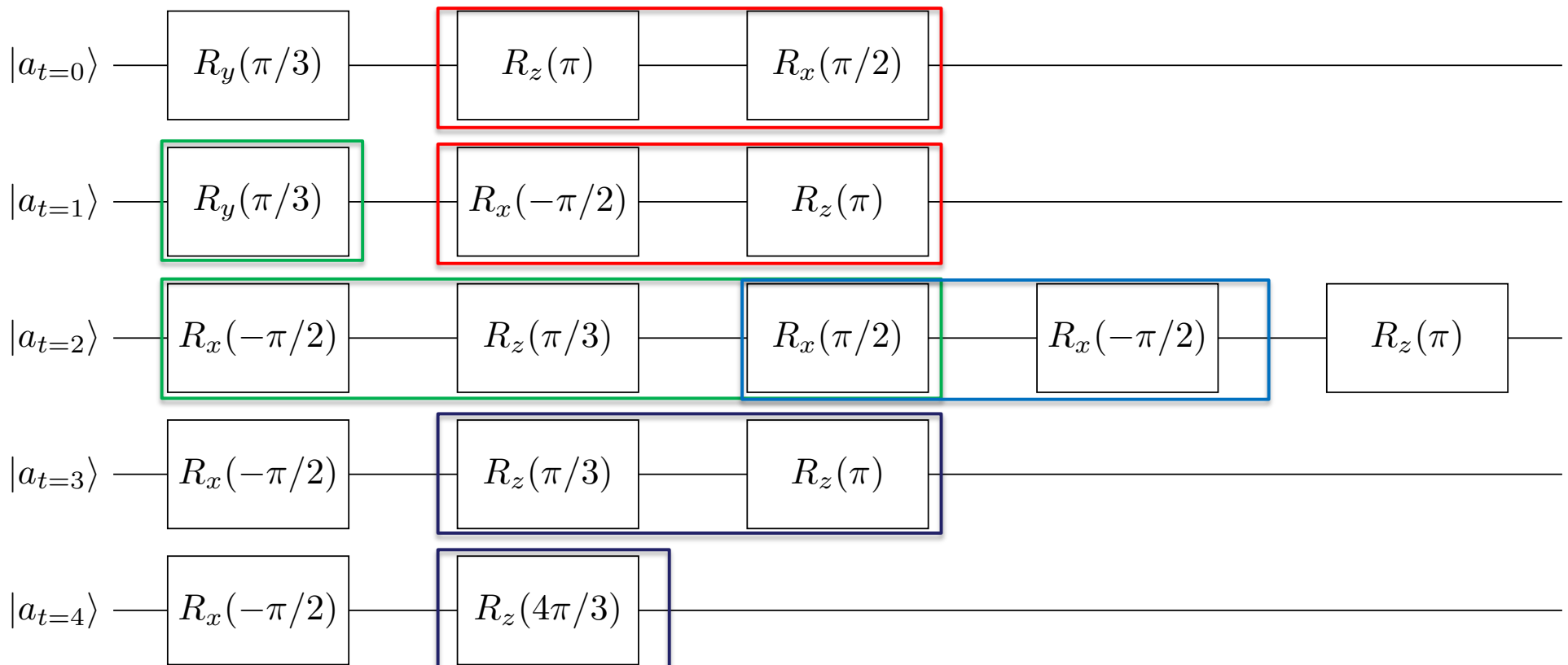
Example: Qubit Operator Optimization

Goal: Maximize $R_z(\delta)$ Rotations:

Key: **Step 1**
Step 2
Step 3
Step 4



Optimization Steps:

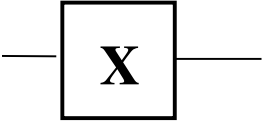
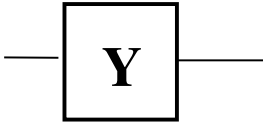
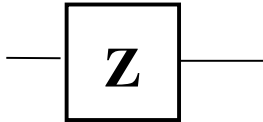


Single Qubit General Rotations

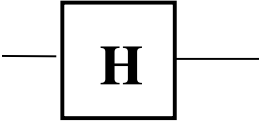
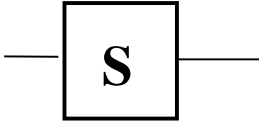
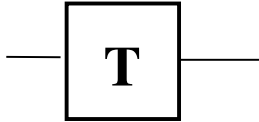
$\mathbf{R}_x(\theta)$		$\begin{bmatrix} \cos\frac{\theta}{2} & -i\sin\frac{\theta}{2} \\ -i\sin\frac{\theta}{2} & \cos\frac{\theta}{2} \end{bmatrix}$
$\mathbf{R}_y(\theta)$		$\begin{bmatrix} \cos\frac{\theta}{2} & -i\sin\frac{\theta}{2} \\ \sin\frac{\theta}{2} & \cos\frac{\theta}{2} \end{bmatrix}$
$\mathbf{R}_z(\theta)$		$\begin{bmatrix} e^{-i\theta/2} & 0 \\ 0 & e^{i\theta/2} \end{bmatrix}$

Single Qubit Pauli Operators

(180-degree rotations about axes in Bloch Sphere)

Pauli- X		$\begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$
Pauli- Y		$\begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix}$
Pauli- Z ($\pi/2$ gate)		$\begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$

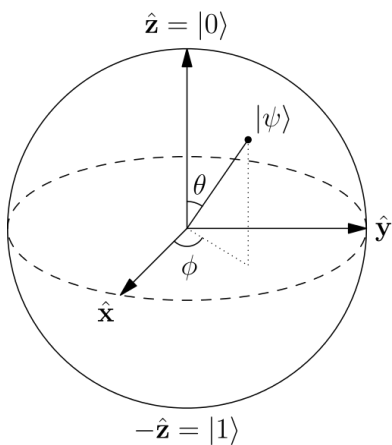
Other Single Qubit Operations

Hadamard		$\frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix}$
Phase ($\pi/4$ gate)		$\begin{bmatrix} 1 & 0 \\ 0 & i \end{bmatrix}$
“T” ($\pi/8$ gate)		$\begin{bmatrix} 1 & 0 \\ 0 & e^{i\pi/4} \end{bmatrix}$

Completely Generalized Rotation

$$0 \leq \theta \leq \pi$$

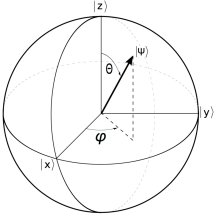
$$0 \leq \phi \leq 2\pi$$



$$\begin{aligned} \mathbf{R}(\theta, \phi) = |\psi\rangle\langle\psi| &= \begin{bmatrix} \cos\frac{\theta}{2} & \\ & e^{i\phi} \sin\frac{\theta}{2} \end{bmatrix} \otimes \begin{bmatrix} \cos\frac{\theta}{2} & e^{-i\phi} \sin\frac{\theta}{2} \\ & \end{bmatrix} \\ &= \begin{bmatrix} \cos^2\frac{\theta}{2} & e^{-i\phi} \cos\frac{\theta}{2} \sin\frac{\theta}{2} \\ e^{i\phi} \cos\frac{\theta}{2} \sin\frac{\theta}{2} & \sin^2\frac{\theta}{2} \end{bmatrix} \\ &= \frac{1}{2} \begin{bmatrix} 1 + \cos\theta & \cos\phi \sin\theta - i \sin\phi \sin\theta \\ \cos\phi \sin\theta + i \sin\phi \cos\theta & 1 - \cos\theta \end{bmatrix} \end{aligned}$$

Binary Basis States (Eigenstates)

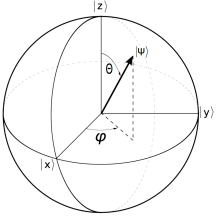
Photonic Example



Pauli Gate	Optical Name	Basis State (Jones Vector)	Bloch Sphere
$\sigma_1 = \sigma_x = \mathbf{X}$	Diag./Anti-dia. Slant-45 X-Basis	$ +\rangle = (0\rangle + 1\rangle) / \sqrt{2}$	$\hat{\mathbf{x}}$
		$ -\rangle = (0\rangle - 1\rangle) / \sqrt{2}$	$-\hat{\mathbf{x}}$
$\sigma_2 = \sigma_y = \mathbf{Y}$	LHP/RHP Circular Y-Basis	$ +i\rangle = (0\rangle + i 1\rangle) / \sqrt{2}$	$\hat{\mathbf{y}}$
		$ -i\rangle = (0\rangle - i 1\rangle) / \sqrt{2}$	$-\hat{\mathbf{y}}$
$\sigma_3 = \sigma_z = \mathbf{Z}$	Horizontal/ Vertical Computational Z-basis	$ 0\rangle$	$\hat{\mathbf{z}}$
		$ 1\rangle$	$-\hat{\mathbf{z}}$

Binary Basis States (Eigenstates)

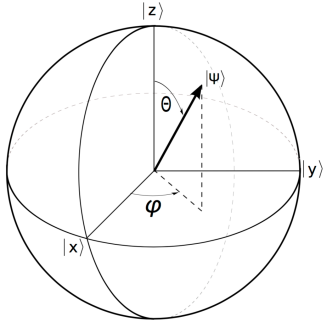
alternative notation



$$|L\rangle = (|H\rangle + i|V\rangle)/\sqrt{2} = |\odot\rangle \quad |R\rangle = (|H\rangle - i|V\rangle)/\sqrt{2} = |\ominus\rangle$$

Pauli Gate	Optical Name	Basis State (Jones Vector)	Jones Vector (standard)
$\sigma_1 = \sigma_x = \mathbf{X}$	Diag./Anti-dia. Slant-45 X-Basis	$ D\rangle = (H\rangle + V\rangle)/\sqrt{2}$	$\hat{\mathbf{x}} = 1/\sqrt{2} \begin{bmatrix} 1 & 1 \end{bmatrix}^T$
		$ A\rangle = (H\rangle - V\rangle)/\sqrt{2}$	$-\hat{\mathbf{x}} = 1/\sqrt{2} \begin{bmatrix} 1 & -1 \end{bmatrix}^T$
$\sigma_2 = \sigma_y = \mathbf{Y}$	LHP/RHP Circular Y-Basis	$ L\rangle = (H\rangle + i V\rangle)/\sqrt{2}$	$\hat{\mathbf{y}} = 1/\sqrt{2} \begin{bmatrix} 1 & i \end{bmatrix}^T$
		$ R\rangle = (H\rangle - i V\rangle)/\sqrt{2}$	$-\hat{\mathbf{y}} = 1/\sqrt{2} \begin{bmatrix} 1 & -i \end{bmatrix}^T$
$\sigma_3 = \sigma_z = \mathbf{Z}$	Horizontal/Ver tical Computational Z-basis	$ 0\rangle = H\rangle$	$\hat{\mathbf{z}} = \begin{bmatrix} 1 & 0 \end{bmatrix}^T$
		$ 1\rangle = V\rangle$	$-\hat{\mathbf{z}} = \begin{bmatrix} 0 & 1 \end{bmatrix}^T$

Single Qubit Operator Relationships



Pauli Gate	Eigenvalue	Eigenvector
$\sigma_1 = \sigma_X = \mathbf{X}$	+1	$ +\rangle = (0\rangle + 1\rangle)/\sqrt{2}$
$\sigma_1 = \sigma_X = \mathbf{X}$	-1	$ -\rangle = (0\rangle - 1\rangle)/\sqrt{2}$
$\sigma_2 = \sigma_Y = \mathbf{Y}$	+1	$ +i\rangle = (0\rangle + i 1\rangle)/\sqrt{2}$
$\sigma_2 = \sigma_Y = \mathbf{Y}$	-1	$ -i\rangle = (0\rangle - i 1\rangle)/\sqrt{2}$
$\sigma_3 = \sigma_Z = \mathbf{Z}$	+1	$ 0\rangle$
$\sigma_3 = \sigma_Z = \mathbf{Z}$	-1	$ 1\rangle$

(Diagonal)

(Anti-Diagonal)

Quantum Logic Gates and Circuits

Multi-qubit “Gates” or Operators

Multi Qubit Systems (Circuits)

- Multi qubit systems are represented in terms of a “product quantum state”
- Consider a System of Two qubits, the state of this system is a superposition of:

$$|pq\rangle = \alpha|00\rangle + \beta|01\rangle + \chi|10\rangle + \delta|11\rangle$$

α	←	Amplitude for 00
β	←	Amplitude for 01
χ	←	Amplitude for 10
δ	←	Amplitude for 11

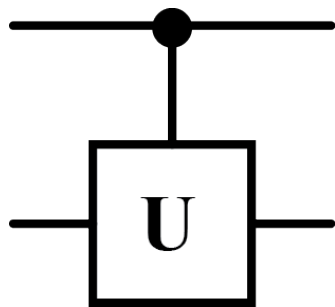
2-Qubit (Controlled) Gates

- General Model where Top Qubit Allows Bottom Qubit to Evolve with U (or not)
- Recall that:



$$U = \begin{bmatrix} u_{00} & u_{01} \\ u_{10} & u_{11} \end{bmatrix}$$

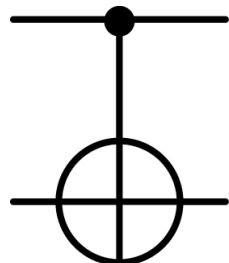
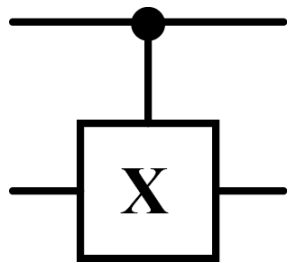
- Controlled-U Gate



$$C_U = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & u_{00} & u_{01} \\ 0 & 0 & u_{10} & u_{11} \end{bmatrix}$$

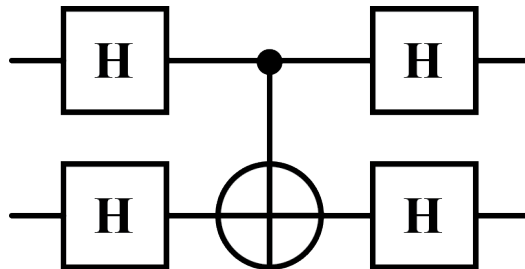
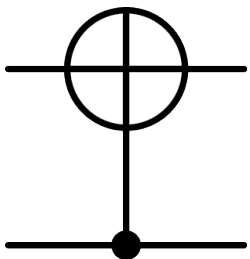
Two-Qubit Gates

- Bloch Sphere can only Represent a Single Qubit
- Feynman, Controlled-NOT, CNOT, Controlled-**X**



$$\text{CNOT} = C_x = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix}$$

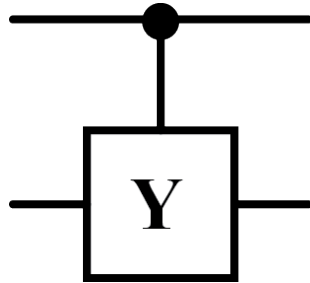
- Inverted SWAP, Inverted Controlled-**X**,
(technically not a different gate)



$$\text{ISWAP} = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \end{bmatrix}$$

Two-Qubit Gates (cont)

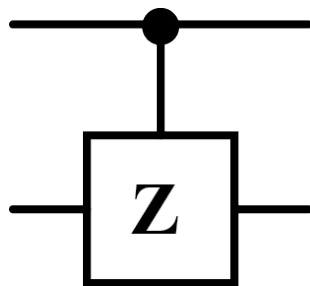
- Controlled-Y



$$C_Y = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & -i \\ 0 & 0 & i & 0 \end{bmatrix}$$

- Controlled-Z, Controlled-Sign, Controlled-Phase, **CPHASE**, **CSIGN**

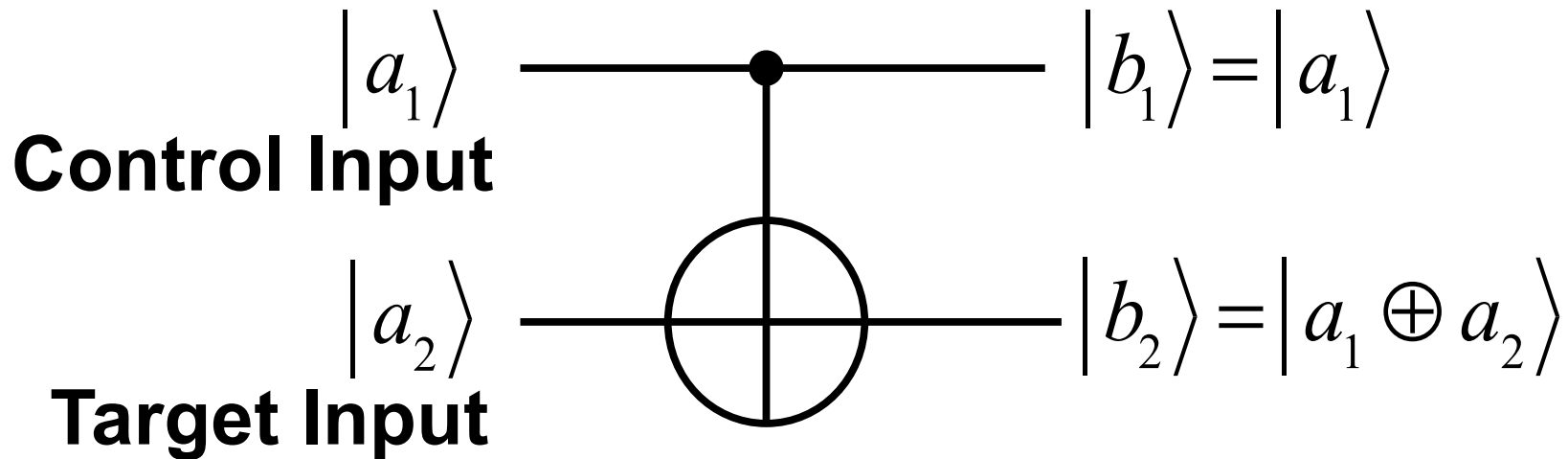
– arises in linear optical computing



$$C_Z = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \end{bmatrix}$$

Controlled-NOT (C_{NOT}) Gate

(aka Feynman, Controlled-X, Quantum XOR Gate)



$$\begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix}$$

$ a_1\rangle$	$ a_2\rangle$	$ b_1\rangle$	$ b_2\rangle$
$ 0\rangle$	$ 0\rangle$	$ 0\rangle$	$ 0\rangle$
$ 0\rangle$	$ 1\rangle$	$ 0\rangle$	$ 1\rangle$
$ 1\rangle$	$ 0\rangle$	$ 1\rangle$	$ 1\rangle$
$ 1\rangle$	$ 1\rangle$	$ 1\rangle$	$ 0\rangle$

Controlled-NOT (C_{NOT}) Gate

(aka Feynman, Controlled-X, Quantum XOR Gate)

$$C_{NOT} = C_X = |00\rangle\langle 00| + |01\rangle\langle 01| + |10\rangle\langle 11| + |11\rangle\langle 10|$$

$$|00\rangle = |0\rangle \otimes |0\rangle = \begin{bmatrix} 1 \\ 0 \end{bmatrix} \otimes \begin{bmatrix} 1 \\ 0 \end{bmatrix} = \begin{bmatrix} 1 \\ 0 \\ 0 \\ 0 \end{bmatrix}$$

$$\langle 00| = \langle 0| \otimes \langle 0| = \begin{bmatrix} 1 & 0 \end{bmatrix} \otimes \begin{bmatrix} 1 & 0 \end{bmatrix} = \begin{bmatrix} 1 & 0 & 0 & 0 \end{bmatrix}$$

$$|01\rangle = |0\rangle \otimes |1\rangle = \begin{bmatrix} 1 \\ 0 \end{bmatrix} \otimes \begin{bmatrix} 0 \\ 1 \end{bmatrix} = \begin{bmatrix} 0 \\ 1 \\ 0 \\ 0 \end{bmatrix}$$

$$\langle 01| = \langle 0| \otimes \langle 1| = \begin{bmatrix} 1 & 0 \end{bmatrix} \otimes \begin{bmatrix} 0 & 1 \end{bmatrix} = \begin{bmatrix} 0 & 1 & 0 & 0 \end{bmatrix}$$

Controlled-NOT (C_{NOT}) Gate

(aka Feynman, Controlled-X, Quantum XOR Gate)

$$C_x = |00\rangle\langle 00| + |01\rangle\langle 01| + |10\rangle\langle 11| + |11\rangle\langle 10|$$

$$|10\rangle = |1\rangle \otimes |0\rangle = \begin{bmatrix} 0 \\ 1 \end{bmatrix} \otimes \begin{bmatrix} 1 \\ 0 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 1 \\ 0 \end{bmatrix}$$

$$\langle 10| = \langle 1| \otimes \langle 0| = \begin{bmatrix} 0 & 1 \end{bmatrix} \otimes \begin{bmatrix} 1 & 0 \end{bmatrix} = \begin{bmatrix} 0 & 0 & 1 & 0 \end{bmatrix}$$

$$|11\rangle = |1\rangle \otimes |1\rangle = \begin{bmatrix} 0 \\ 1 \end{bmatrix} \otimes \begin{bmatrix} 0 \\ 1 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \\ 1 \end{bmatrix}$$

$$\langle 11| = \langle 1| \otimes \langle 1| = \begin{bmatrix} 0 & 1 \end{bmatrix} \otimes \begin{bmatrix} 0 & 1 \end{bmatrix} = \begin{bmatrix} 0 & 0 & 0 & 1 \end{bmatrix}$$

Controlled-NOT (C_{NOT}) Gate

(aka Feynman, Controlled-X, Quantum XOR Gate)

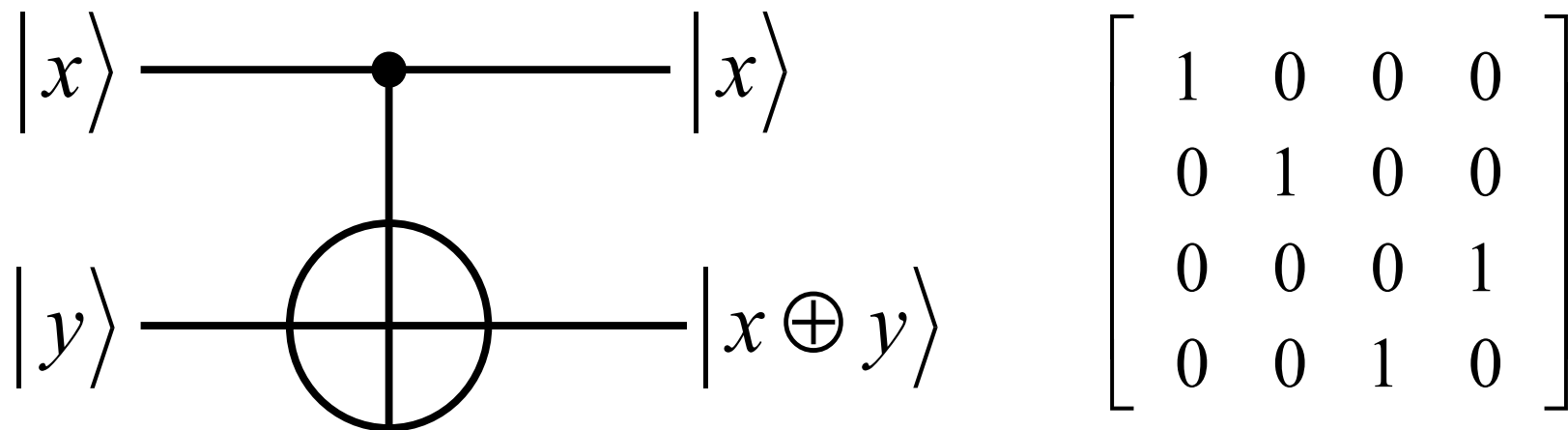
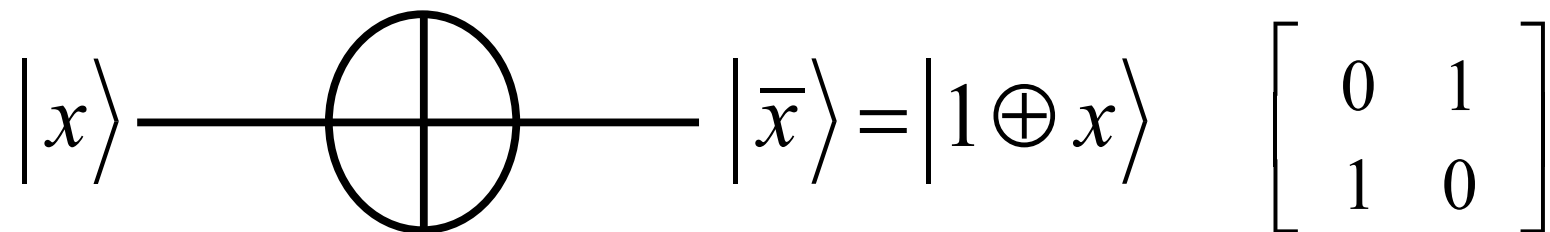
$$C_x = |00\rangle\langle 00| + |01\rangle\langle 01| + |10\rangle\langle 11| + |11\rangle\langle 10|$$

$$C_x = \begin{bmatrix} 1 \\ 0 \\ 0 \\ 0 \end{bmatrix} \otimes [1 \ 0 \ 0 \ 0] + \begin{bmatrix} 0 \\ 1 \\ 0 \\ 0 \end{bmatrix} \otimes [0 \ 1 \ 0 \ 0] \\ + \begin{bmatrix} 0 \\ 0 \\ 1 \\ 0 \end{bmatrix} \otimes [0 \ 0 \ 0 \ 1] + \begin{bmatrix} 0 \\ 0 \\ 0 \\ 1 \end{bmatrix} \otimes [0 \ 0 \ 1 \ 0]$$

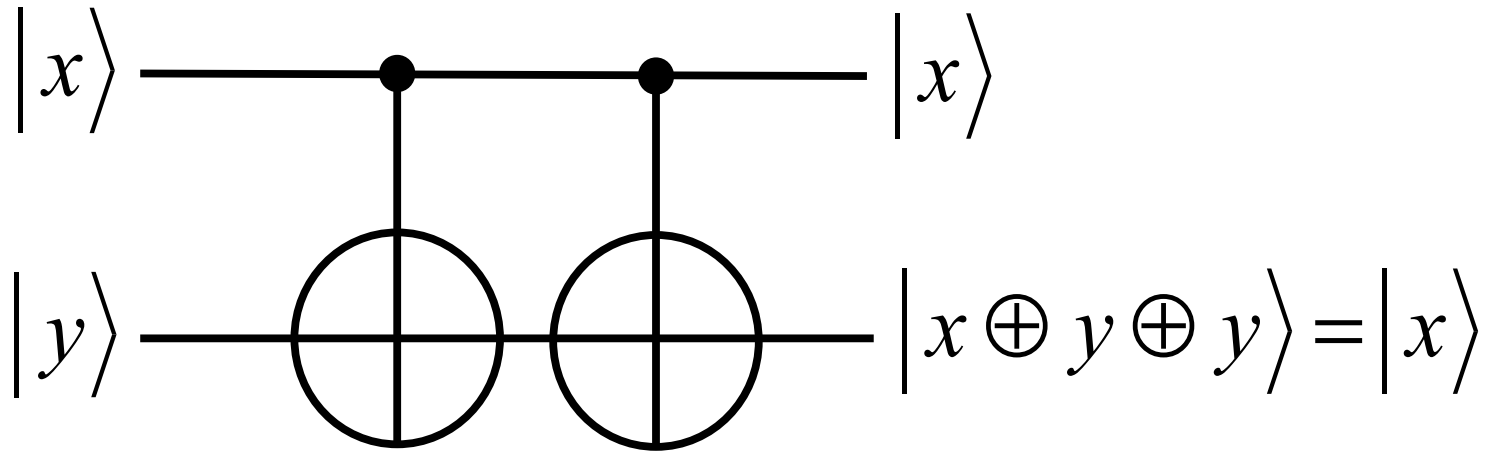
$$C_x = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix} + \begin{bmatrix} 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix} + \begin{bmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & 0 \end{bmatrix} + \begin{bmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 \end{bmatrix} = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix}$$

X and C_X Gates

(aka Pauli-X/NOT and Controlled-NOT)



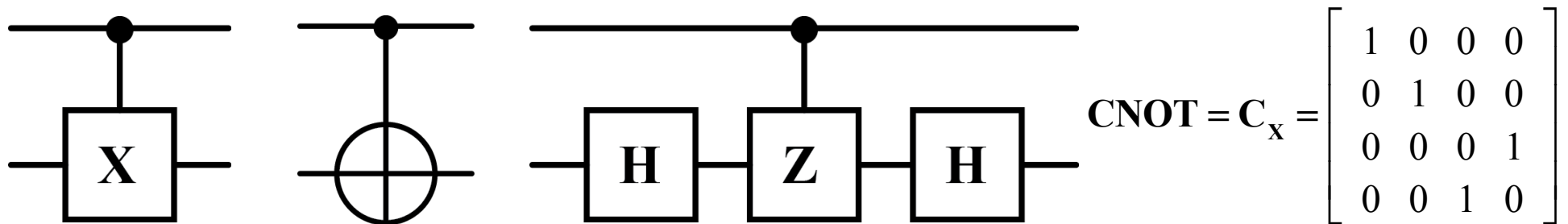
Reversibility of C_x Gate



$$\begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix} \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix} = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix}$$

CNOT Using CPHASE and H

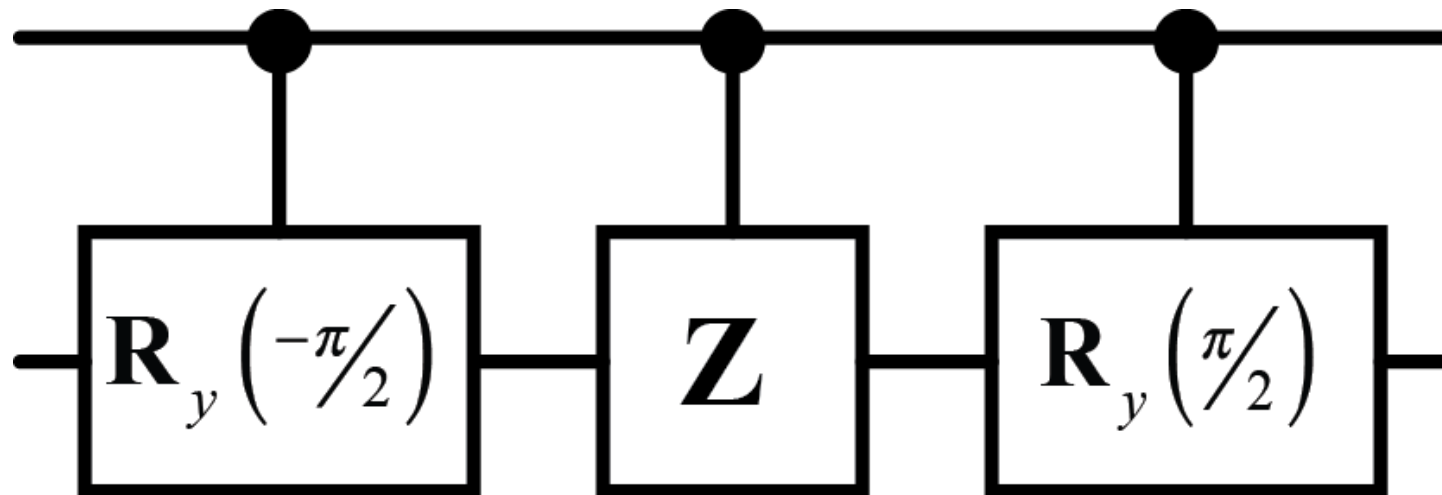
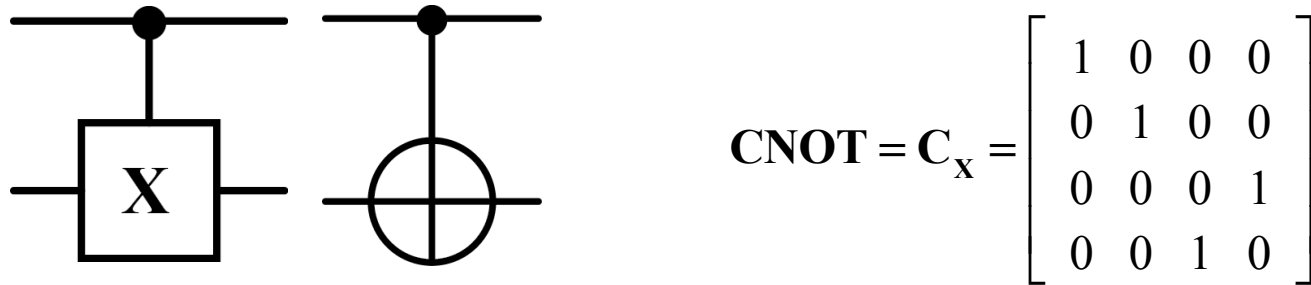
- Can Build CNOT with CPHASE and H



$$\begin{aligned}
 (\mathbf{I} \otimes \mathbf{H})C_Z(\mathbf{I} \otimes \mathbf{H}) &= \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 & 0 & 0 \\ 1 & -1 & 0 & 0 \\ 0 & 0 & 1 & 1 \\ 0 & 0 & 1 & -1 \end{bmatrix} \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \end{bmatrix} \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 & 0 & 0 \\ 1 & -1 & 0 & 0 \\ 0 & 0 & 1 & 1 \\ 0 & 0 & 1 & -1 \end{bmatrix} \\
 &= \frac{1}{2} \begin{bmatrix} 1 & 1 & 0 & 0 \\ 1 & -1 & 0 & 0 \\ 0 & 0 & 1 & 1 \\ 0 & 0 & 1 & -1 \end{bmatrix} \begin{bmatrix} 1 & 1 & 0 & 0 \\ 1 & -1 & 0 & 0 \\ 0 & 0 & 1 & 1 \\ 0 & 0 & -1 & 1 \end{bmatrix} = \frac{1}{2} \begin{bmatrix} 2 & 0 & 0 & 0 \\ 0 & 2 & 0 & 0 \\ 0 & 0 & 0 & 2 \\ 0 & 0 & 2 & 0 \end{bmatrix} = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix}
 \end{aligned}$$

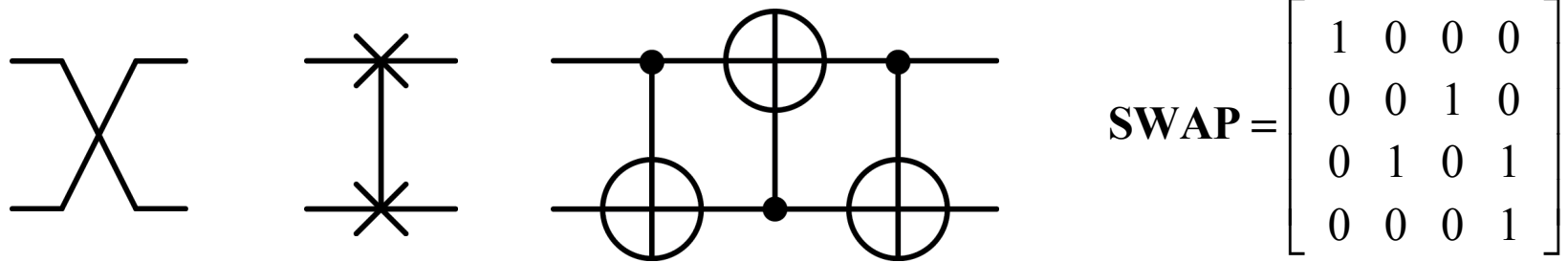
CNOT Using CPHASE and R_y

- Can Build CNOT with CPHASE and R_y



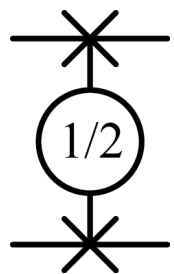
Two-Qubit Gates (cont)

- SWAP Gate



- Square Root of $\text{SWAP}_{\sqrt{\text{SWAP}}}$

– spintronic-based circuits (“exchange interaction”)

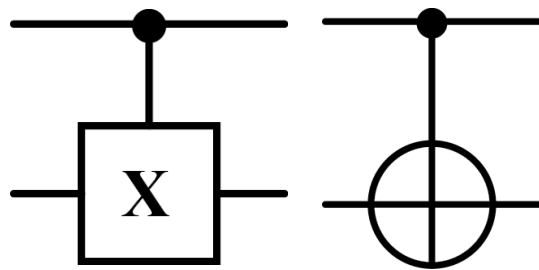


$$Q_{\text{SWP}} = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & \frac{1}{2}(1+i) & \frac{1}{2}(1-i) & 0 \\ 0 & \frac{1}{2}(1-i) & \frac{1}{2}(1+i) & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix}$$

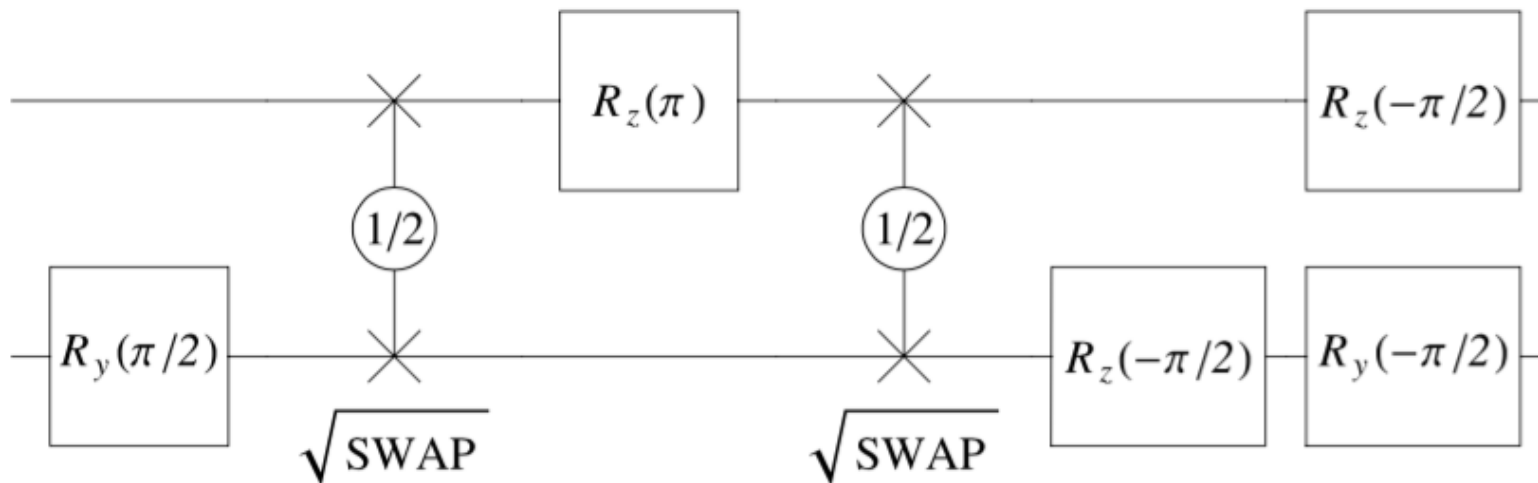
CNOT Using CPHASE and Sq. root

SWAP, R_y , and R_z

- Can Build CNOT with CPHASE, Sq. root SWAP, R_y and R_z



$$\text{CNOT} = C_x = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix}$$



Two-Qubit Gates (cont)

- *i*SWAP Gate

- arises in superconducting quantum computing with Hamiltonians that implement the so-called **XY Model**

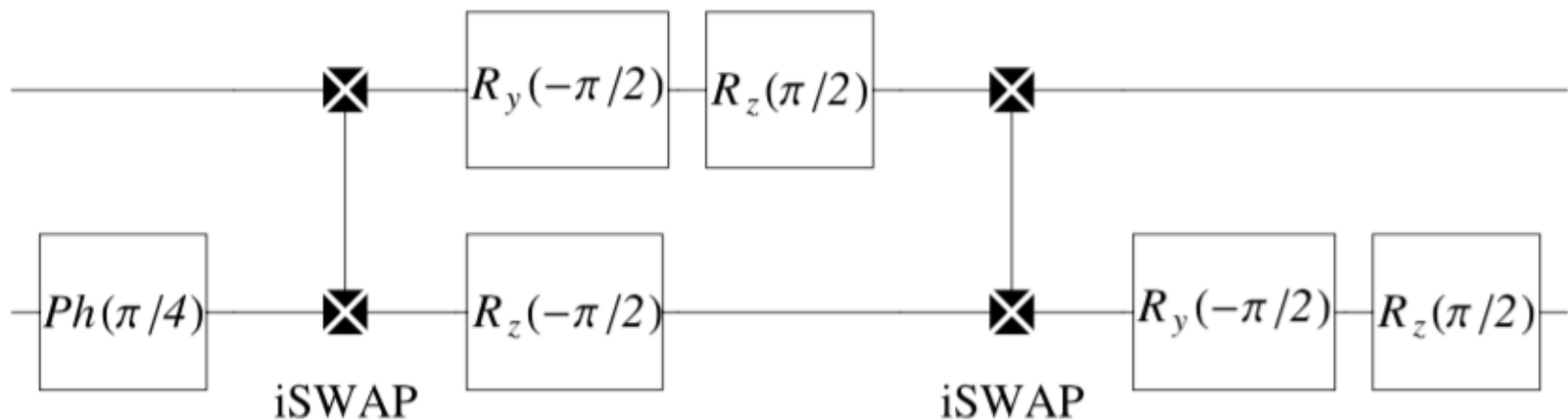
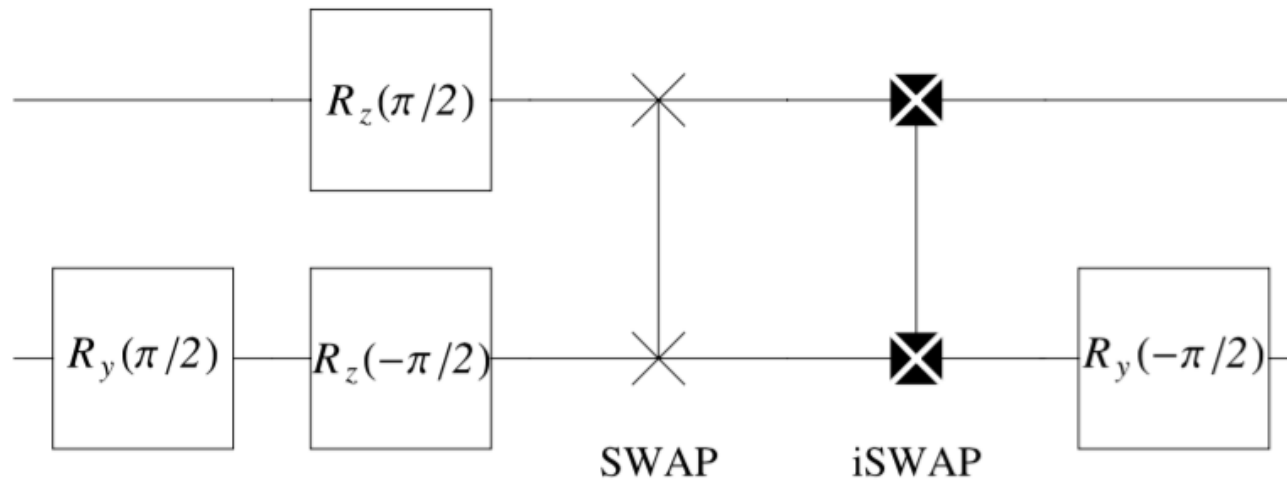
$$i\text{SWAP} = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & i & 0 \\ 0 & i & 0 & 1 \\ 0 & 0 & 0 & 1 \end{bmatrix}$$

- Square Root of **SWAP**^α

- spintronic-based circuits (“exchange interaction”)
- duration of exchange determines the exponent

$$\text{SWAP}^\alpha = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & \frac{1}{2}(1 + e^{i\pi\alpha}) & \frac{1}{2}(1 - e^{i\pi\alpha}) & 0 \\ 0 & \frac{1}{2}(1 - e^{i\pi\alpha}) & \frac{1}{2}(1 + e^{i\pi\alpha}) & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix}$$

CNOT Using CPHASE and Sq. root SWAP, i SWAP, R_y , and R_z



Two-Qubit Gates (cont)

- **BARENCO Gate**

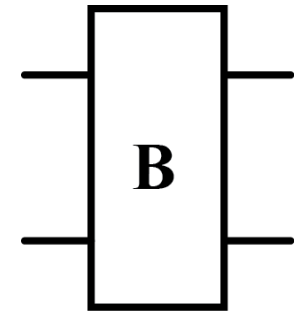
- ϕ , α , and θ are fixed irrational multiples of each and with π

- In Form of a Controlled-U Gate

$$\mathbf{BARENCO} = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & e^{i\alpha} \cos(\theta) & -ie^{i(\alpha-\phi)} \sin(\theta) \\ 0 & 0 & -ie^{i(\alpha-\phi)} \sin(\theta) & e^{i\alpha} \cos(\theta) \end{bmatrix}$$

Berkeley **B** Gate

- Hamiltonian is: $\mathcal{H} = \frac{\pi}{8}(2\mathbf{X} \otimes \mathbf{X} + \mathbf{Y} \otimes \mathbf{Y})$
- Transfer Matrix: $\mathbf{B} = e^{i\mathcal{H}} = e^{i\frac{\pi}{8}(2\mathbf{X} \otimes \mathbf{X} + \mathbf{Y} \otimes \mathbf{Y})}$



$$\mathbf{B} = \begin{bmatrix} \cos\left(\frac{\pi}{8}\right) & 0 & 0 & i\sin\left(\frac{\pi}{8}\right) \\ 0 & \cos\left(\frac{3\pi}{8}\right) & i\sin\left(\frac{3\pi}{8}\right) & 0 \\ 0 & i\sin\left(\frac{3\pi}{8}\right) & \cos\left(\frac{3\pi}{8}\right) & 0 \\ i\sin\left(\frac{\pi}{8}\right) & 0 & 0 & \cos\left(\frac{\pi}{8}\right) \end{bmatrix} = \frac{\sqrt{2-\sqrt{2}}}{2} \begin{bmatrix} 1+\sqrt{2} & 0 & 0 & i \\ 0 & 1 & i(1+\sqrt{2}) & 0 \\ 0 & i(1+\sqrt{2}) & 1 & 0 \\ i & 0 & 0 & 1+\sqrt{2} \end{bmatrix}$$

Quantum Logic Gates and Circuits

Three-qubit Gates

Three-Qubit Gates

- Toffoli Gate $\text{Toff}|a,b,c\rangle = |a,b,ab \oplus c\rangle$
 - universal for reversible with Pauli-X and CNOT
- Peres Gate $\text{P}|a,b,c\rangle = |a,a \oplus b,ab \oplus c\rangle$
- Fredkin (Controlled-SWAP)
- Deutsch Gate
 - universal

Deutsch Gate

- θ is any constant angle, such that $\frac{2\theta}{\pi}$ is an irrational number.

$$\mathbf{DEUTSCH} = \begin{bmatrix} 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & i \cos(\theta) & \sin(\theta) \\ 0 & 0 & 0 & 0 & 0 & 0 & \sin(\theta) & i \cos(\theta) \end{bmatrix}$$

Generalized Toffoli Gate

- Generalized Toffoli Gate
 - Some Papers Refer to Generalized Toffoli as Toffoli
- NOT (Pauli-X), CNOT (Feynman), Toffoli are Universal Set of Reversible Logic Gates
 - Reversible Logic assumes Only Basis States are Used
 - Form of Classical Switching Theory Based Computation (e.g., electronic adiabatic circuits)

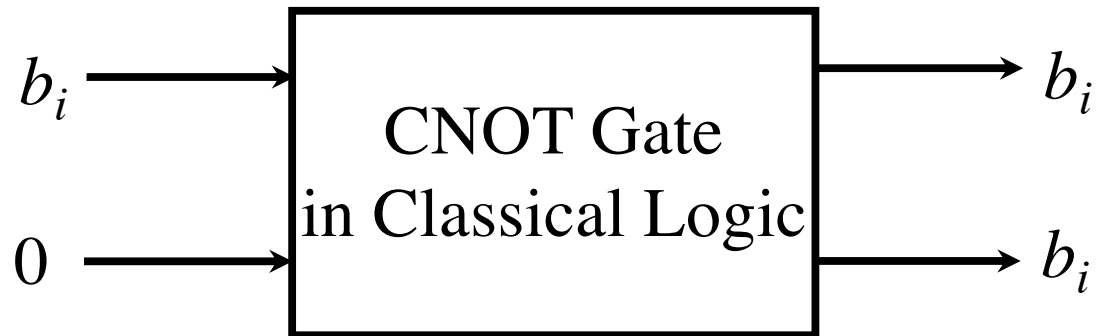
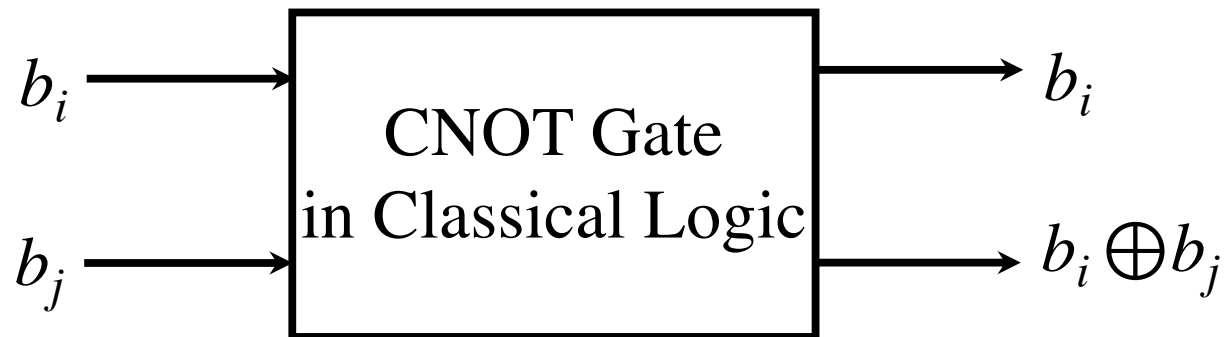
Decompositions: Barenco's Theorem

- All Multi-qubit Gates can be Decomposed into a Cascade of Two-qubit Controlled and Single Qubit Rotations
- “Quantum Cost” is the Total Number of Single and Two-qubit Controlled Gates Required to Perform a Computation
- Other Decompositions are Known
- New Decompositions are an Open Area of Research

Quantum Logic Gates and Circuits

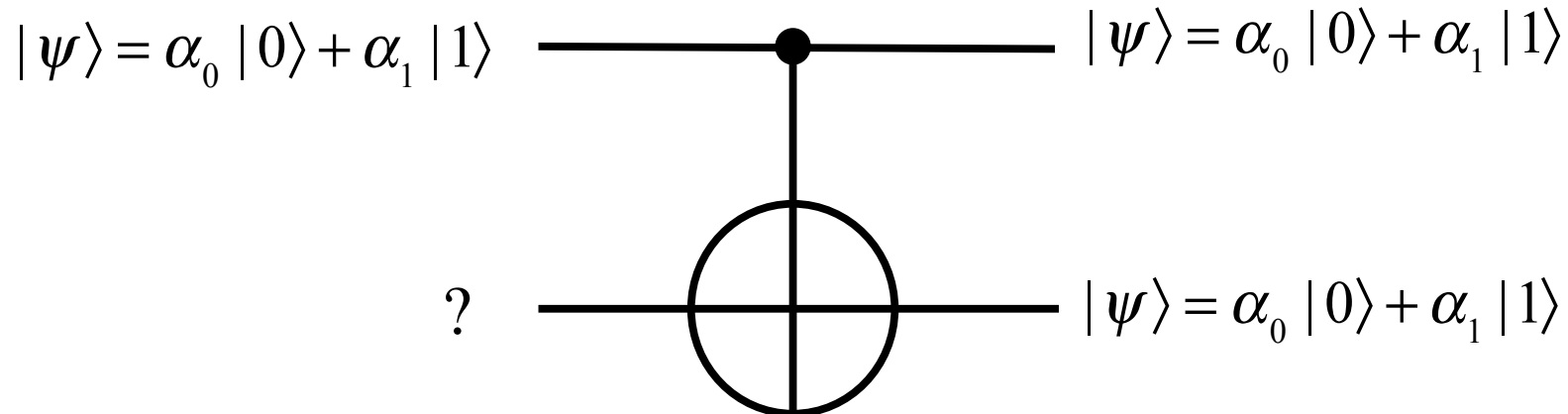
The “No-Cloning” Theorem

Bit Copying



Can Produce a Fanout (copy) with CNOT Gate implemented in CLASSICAL Logic

Is Qubit Copying Possible?



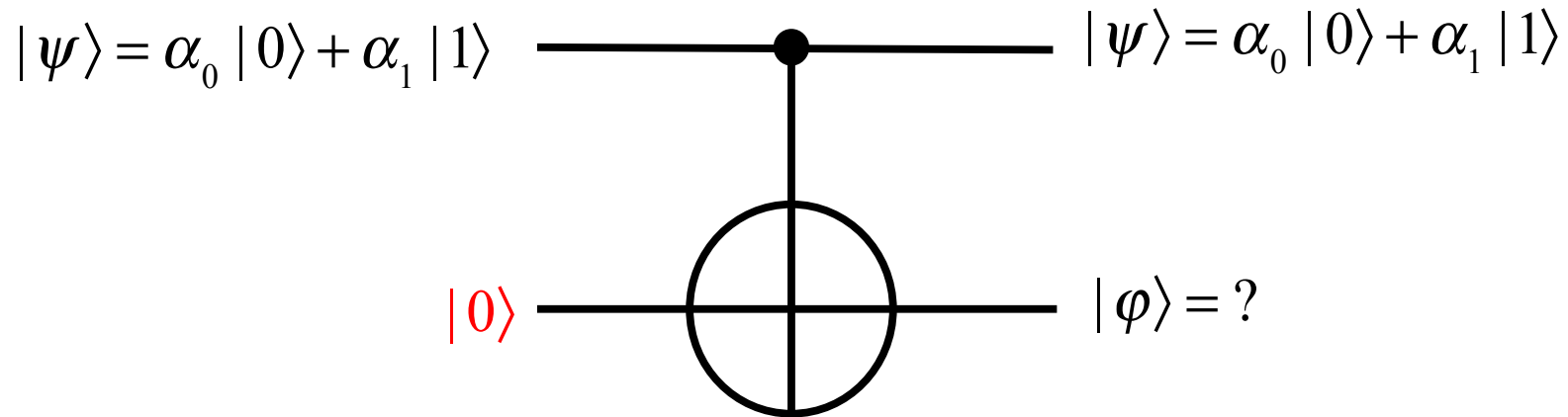
- If **Qubit Copying Possible**, Output Quantum State is:

$$\begin{aligned}
 |\psi\rangle \otimes |\psi\rangle &= |\psi\psi\rangle = (\alpha_0 |0\rangle + \alpha_1 |1\rangle)(\alpha_0 |0\rangle + \alpha_1 |1\rangle) \\
 &= |\psi\psi\rangle = \alpha_0^2 |00\rangle + \alpha_0\alpha_1 |01\rangle + \alpha_0\alpha_1 |10\rangle + \alpha_1^2 |11\rangle
 \end{aligned}$$

- Or Using Tensor Product:

$$|\psi\rangle \otimes |\psi\rangle = \begin{bmatrix} \alpha_0 \\ \alpha_1 \end{bmatrix} \otimes \begin{bmatrix} \alpha_0 \\ \alpha_1 \end{bmatrix} = \begin{bmatrix} \alpha_0^2 \\ \alpha_0\alpha_1 \\ \alpha_0\alpha_1 \\ \alpha_1^2 \end{bmatrix}$$

Qubit Copying?



- Input Quantum State:

$$|\psi\rangle \otimes |0\rangle = |\psi 0\rangle = (\alpha_0 |0\rangle + \alpha_1 |1\rangle)(|0\rangle) = \alpha_0 |00\rangle + \alpha_1 |10\rangle$$

- Output Quantum State:

$$|\psi\varphi\rangle = \mathbf{CNOT} |\psi 0\rangle = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix} \begin{bmatrix} \alpha_0 \\ 0 \\ \alpha_1 \\ 0 \end{bmatrix} = \begin{bmatrix} \alpha_0 \\ 0 \\ 0 \\ \alpha_1 \end{bmatrix} = \alpha_0 |00\rangle + \alpha_1 |11\rangle$$

- Qubit is **NOT COPIED**:

$$\alpha_0 |00\rangle + \alpha_1 |11\rangle \neq \alpha_0^2 |00\rangle + \alpha_0 \alpha_1 |01\rangle + \alpha_0 \alpha_1 |10\rangle + \alpha_1^2 |11\rangle$$

No-Cloning Theorem

- Transformations Carried out by Quantum Gates are Unitary
- Cloning of a Quantum State is a Non-Unitary and Non-Linear Process
- Information Point of View is Two Copies of Quantum State Embody MORE Information than Available in One Copy
- IS POSSIBLE to Clone States After a Measurement has Occurred
- No Cloning Theorem Applies to UNKNOWN Quantum States

No-Cloning Theorem

- Proof by Contradiction
- Assume a Cloning Gate Exists Characterized by Transform Matrix G
- Assume Two Orthogonal Quantum States are Cloned, One after the Other

$$G(|\psi\rangle \otimes |0\rangle) = |\psi\rangle \otimes |\psi\rangle \quad \text{OR} \quad G(|\psi 0\rangle) = |\psi\psi\rangle$$

$$G(|\varphi\rangle \otimes |0\rangle) = |\varphi\rangle \otimes |\varphi\rangle \quad \text{OR} \quad G(|\varphi 0\rangle) = |\varphi\varphi\rangle$$

- These Two Equations State that G Performs a Cloning Operation When the Second Qubit is ket-zero

No-Cloning Theorem

- Consider Another Quantum State:

$$|\xi\rangle = (1/\sqrt{2})(|\psi\rangle + |\varphi\rangle)$$

- Applying the Cloning Transform:

$$\mathbf{G}(|\xi 0\rangle) = \frac{1}{\sqrt{2}}[\mathbf{G}(|\psi 0\rangle) + \mathbf{G}(|\varphi 0\rangle)] = \frac{1}{\sqrt{2}}[|\psi\psi\rangle + |\varphi\varphi\rangle]$$

- If G is Truly a Cloning Gate then:

$$\mathbf{G}(|\xi 0\rangle) = |\xi\xi\rangle$$

- But:

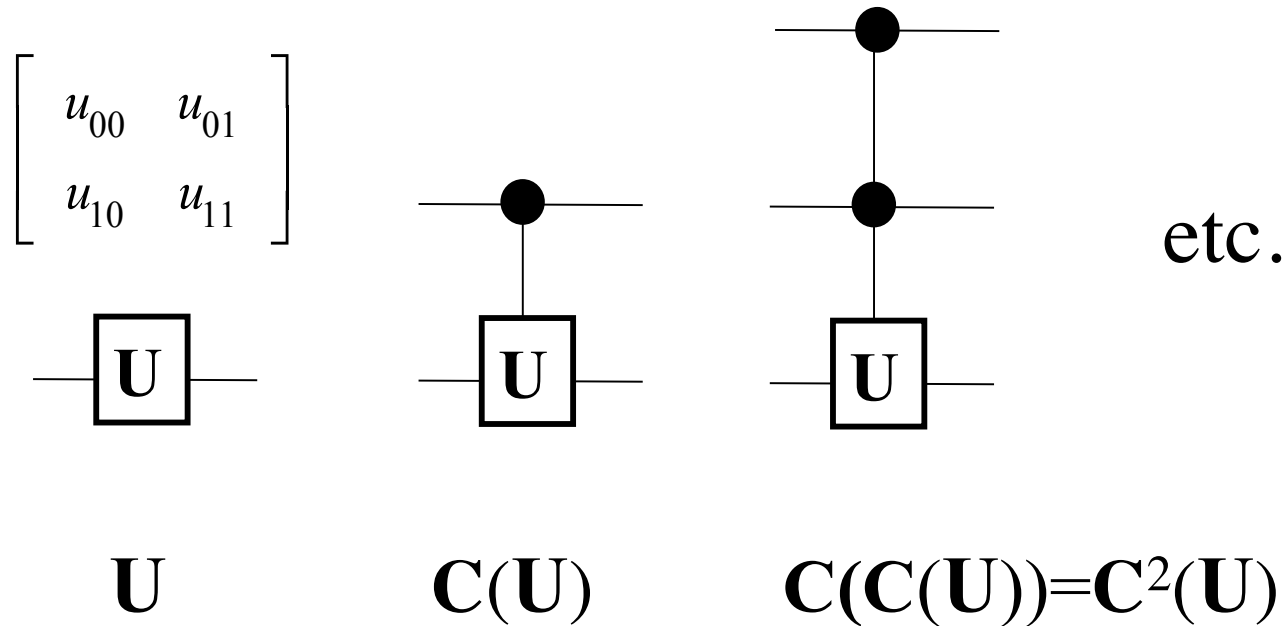
$$|\xi\xi\rangle = \left(\frac{|\psi\rangle + |\varphi\rangle}{\sqrt{2}}\right)\left(\frac{|\psi\rangle + |\varphi\rangle}{\sqrt{2}}\right) = \frac{1}{2}(|\psi\psi\rangle + |\psi\varphi\rangle + |\varphi\psi\rangle + |\varphi\varphi\rangle)$$

CONTRADICTION!!!



Quantum Gates*

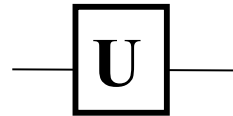
- General controlled gates that control some 1-qubit unitary operation U are useful



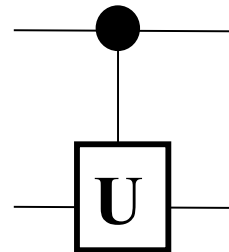
Quantum Gates

- General controlled gates that control some 1-qubit unitary operation U are useful

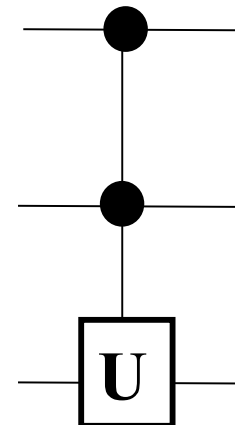
$$\begin{bmatrix} u_{00} & u_{01} \\ u_{10} & u_{11} \end{bmatrix}$$



U



$C(U)$



$C(C(U))=C^2(U)$

$$C(U) = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & u_{00} & u_{01} \\ 0 & 0 & u_{10} & u_{11} \end{bmatrix}$$

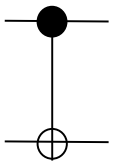
$$C^2(U) = \begin{bmatrix} 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & u_{00} & u_{01} \\ 0 & 0 & 0 & 0 & 0 & 0 & u_{10} & u_{11} \end{bmatrix}$$

Quantum Gates*

Discrete Universal Gate Set Example

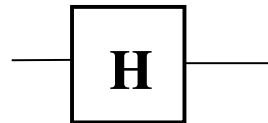
- **Example 1:** Four-member “standard” gate set, $\{C_X, H, S, T\}$

$$\begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix}$$



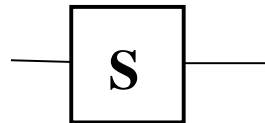
C_X

$$\frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix}$$



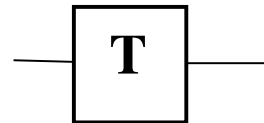
Hadamard

$$\begin{bmatrix} 1 & 0 \\ 0 & i \end{bmatrix}$$



Phase

$$\begin{bmatrix} 1 & 0 \\ 0 & e^{i\pi/4} \end{bmatrix} = e^{i\pi/8} \begin{bmatrix} e^{-i\pi/8} & 0 \\ 0 & e^{i\pi/8} \end{bmatrix}$$



$\pi/8$ (T) gate

- **Example 2:** $\{X, C_X, H, \text{Toffoli}\}$

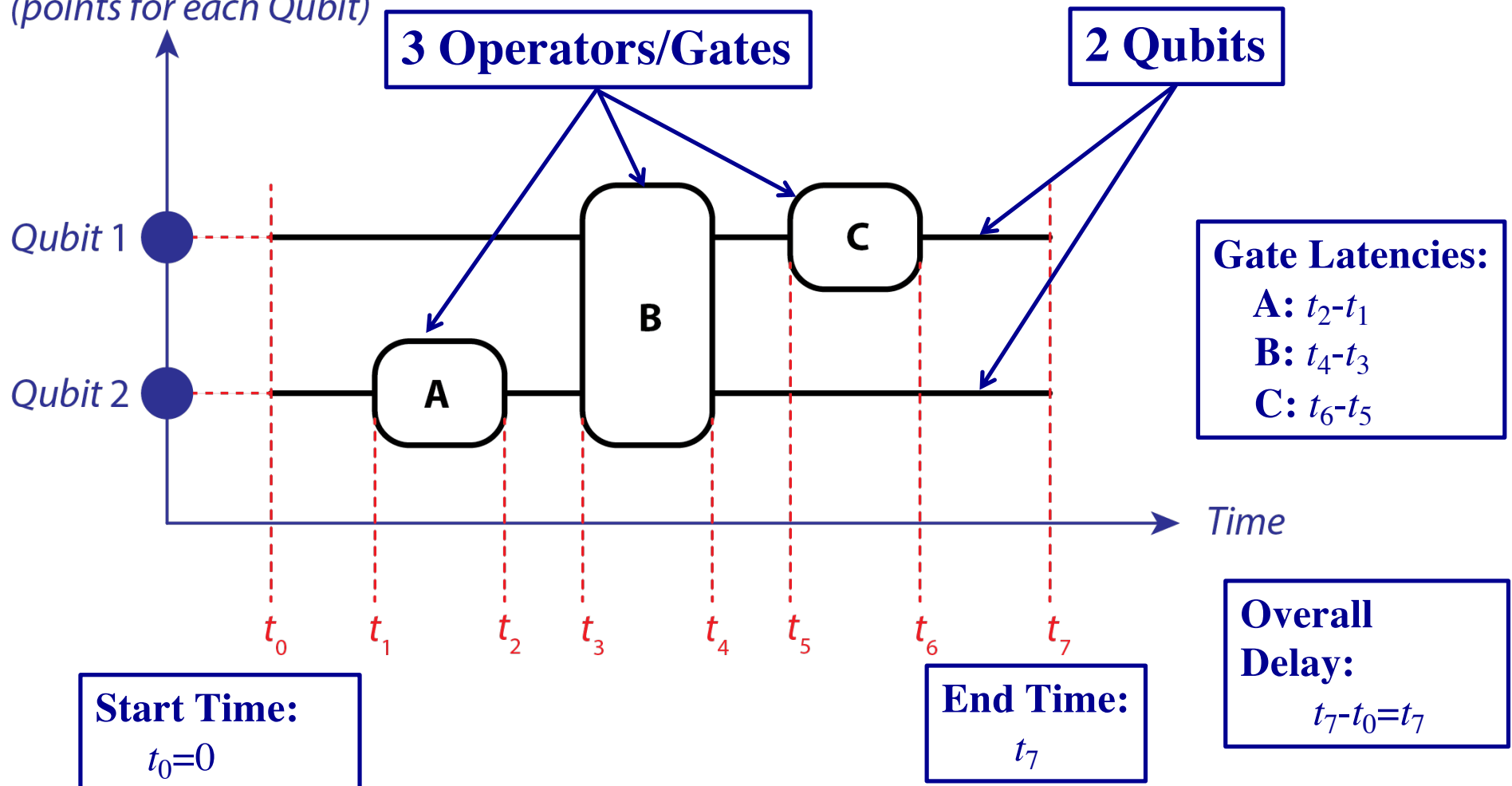
Quantum Logic Gates and Circuits

Quantum Programs

Graphical Representation

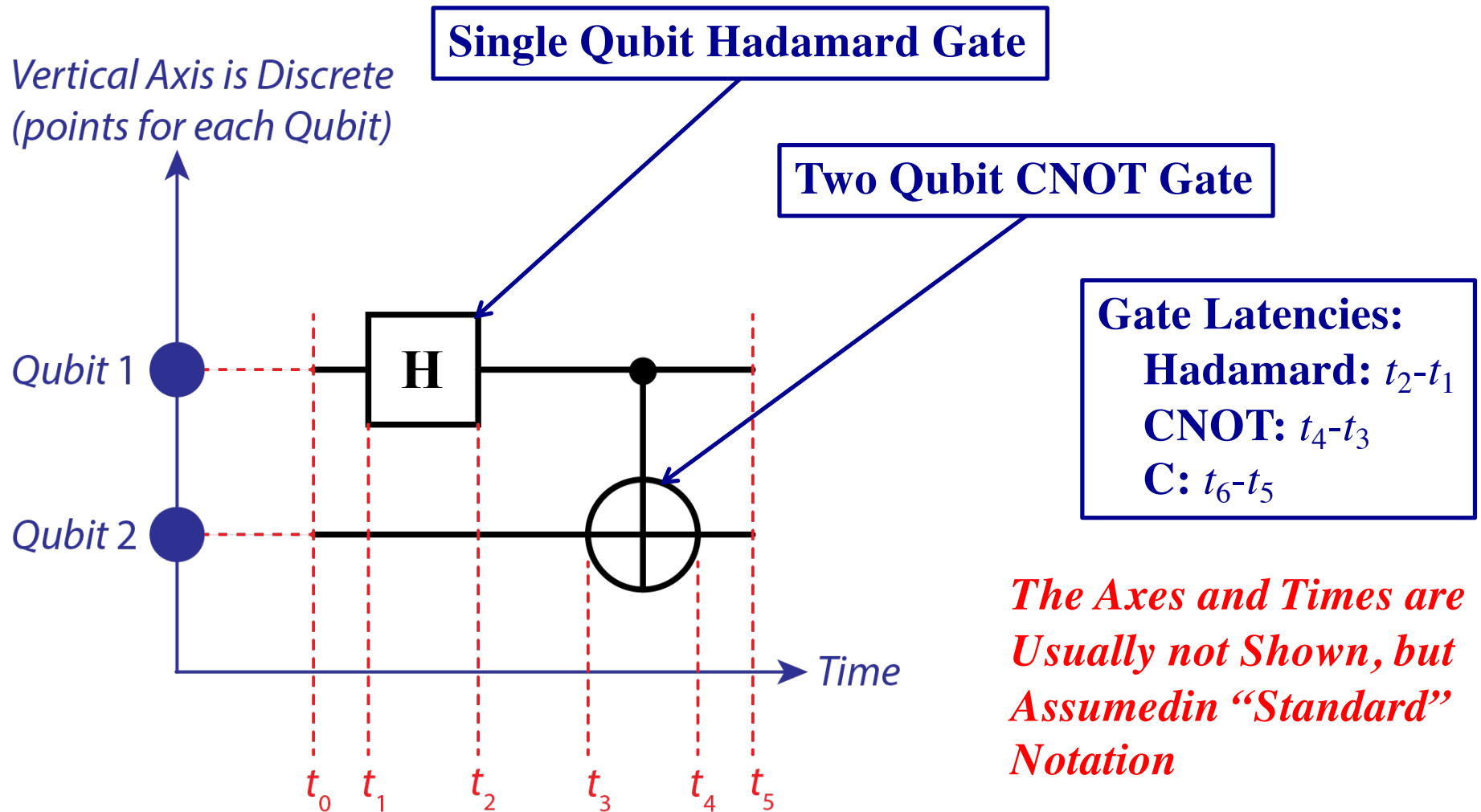
- Horizontal Lines Represent a Qubit
- Symbols that Vertically Span one or more Qubits are Quantum Gates, Operators, or QC Instructions

Vertical Axis is Discrete
(points for each Qubit)



Graphical Representation Example

- This is Well-Known “Bell State Generator”



Start Time:

$$t_0=0$$

End Time:

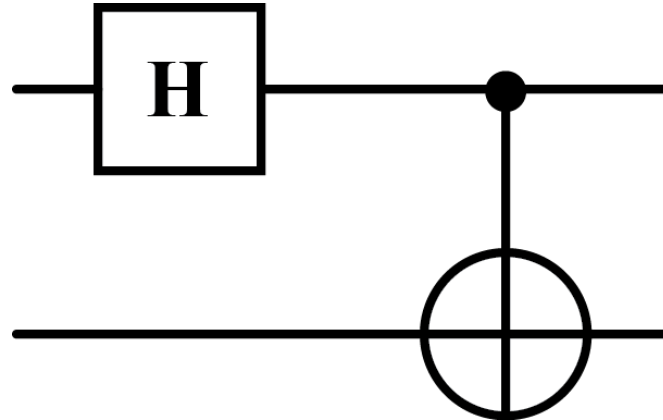
$$t_5$$

Overall Delay:

$$t_5-t_0=t_5$$

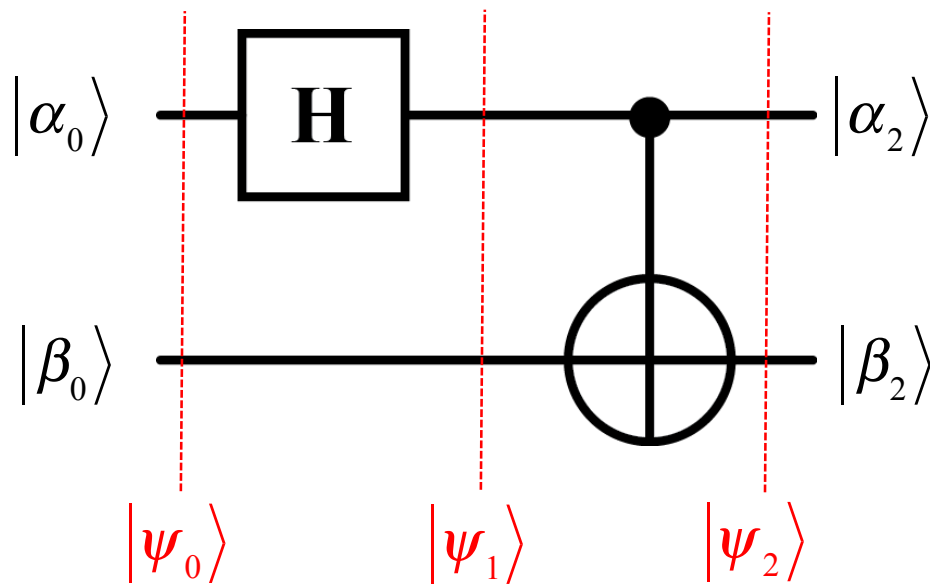
Graphical Representation Example (cont.)

- This is Well-Known “Bell State Generator”



- How is such a Circuit analyzed?
 - Qubit Values Represented as 2-dim. Column Vectors
 - Operators Represented as Linear Transformation Matrices
- Combining Qubits Accomplished with Outer Product Multiplication
- Combining Operators Accomplished with Direct Matrix Products
- Generally use Bra-Ket Notation for Conciseness

Bell State Generator Analysis

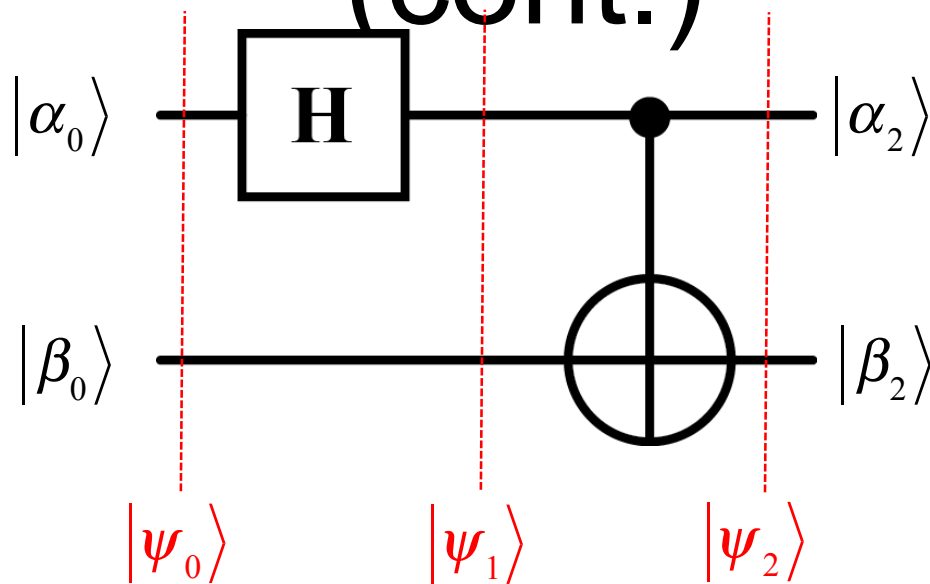


- I Prefer to Denote the Quantum State (Wave Function) at Discrete Points in Time with Dashed Lines
- STEP 1: Initialize Qubits at Time Zero to a Particular Value
 - Typically a “Ground” Basis State,

$$|\psi_0\rangle = |\alpha_0\rangle|\beta_0\rangle = |\alpha_0\rangle \otimes |\beta_0\rangle = |\alpha_0\beta_0\rangle = |00\rangle$$

Bell State Generator Analysis

(cont.)



- STEP 2: Calculate Evolved Quantum States at Each Denoted Time Instant

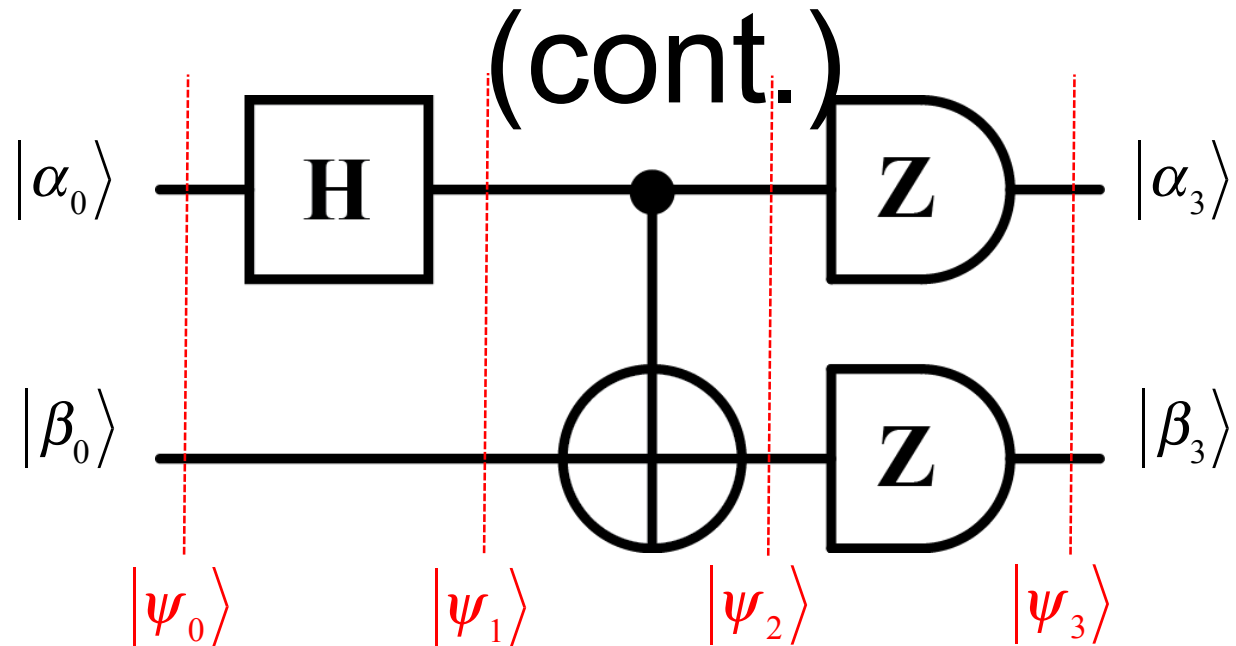
$$|\psi_2\rangle = \mathbf{C}_N |\psi_1\rangle \quad |\psi_1\rangle = (\mathbf{H} \otimes \mathbf{I}) |\psi_0\rangle \quad |\psi_2\rangle = \mathbf{C}_N |\psi_1\rangle = \mathbf{C}_N (\mathbf{H} \otimes \mathbf{I}) |\psi_0\rangle$$

- Overall Circuit/Program Transfer Matrix, \mathbf{U} , is:

$$\mathbf{U} = \mathbf{C}_N (\mathbf{H} \otimes \mathbf{I})$$

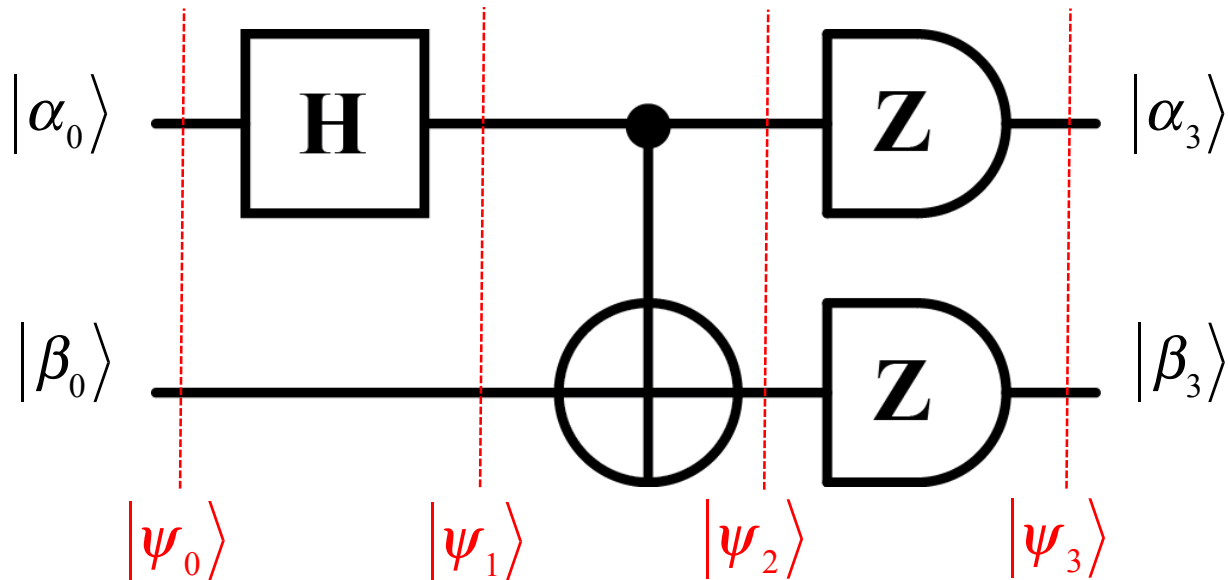
$$|\psi_2\rangle = \mathbf{U} |\psi_0\rangle$$

Bell State Generator Analysis



- STEP 3: Measure (Observe) the Overall Evolved Quantum State (using **Z** Detectors)
 - For Photonic Circuit, use Detectors at $|\psi_2\rangle$
- For this Example, $|\psi_2\rangle$ is an Entangled & Superimposed State. Measuring $|\psi_2\rangle$ is an Evolution Yielding $|\psi_3\rangle$
 - Probability of Measuring $|\psi_3\rangle = |00\rangle$ is 0.5
 - Probability of Measuring $|\psi_3\rangle = |11\rangle$ is 0.5

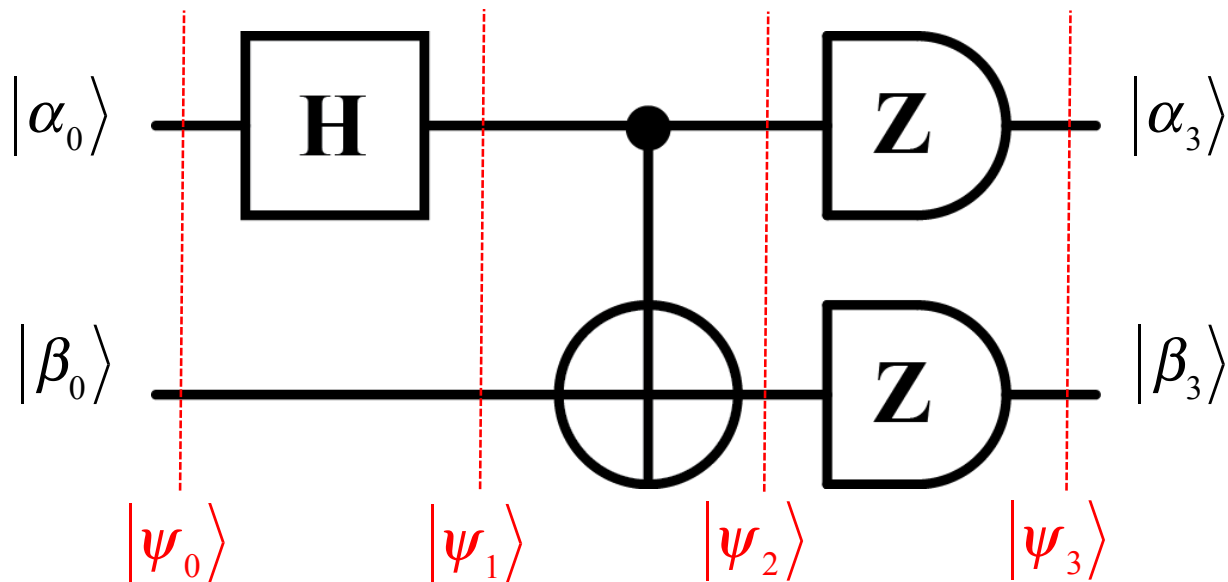
Numerical Example



$$|\psi_0\rangle = |\alpha_0\rangle|\beta_0\rangle = |\alpha_0\rangle \otimes |\beta_0\rangle = |\alpha_0\beta_0\rangle = |00\rangle$$

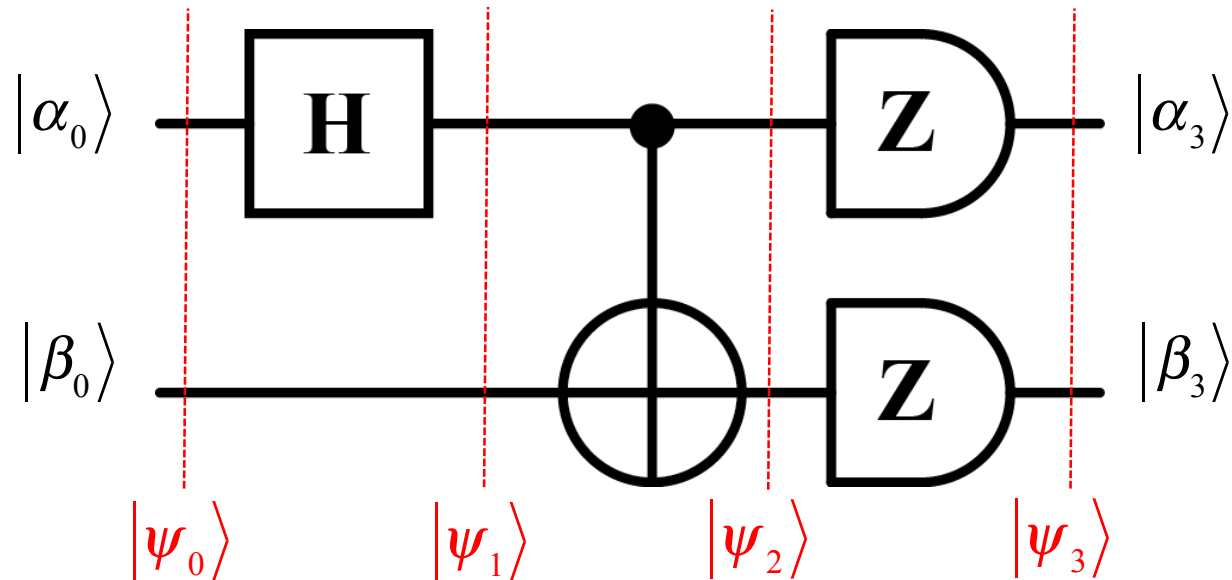
$$\mathbf{U} = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix} \left(\frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} \otimes \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right)$$

Numerical Example (cont.)



$$\mathbf{U} = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix} \left(\frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} \otimes \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \\
 = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix} \begin{bmatrix} 1 & 0 & 1 & 0 \\ 0 & 1 & 0 & 1 \\ 1 & 0 & -1 & 0 \\ 0 & 1 & 0 & -1 \end{bmatrix} = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 0 & 1 & 0 \\ 0 & 1 & 0 & 1 \\ 0 & 1 & 0 & -1 \\ 1 & 0 & -1 & 0 \end{bmatrix}$$

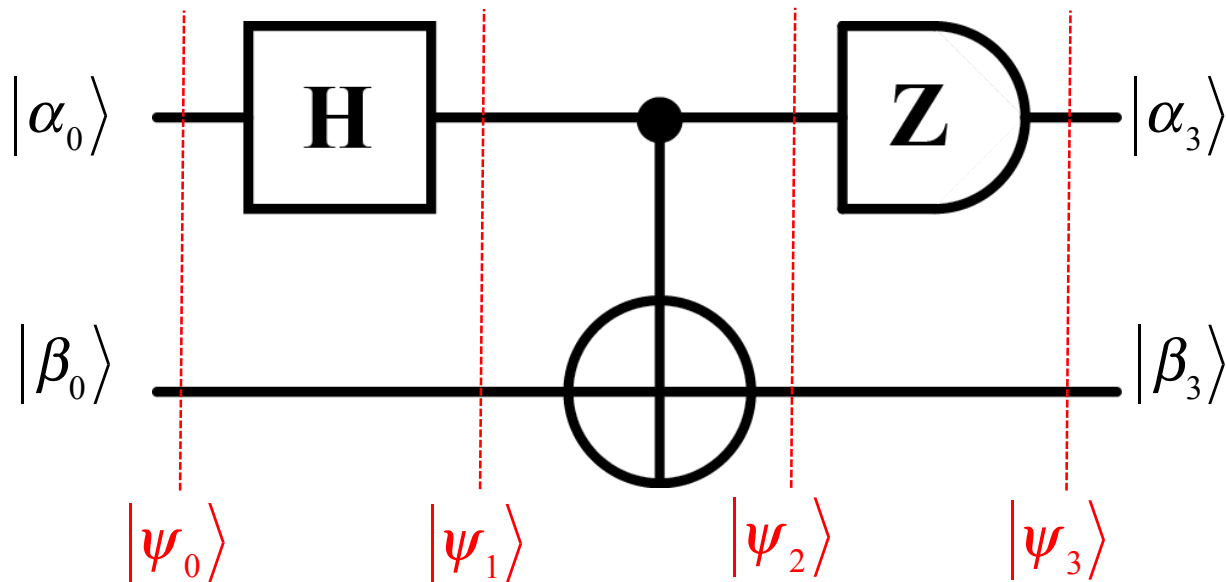
Numerical Example (cont.)



$$|\psi_2\rangle = \mathbf{U}|\psi_0\rangle = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 0 & 1 & 0 \\ 0 & 1 & 0 & 1 \\ 0 & 1 & 0 & -1 \\ 1 & 0 & -1 & 0 \end{bmatrix} \begin{bmatrix} 1 \\ 0 \\ 0 \\ 0 \end{bmatrix} = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 \\ 0 \\ 0 \\ 1 \end{bmatrix} = \frac{|00\rangle + |11\rangle}{\sqrt{2}}$$

- This is an Entangled Bell State since Impossible to Have $|\psi_3\rangle = |01\rangle$ or $|\psi_3\rangle = |10\rangle$
- Need Only Measure One Qubit to “Know” the Other

Numerical Example (cont.)



- Measurement with Respect to Computational Basis (**Z**-basis) Causes Probabilistic Collapse to One of the Eigenstates (i.e., eigenvectors of Pauli-**Z**)
- If $|\alpha_3\rangle = |0\rangle$, Then $|\beta_3\rangle = |0\rangle$ Also
- If $|\alpha_3\rangle = |1\rangle$, Then $|\beta_3\rangle = |1\rangle$ Also